

# $B \rightarrow \pi$ semileptonic form factors from Lattice QCD

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**Fermilab Lattice and MILC Collaborations**

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# Motivation

- Discrepancy between inclusive/exclusive determination of  $|V_{ub}|$ .  
Will it stay?
- Continuous experimental effort on  $B \rightarrow \pi \ell \nu$  decays .  
BaBar 1208.1253, Belle PRD83(071101)  
Expect  $\sim 4\%$  exp. error at Belle II. Browder, ANL snowmass
- Possible NP contribution in rare decays  $B \rightarrow K(\pi)\ell^+\ell^-$  (more by Andreas Kronfeld and Chris Bouchard).  
First observation of  $B^+ \rightarrow \pi^+ \mu^+ \mu^-$  by LHCb. 1210.2645  
But not seen in BaBar 1303.6010 and Belle 0804.3656v2 . Desy?
- **No** published update from LQCD since 2009. FNAL/MILC 2008

Quantity	CKM element	Present expt. error	2007 forecast lattice error	Present lattice error	2018 lattice error
$B \rightarrow \pi \ell \nu$	$ V_{ub} $	4.1%	–	8.7%	2%

# Introduction and Notations

$$\frac{d\Gamma}{dq^2} (\boxed{B \rightarrow \pi \ell \nu}) = \text{Phase space} \times |V_{ub}|^2 |f_+|^2$$

$$\begin{aligned} \langle \pi | \mathcal{V}^\mu | B \rangle &= f_+(q^2) \left( p_B^\mu + p_\pi^\mu - \frac{M_B^2 - M_\pi^2}{q^2} q^\mu \right) + f_0(q^2) \frac{M_B^2 - M_\pi^2}{q^2} q^\mu \\ &= \sqrt{2M_B} \left[ \underline{v^\mu f_\parallel(E_\pi) + p_\perp^\mu f_\perp(E_\pi)} \right] \end{aligned}$$

Convenient for lattice

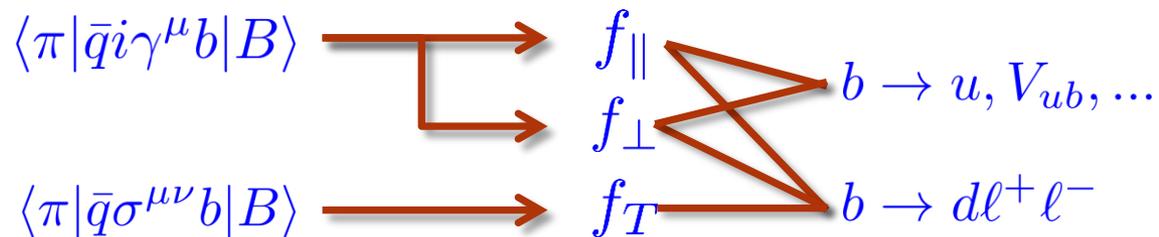
$$\frac{d\Gamma}{dq^2} (\boxed{B \rightarrow \pi \ell^+ \ell^-}) = \text{Phase space} \times |V_{tb} V_{td}^*|^2 \left\{ |C_9^{eff} f_+ - C_7 \frac{2m_b f_T}{M_B + M_\pi}|^2 + |C_{10} f_+|^2 \right\}$$

Aliev & Savci 1999

$$\langle \pi | \bar{d} \sigma^{\mu\nu} b | B \rangle q_\nu = - \frac{q^2 (p + k)^\mu - (M_B^2 - M_\pi^2) q^\mu}{M_B + M_\pi} f_T(q^2)$$

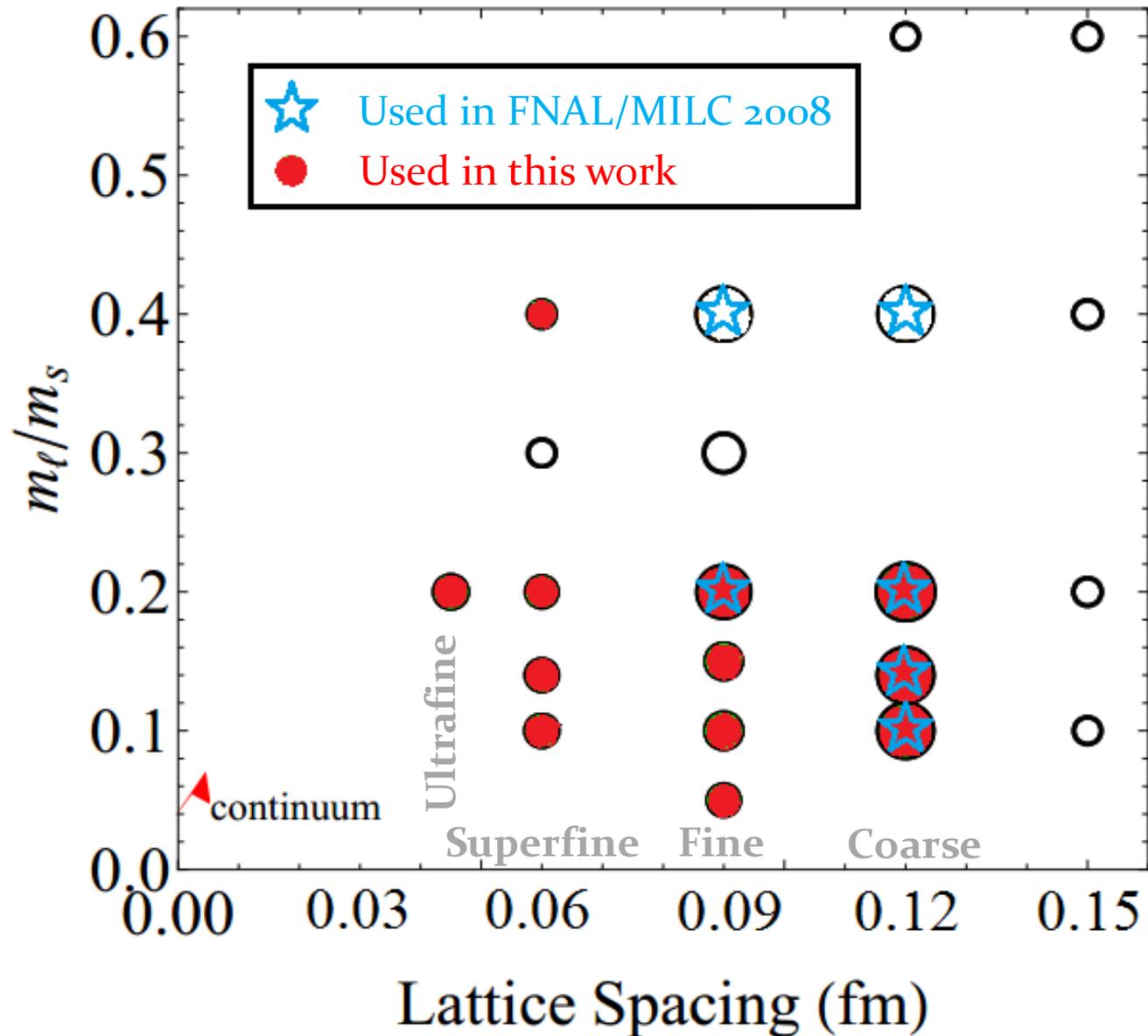
# Overview

- What do we calculate?



- What is **new** (compared to **FNAL/MILC 2008**)?
  - ✓ Major set of the new asqtad data (no  $u_0$  tuning error)
  - ✓ **12** ensembles vs 6
  - ✓ **4** lattice spacings vs 2
  - ✓ Mostly 3 times more configurations for the original ensembles still used
  - ✓ Better understanding of the simulation parameters and the systematic errors (heavy quark mass,  $Z_V$ , etc)
  - ✓ New analysis methods (due to improved statistics)

# MILC asqtad ensembles

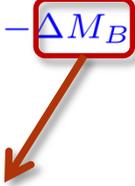


# Fitting correlation functions

- Method to extract the form factors  $f_{\parallel}, f_{\perp}, f_T$ 
  - Use the  $3pt/\sqrt{2pts}$  ratio to avoid using wave function overlaps. Iteratively average over two sink-source separations to suppress the contributions from oscillating states (opposite parity to the physical particles). **Fermilab-MILC 2008**

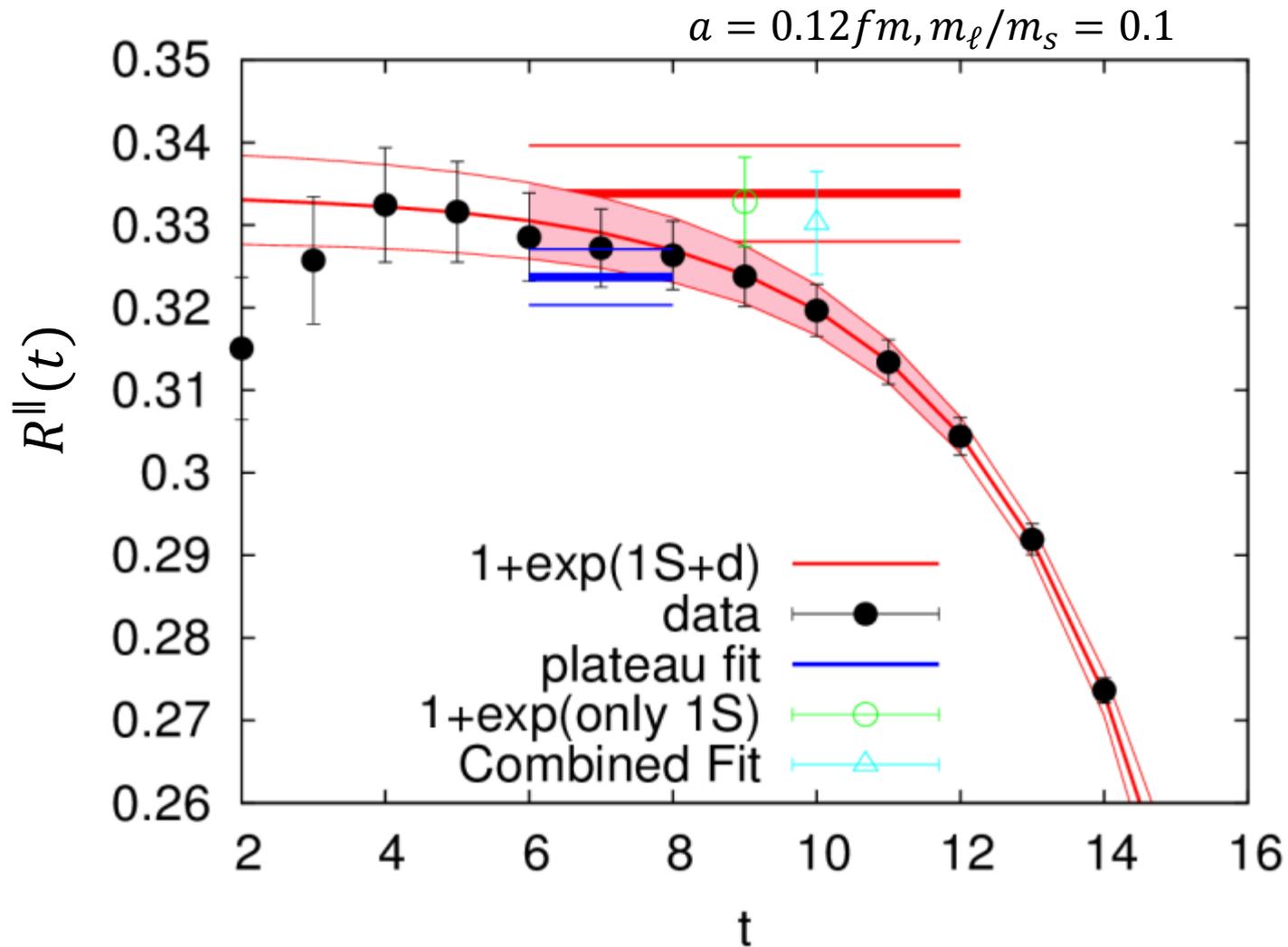
$$R(t) = \frac{\overline{C}_{3pt}(0, t, t_s; \mathbf{p}_{\pi})}{\sqrt{\overline{C}_{2pt}^{\pi}(t; \mathbf{p}_{\pi})\overline{C}_{2pt}^B(t_s - t)}} \sqrt{2E_{\pi} e^{E_{\pi,0}t + m_{B,0}(t_s - t)}}$$

- Fit to **plateau + exponential** function reflecting the B meson excited state contribution.

$$R_{(\parallel, \perp, T)}(t) = f_{(\parallel, \perp, T)}^{\text{LAT}} \times \left( 1 + \mathcal{A}_{(\parallel, \perp, T)} e^{-\Delta M_B(t_s - t)} \right)$$


- Combine the fit of ratio and B meson two-point function.

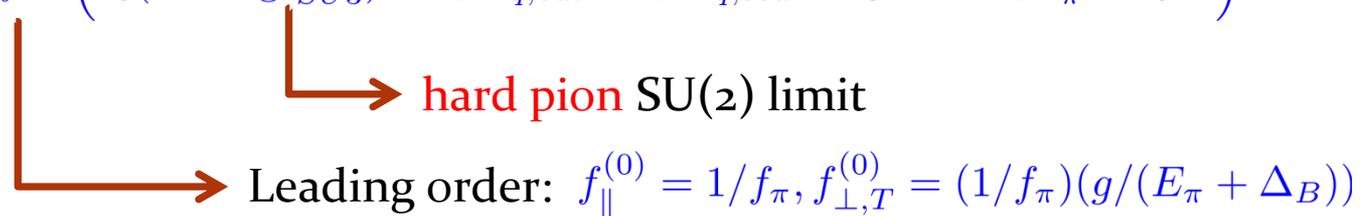
# Fitting correlation functions



# Chiral/continuum extrapolation

- The one-loop (NLO) HMs $\chi$ PT expansion

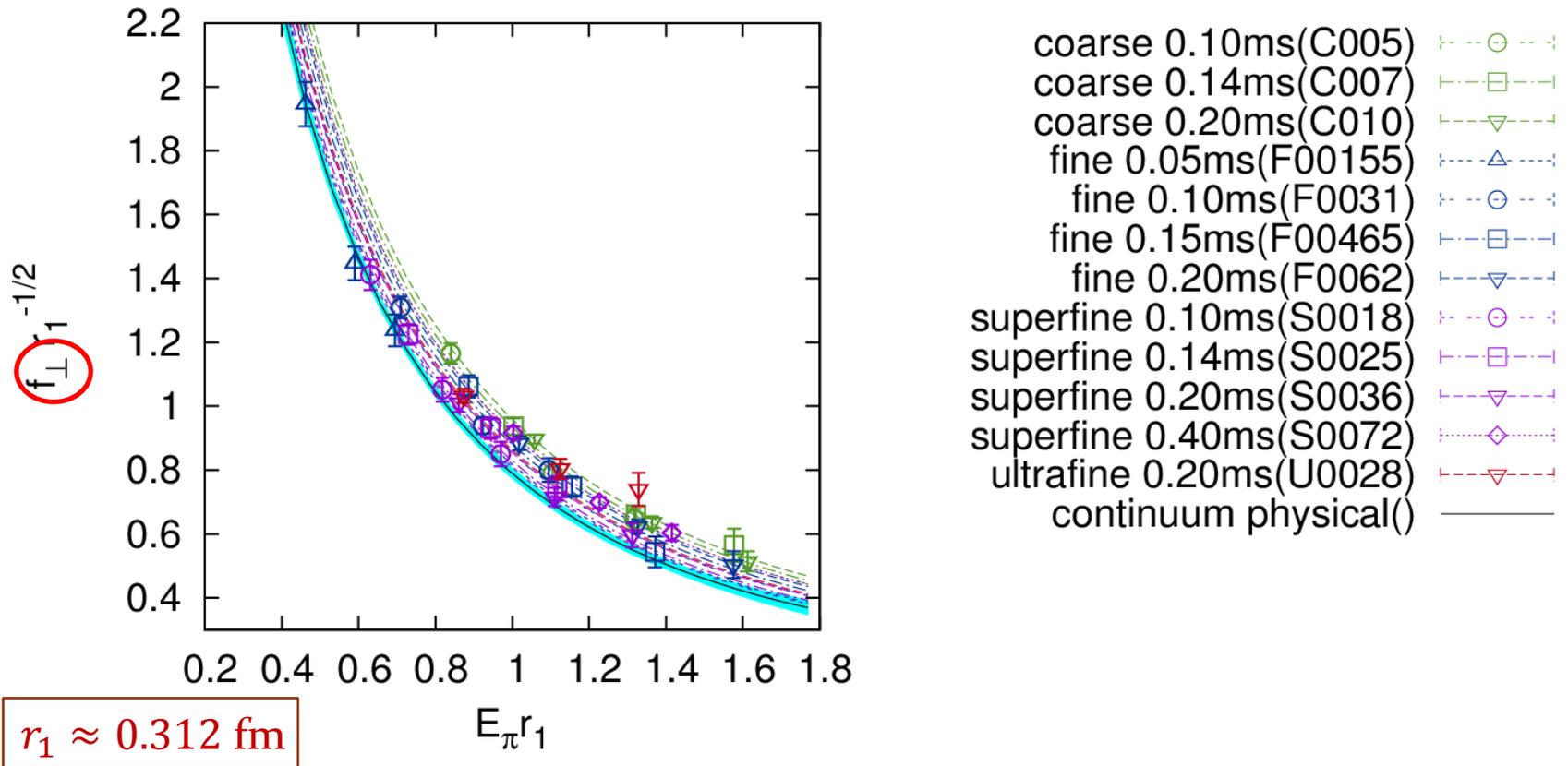
$$f = f^{(0)} \left( c_0(1 + \text{logs}_{SU3}) + c_1 m_{q, \text{val}} + c_2 m_{q, \text{sea}} + c_3 E_\pi + c_4 E_\pi^2 + c_5 a^2 \right)$$



- The non-analytic term  $\text{logs}$  are replaced by their **hard pion 1011.6531** and SU(2) limit, because it fails to describe  $f_{\parallel}$  data (even with NNLO analytic terms!)
- Pions in the simulation are too energetic. Hard pion  $\chi$ PT integrates them out. So the  $E_\pi$  dependence is essentially a phenomenological expansion.
- In hard pion limit,  $f_{\parallel}, f_{\perp}, f_T$  share the same non-analytic term.
- The variables are normalized with the  $\chi$ PT breaking scale  $\Lambda_\chi$ .
- The  $B$ - $B^*$ - $\pi$  coupling  $g = 0.45(8)$ . It has a small effect.

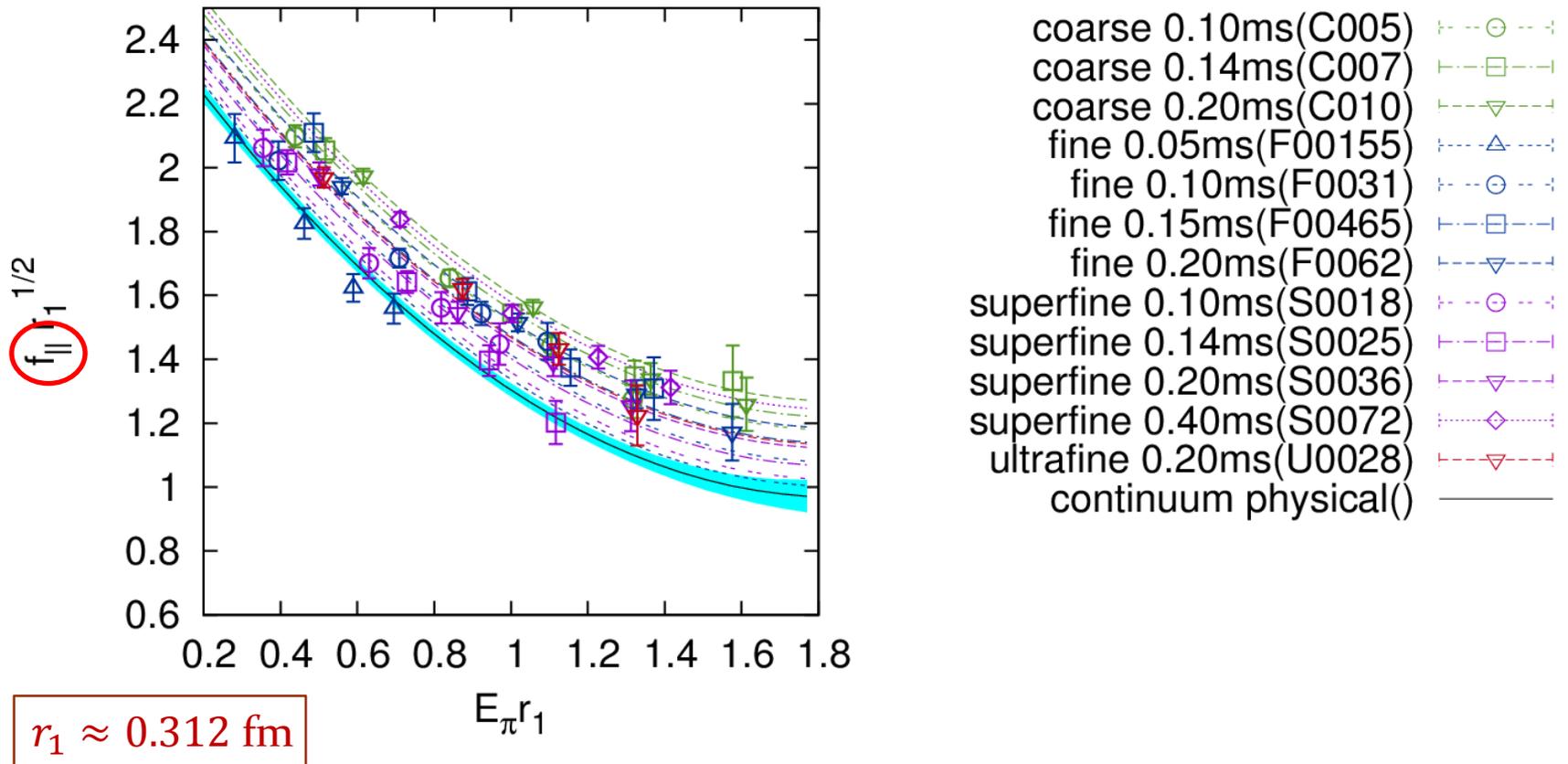
# Chiral/continuum extrapolation

chiral/continuum fit:  $\chi^2/[\text{dof}] = 0.94[36]$ ,  $p\text{value} = 0.56$



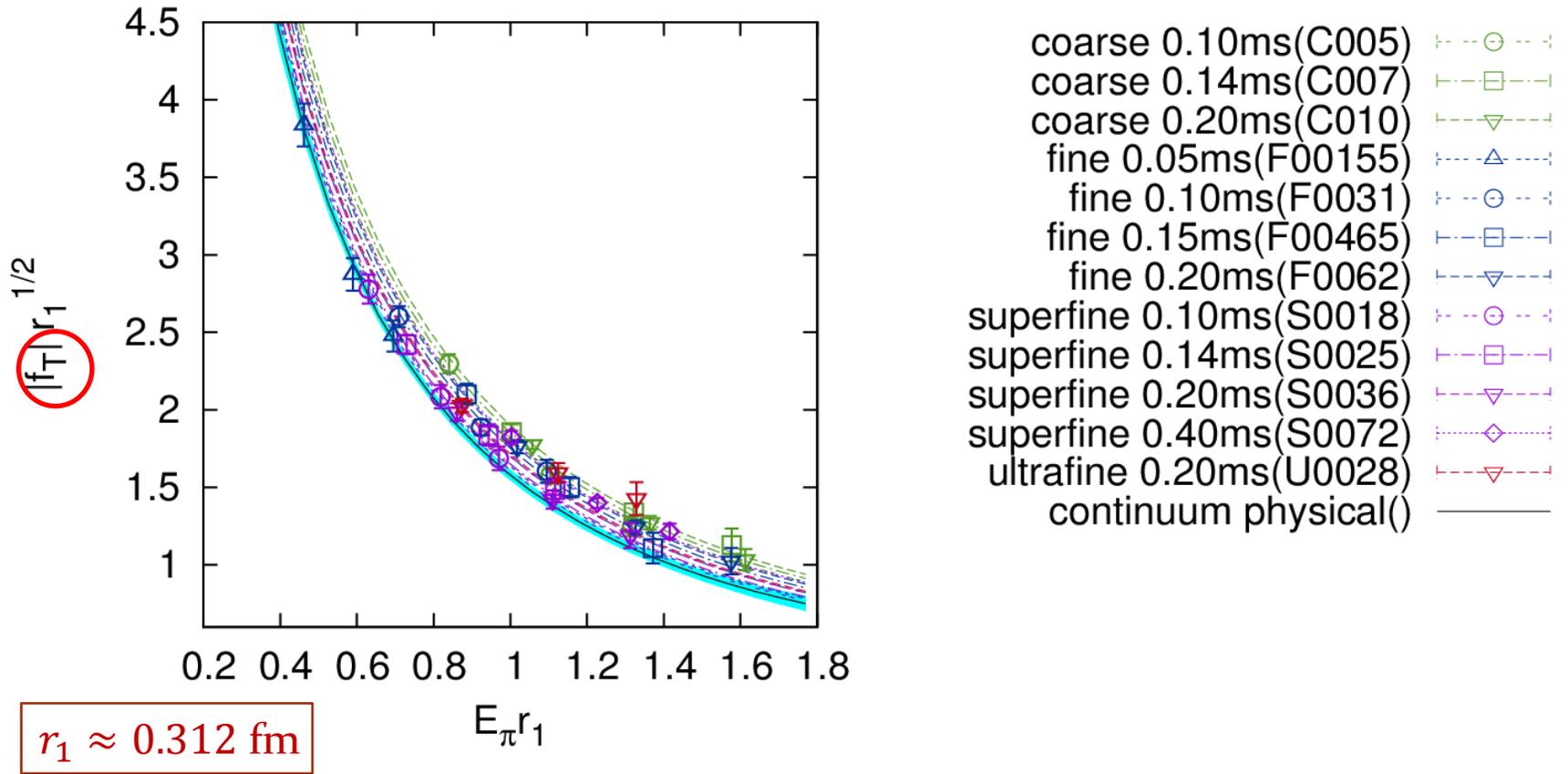
# Chiral/continuum extrapolation

chiral/continuum fit:  $\chi^2/[\text{dof}] = 1.2$  [48],  $p\text{value} = 0.2$

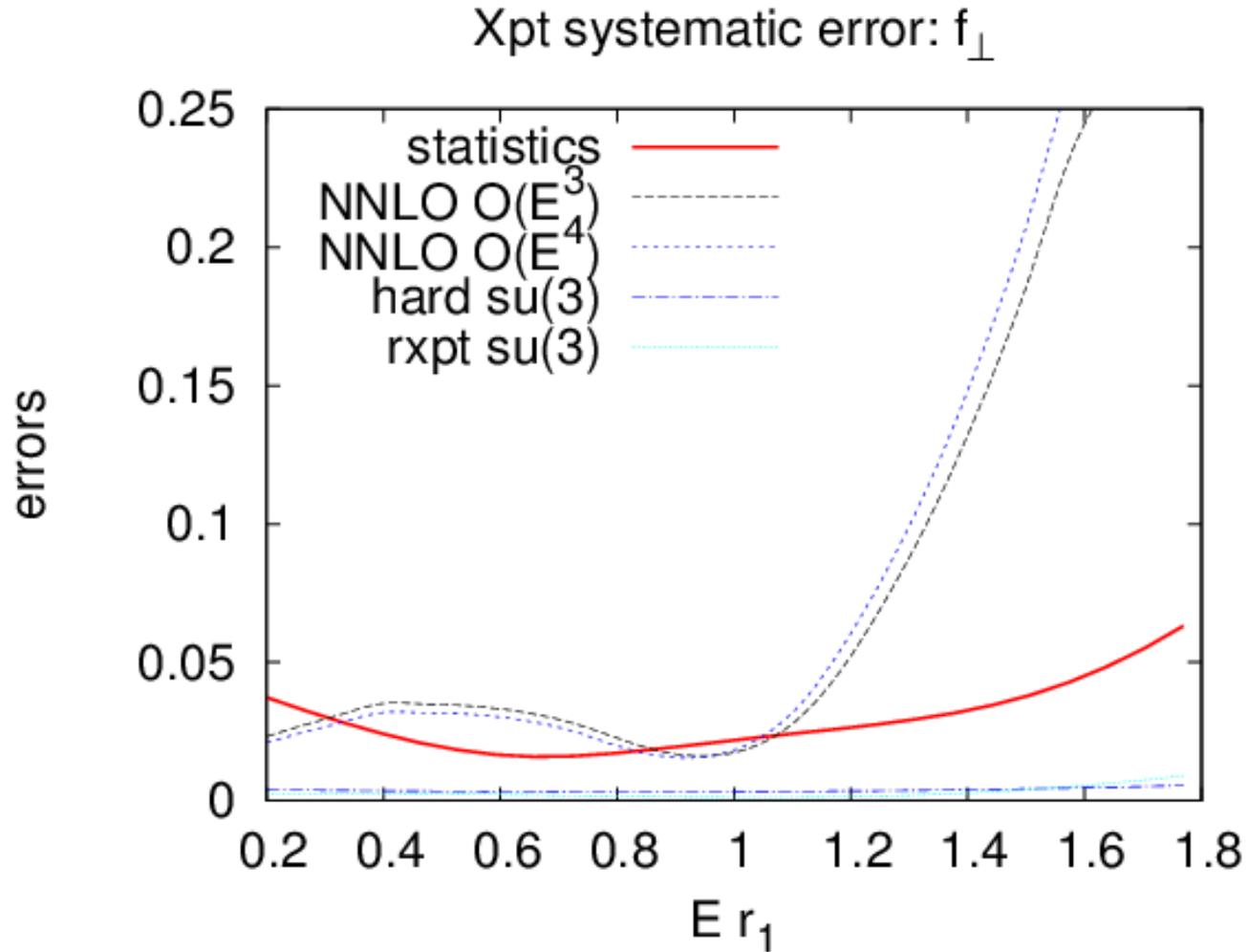


# Chiral/continuum extrapolation

chiral/continuum fit:  $\chi^2/[\text{dof}] = 0.8$  [36],  $p\text{value} = 0.8$



# $\chi$ PT systematic errors



# z-expansion

- Measured partial rates are more reliable in low- $q^2$  region, while lattice has better confidence in large- $q^2$  region. To determine  $|V_{ub}|$ , need lattice results in full- $q^2$  range.
- The model-independent z-expansion is more natural than the  $\chi$ PT expansion ( $|z| < 0.4$ ), plus important physical constraints (analyticity, unitarity, ...).
- A sequential fit (  $\chi$ PT/continuum  $\rightarrow$  z expansion). Correcting an ill expansion with a well behaving expansion. **Fitting a function variable to another function form!**
- **Synthetic data** approach: sampling the  $\chi$ PT results at “certain” locations  $\longrightarrow$  a regular curve fitting problem.

**Issues:** how are points taken (number and locations)?

Is correlation correctly accounted? Systematic error?

# z-expansion: a functional approach

- A generalization of data-point fitting to functional fitting.
- Covariance function (surface)  $K_f(s, t)$  of a function  $f(t)$  is a **Mercer kernel** (continuous, symmetry, positive semi-definite). By Mercer's theorem

$$K(s, t) = \sum_{i=1}^L \lambda_i \psi_i(s) \psi_i(t)$$

$$\mathcal{O}_K \psi_i = \lambda_i \psi_i$$

$$\mathcal{O}_K \phi(s) \equiv \int K(s, t) \phi(t) dt$$

If  $f(t)$  is a finite asymptotic expansion  $\rightarrow L$  finite.

(Pseudo) inverse of  $K(s, t)$ ,  $J(s, t)$

- Define an objective function (" $\chi^2$ ")

$$\begin{aligned} \mathcal{L} &= \int_{z_1}^{z_2} \int_{z_1}^{z_2} ds dt [f(s) - g_f(s)] J(s, t) [f(t) - g_f(t)] \\ &= \sum_{i=1}^L \frac{1}{\lambda_i} \left[ \int dt (f(t) - g_f(t)) \psi_i(t) \right]^2 \end{aligned}$$

"data"      Fit function



- Data point samples:  $\{x_i, y_i + \delta y_i\}$
- Target function:  $g(x; a_0, \dots)$
- Covariance matrix  $\text{Cov}$
- Objective function:  $\chi^2$

$$\chi^2 = \sum_{i,j} [y_i - g(x_i)] \text{Cov}_{ij}^{-1} [y_j - g(x_j)]$$

- Summation over all data point  $\{x_i\}$
- Principal decomposition:  
The eigen modes of  $\text{Cov}$

$$\text{cov } v_i = \lambda_i v_i$$

- Degrees of freedom:  $N - n$   
 $N$  - # of non-singular modes of  $\text{Cov}$   
 $n$  - # of fit parameters

- Function samples:  $\{x, f(x) + \delta f(x)\}$
- Target function:  $g_f(x; a_0, \dots)$
- Covariance function  $K(s, t)$
- Objective function:  $\mathcal{L}$

$$\mathcal{L} = \iint ds dt [f(s) - g_f(s)] K^{-1}(s, t) [f(t) - g_f(t)]$$

- Integral over the data range  $[a, b]$
- Principal decomposition:  
The eigen modes of  $K(s, t)$

$$\int dt K(s, t) \psi_i(t) = \lambda_i \psi_i(s)$$

- Degrees of freedom:  $L - n$   
 $L$  - # of non-zero eigenmodes of  $K(s, t)$   
 $n$  - # of fit parameters

# z-expansion: a functional approach

- Getting the covariance function  $K(s, t)$ .

$$f(t) = \sum_i c_i y_i(t) \quad K_f(s, t) = \sum_{i,j} y_i(s) \text{Cov}_{ij} y_j(t)$$

$\langle \delta c_i \delta c_j \rangle$



- z-expansion

$$z = \frac{\sqrt{t_+ - q^2} - \sqrt{t_+ - t_0}}{\sqrt{t_+ - q^2} + \sqrt{t_+ - t_0}}$$

$$t_+ = (M_B + M_\pi)^2$$

$t_0$  choice to be optimized

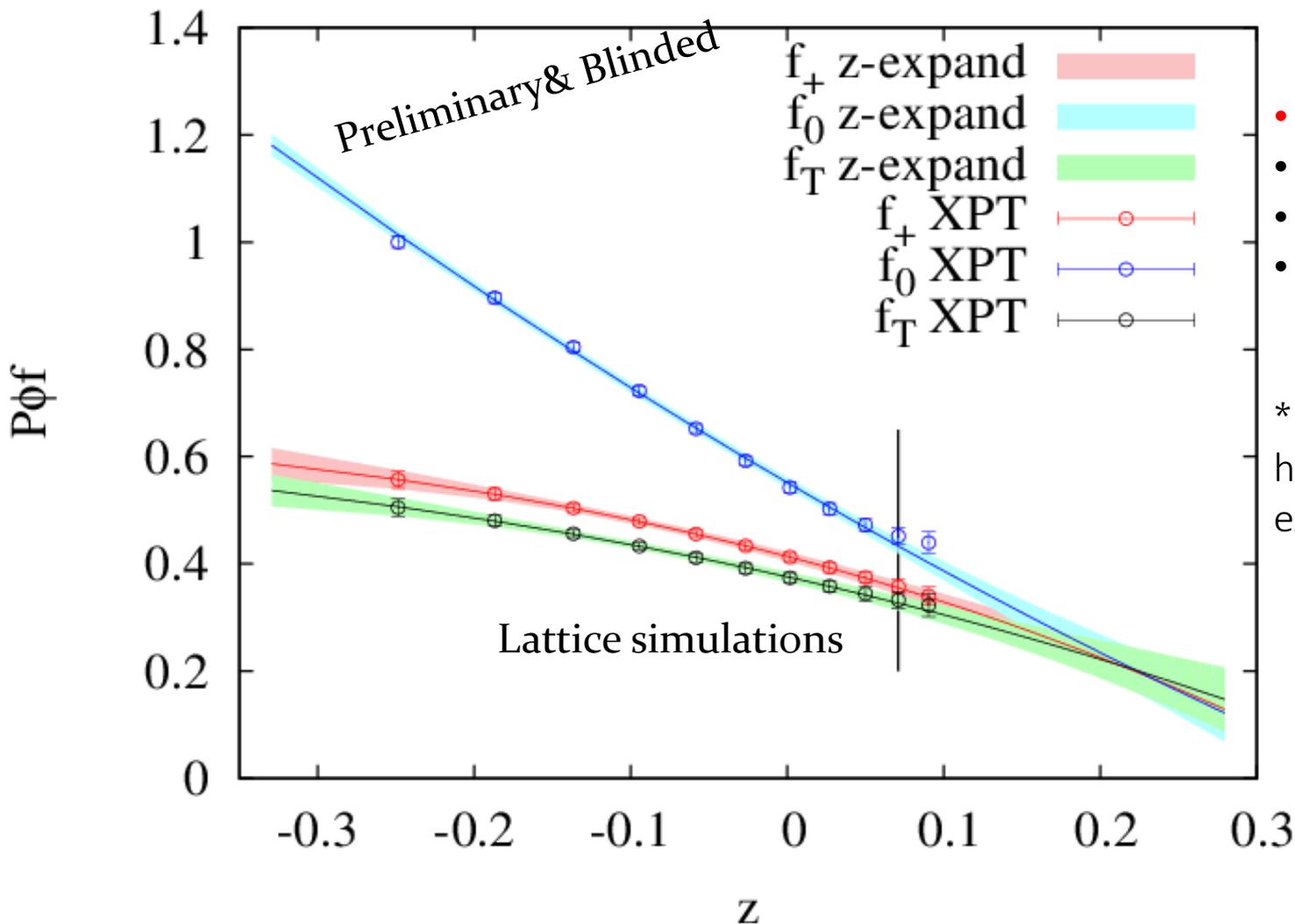
- BGL** (Boyd-Grinstein-Lebed) expansion [hep-ph/0504209](#)

$$g_f(z) = \frac{1}{P(z)\phi(z)} \sum_{n=0} a_n z^n$$

- BCL** (Bourenly-Caprini-Lellouch) expansion: correcting large  $q^2$  behavior [0807.2722v3](#)

$$g_f(z) = \frac{1}{1 - q^2/M_{B^*}^2} \sum_{n=0}^{K-1} a_n \left( z^n - (-1)^{n-K} \frac{n}{K} z^K \right)$$

# $z$ -expansion: a functional approach



- BCL expansion
  - Individual fits
  - Expansion up to  $z^3$
  - One singular mode is cut from  $K(s, t)$ .
- \* The expansion for  $f_0$  here is a polynomial expansion with no pole.

# Adding Constraints

- Use Lagrange multiplier to add bounds on expansion coefficients.

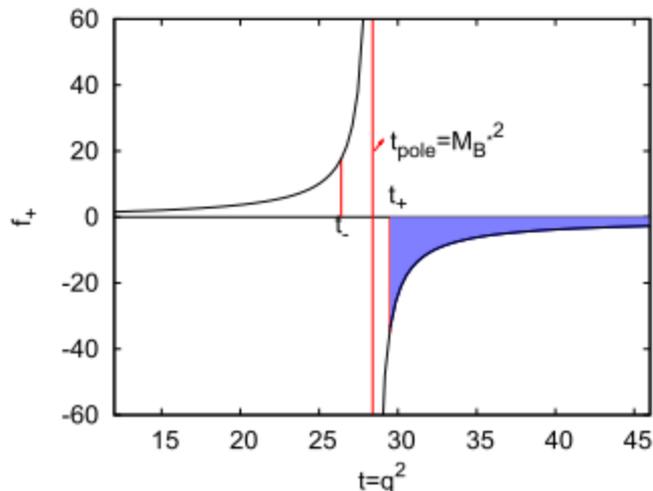
$$\sum_n |a_n|^2 \leq A \quad \text{BGL}$$

$$\sum_{m,n} B_{mn} b_m b_n \leq A \quad \text{BCL}$$

For **Unitarity** bound  $A = 1$ , which is usually not saturated (bound is loose).

- Analyticity** of form factors constrains the coefficients more tightly!

**Becher & Hill hep-ph/0509090**



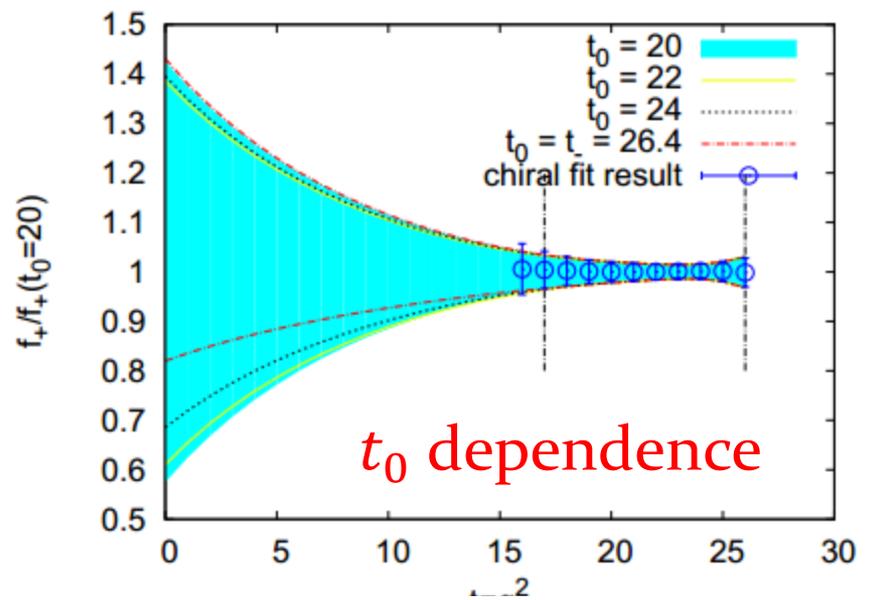
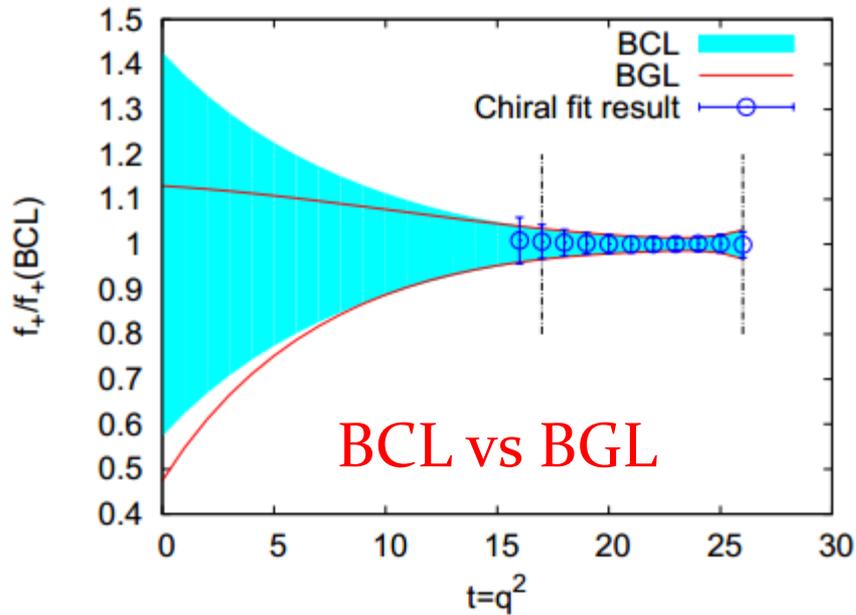
$$A_+ \sim \frac{m_b^2}{3} \int_{t_+}^{\infty} \frac{[(t - t_+)(t - t_-)]^{3/2}}{t^5} |f_+|^2$$

$$\sim 0.03$$

Replaced by leading order

A consistent check  
Important for truncation error estimate.

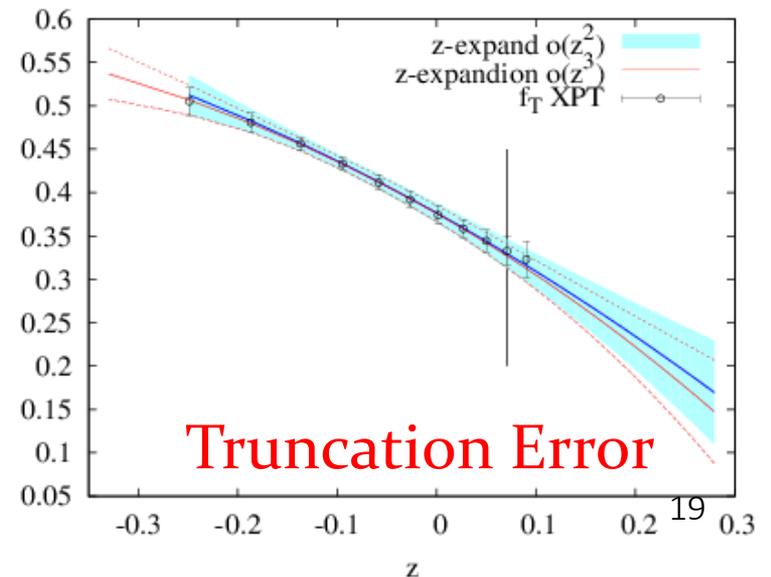
# z-expansion systematic errors



$[z_1, z_2]$	$[t_1, t_2]$	$\tilde{a}_0$	$\tilde{a}_1$	$\tilde{a}_2$	$f_+(q^2 = 0)$
[-0.25, 0.07]	[17,26]	0.414(8)	-0.78(7)	-0.8(5)	0.133(57)
[-0.19, 0.05]	[18,25]	0.412(9)	-0.74(7)	-0.6(6)	0.146(65)
[-0.14, 0.03]	[19,24]	0.413(9)	-0.74(7)	-0.6(5)	0.155(70)
[-0.09, 0.00]	[20,23]	0.411(9)	-0.78(7)	-0.5(6)	0.160(74)

Table I:

Integration Range



# Error Estimate

In the region with lattice data:

- Statistics: 2~4%
- $\chi$ PT systematic: 2~5% (the blowup is cured by z expansion)
- Kappa tuning: 0.5%
- z-expansion: < 1%
- Heavy quark discretization: ?%

# Summary and outlook

- We are working on an update to the lattice form factor calculations for semileptonic  $B \rightarrow \pi$  decays. We also include the tensor form factor in the calculation. The improvement on the error is promising.
- While  $\chi$ PT systematic is a dominant source of systematic errors in low- $q^2$  region, it can be improved by z-expansion.
- A new functional z-expansion method is used to reparameterize the  $\chi$ PT results.
- Full error budget is in progress.
- $|V_{cb}|$  and  $B \rightarrow \pi \ell^+ \ell^-$  prediction.