

On the N_f -dependence of gluonic observables

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ALPHA Collaboration

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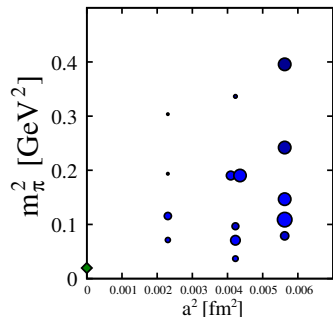
Lattice 2013 - Mainz

July 29, 2013



CLS Ensembles

β	ID	m_π [MeV]	MDU	stat/ τ_{exp}	$m_\pi L$	
5.2	A2	630	8000	121	7.7	
	A3	490	8032	121	6.0	
	A4	380	8096	122	4.6	
	A5	330	4004	163	4.0	
	B6	280	1272	52	5.1	
	5.3	E4	580	2784	10	6.1
E5f		430	16000	60	4.6	
E5g		430	16000	120	4.6	
F6		310	4800	36	4.9	
F7		260	9616	72	4.2	
G8		190	1114	23	4.0	
5.5		N4	550	6552	7	6.5
		N5	440	6208	7	5.2
	N6	340	8040	40	4.0	
	O7	260	3920	20	4.2	

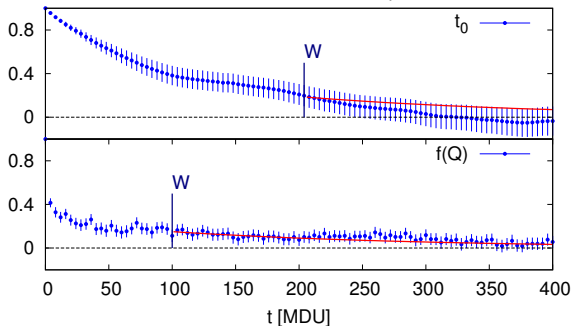




Errors and Autocorrelations

$$\sigma^2 = 2 \frac{\tau_{int}}{N} \Gamma(0), \quad \tau_{int} = \left(\frac{1}{2} + \sum_{t=1}^{W-1} \rho(t) \right) + \tau_{exp} \rho(W)$$

Normalized autocorrelation ρ

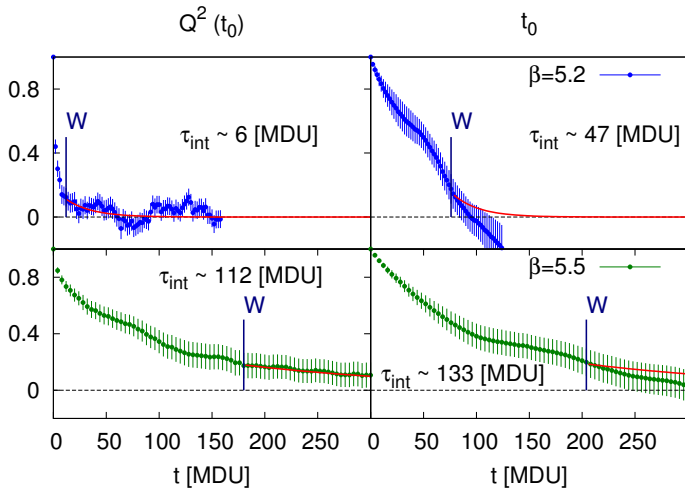


- ▶ $m_\pi = 340$ MeV
- ▶ $L = 2.3$ fm
- ▶ MDU ≈ 8000
- ▶ $stat/\tau_{exp} \approx 40$
- ▶ $f(Q)$:
 $\tau_{exp}\rho(W) \approx 50\%$
- ▶ t_0 : $\tau_{exp}\rho(W) \approx 30\%$

Definition of t_0 : $t_0^2 \langle E(t_0) \rangle = 0.3$, Wilson flow [\[Lüscher, '10\]](#)



Autocorrelations



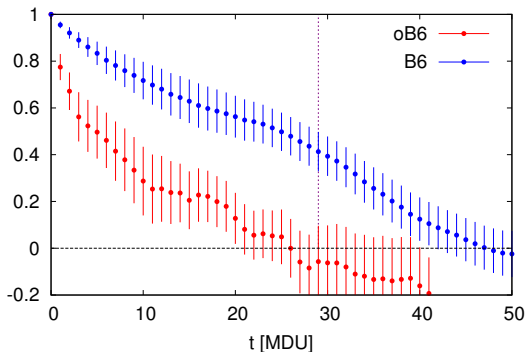
m_π	$stat$
MeV	τ_{exp}
280	50
340	40

Both in Q^2 and t_0 : tail contribution to τ_{int} is at most $\approx 30\%$



Open boundary conditions

Normalized autocorrelation ρ of t_0



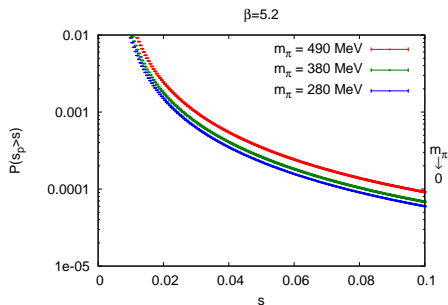
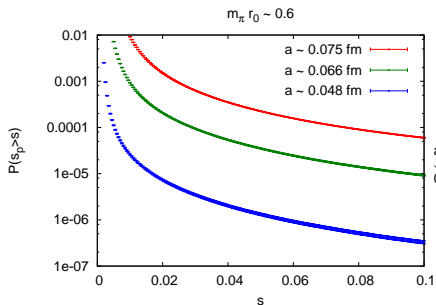
	B6	oB6
β	5.2	5.2
κ	0.13597	0.13597
MDU	1272	1000
L	48	48
T	96	192
plateau	1...96	40...156
t_0	3.3292(80)	3.3317(54)

τ_{int} reduced by factor ≈ 2



Separation between sectors

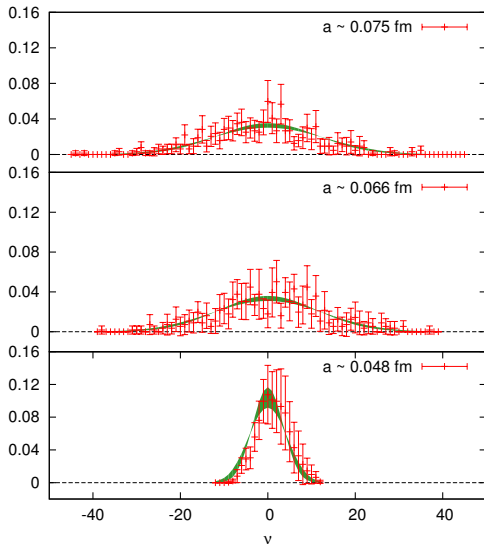
- ▶ Defining $h = \max s_p$, with s_p value of plaquette p at t_0
- ▶ configurations “between the sectors” characterized by $h > 0.067$ [Lüscher, '10]
- ▶ occur less when reducing a or m_π :



With dynamical fermions $P(s_p) \sim (a/r_0(m))^{10}$



How well is the topological charge Q sampled?



Q : top. charge measured at t_0

$$P_\nu = \frac{e^{-\nu^2}}{\sqrt{2\pi Q^2}} (1 + O(V^{-1}))$$

[Giusti, Lüscher, Weisz, Wittig, '03]

1. $m_\pi = 280$ MeV, $L = 3.6$ fm
2. $m_\pi = 190$ MeV, $L = 4.2$ fm
3. $m_\pi = 340$ MeV, $L = 2.3$ fm

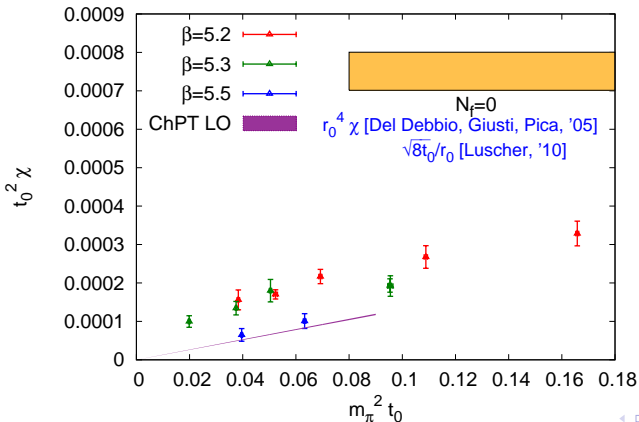
Reasonable sampling of charge distribution



Susceptibility

We computed $t_0^2 \chi = t_0^2 \frac{\langle Q^2 \rangle}{V}$ and we compare it with

$$\text{ChPT LO} : \chi = \frac{m}{2} \Sigma (1 + O(m)) = \frac{1}{8} f_\pi^2 m_\pi^2 (1 + O(m_\pi^2))$$



Input:

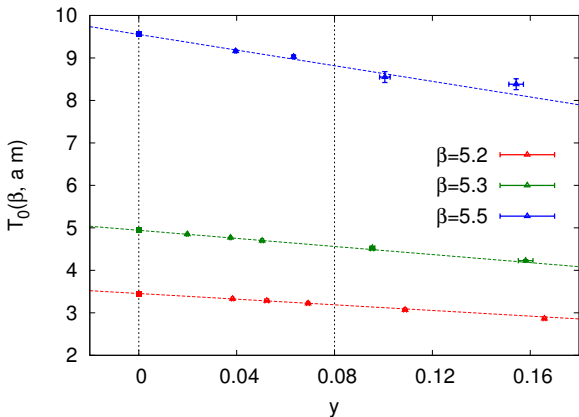
$$f_\pi = 130.4 \text{ MeV}$$

$$t_0 = 0.02396(37) \text{ fm}^2$$

Suppression by fermion determinant is clearly seen



t_0 : chiral extrapolation



Definitions:

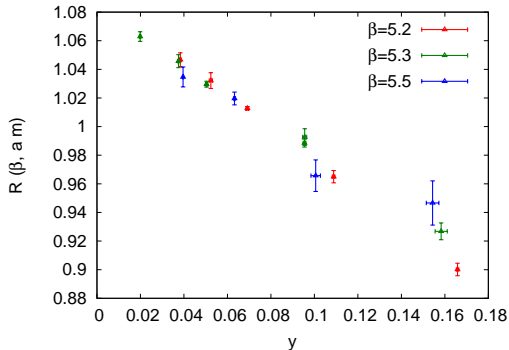
- ▶ $am = am_{PCAC}$
- ▶ $T_0(\beta, am) = \frac{t_0}{a^2}$
- ▶ $M_\pi(\beta, am) = am_\pi$
- ▶ $y \equiv T_0 M_\pi^2$

$$\text{▶ } T_0(\beta, am_{12}) = A(\beta)(1 + By) \rightarrow \left(\frac{t_0}{a^2}\right)^{chiral} \equiv T_0(\beta, 0) = A(\beta)$$



Remaining $O(a)$ effects?

- Does B depend on a : $B(a) = B + O(a)$?



$$R(\beta, am) = \frac{T_0(\beta, am)}{T_0(\beta, am)|_{y=0.08}}$$

We see neither a nor a^2 effects with few per-mille precision

Symanzik effective theory $\rightarrow \tilde{g}_0 = g_0(1 + b_g am_q)$

- ▶ $b_g = (0.012 \times N_f)g_0^2 + O(g_0^4)$ in PT
- ▶ we used $b_g = 0$ because of the smallness of 1-loop effects and it is confirmed to be ok non-perturbatively



t_0 : systematics

We checked the extrapolations are stable by:

- ▶ cutting the pion mass

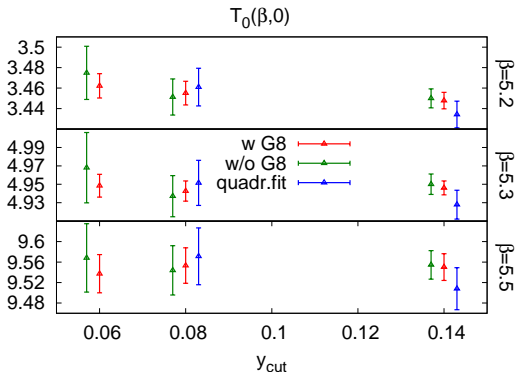
$$m_\pi < \text{cut}$$

cut/MeV	y_{cut}
520	0.14
400	0.08
330	0.06

- ▶ cutting the pion mass

$$m_\pi > 190 \text{ MeV}$$

- ▶ adding higher terms
(e.g. $O(y^2)$)

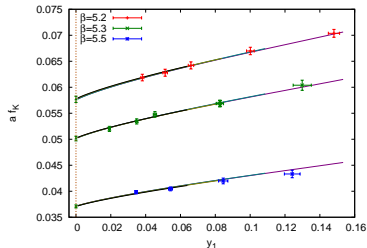




t_0 : continuum limit

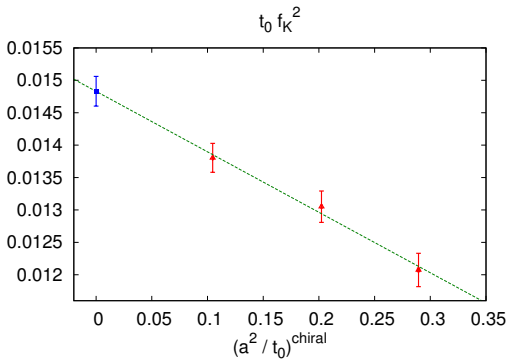
Input:

- ▶ af_K at phys. quark masses,

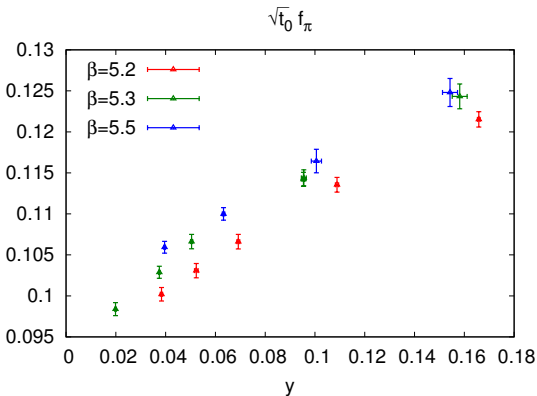


see [S. Lottini's talk](#)

- ▶ $f_{K,phys} = 155\text{MeV}$



$$t_0 = \frac{\lim_{a \rightarrow 0} (t_0 f_K^2)}{f_{K,phys}^2} = 0.02396(37) \text{ fm}^2$$

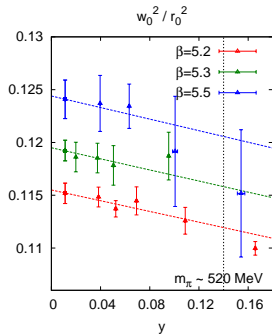
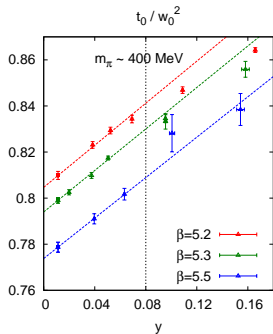
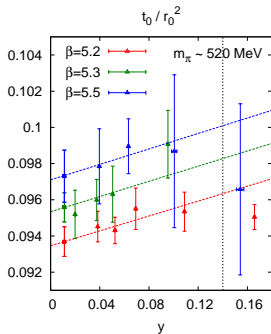
 t_0 : possible uses

- ▶ small discretisation effects
- ▶ more precise extrapolations w.r.t. $r_0 f_\pi$
- ▶ see [S. Lottini's talk](#)



Dynamical fermion effects in gluonic observables

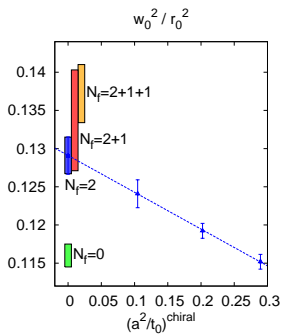
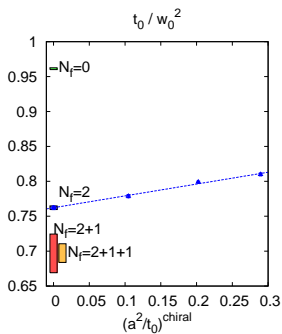
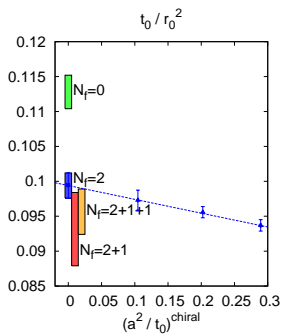
- ▶ 4% discretis. effects
- ▶ weak quark mass dependence
- ▶ 5% discretis. effects
- ▶ **stronger** quark mass dependence
- ▶ 9% discretis. effects
- ▶ weak quark mass dependence



Definition w_0 : $t \frac{d}{dt} t^2 \langle E(t) \rangle \Big|_{t=w_0^2} = 0.3$ [BMW, '12]



N_f dependence



- $N_f = 0$ Lüscher, '10
- $N_f = 2 + 1$

Quantity [fm]	ref.
$r_0 = 0.480(10)(4)$	RBC, '12
$\sqrt{t_0} = 0.1465(21)(13)$	BMW, '12
$w_0 = 0.1755(18)(4)$	BMW, '12

- $N_f = 2 + 1 + 1$

Quantity	ref.
$r_0/r_1 = 1.508$	HotQCD, '11
$\sqrt{t_0}/w_0 = 0.835(8)$	HPQCD, '13
$r_1/w_0 = 1.790(25)$	HPQCD, '13



Conclusions

Simulations with $19 < stat/\tau_{exp} < 165$:

- ▶ autocorrelations effects just under control
- ▶ open bc help even at largest lattice spacing
- ▶ weak quark mass dependence of t_0 , and r_0

$$B_{t_0} = -0.96(5), \quad B_{r_0} = -0.7(2)$$

- ▶ weaker quark mass dependence of t_0/r_0^2

$$B_{t_0/r_0^2} = 0.22(16)$$

- ▶ dynamical fermion suppression of topology clearly seen
- ▶ ... but still a long way to quantitative understanding of topology