Testing reweighting method for truncated Overlap fermions

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Plan

- 1. Truncated overlap fermion
- 2. Reweighting
- 3. Implementation
- 4. Results
- 5. Summary

- Overlap fermion
 - Lattice chiral symmetry (Ginsparg-Wilson relation)

$$D_{OV}\gamma_5 + \gamma_5 D_{OV} = 2aD_{OV}\gamma_5 D_{OV}$$

Hasenfraz, Laliena, Niedermayer, PLB427(1998) Lusher, PLB428(1998)

$$\psi' = \psi + i\varepsilon\gamma_5(1 - aD_{OV})\psi \qquad \Leftrightarrow \overline{\psi}' D\psi' = \overline{\psi}D\psi$$

$$\overline{\psi}' = \overline{\psi} + i\varepsilon\overline{\psi}(1 - aD_{OV})\gamma_5$$

• Explicit contruction (Neuberger overlap Dirac fermion)

$$D_{OV} = \frac{1}{2} (1 + \gamma_5 \text{sign}(H_W)), \quad H_W = \gamma_5 (D_W - m_{dw})$$

- negative mass (0< m_{dw} <2) Wilson kernel D_W
- Matrix Sign function



• Dynamical Overlap fermion

- HMC with dynamical overlap fermion
- Difficulty in the sign function: one needs to keep track of the low-modes of H_w (refrection/reflaction)

$$D = \frac{1}{2} \left(1 + \gamma_5 \operatorname{sign}(H_W) \right), \quad H_W = \gamma_5 (D_W - m_{dw})$$

Fodor,Katz,Szabo, JHEP 08(2004),NPB ProcSuppl 140(2005),JHEP 01(2006), DeGrand, Shafer, PRD71(2005),PRD72(2005),PoS(Lat05), Cundy, Krieg, Arnold, Frommer, Lippert, Schilling, NPB ProcSuppl 140(2005), Cundy, Krieg, Lippert PoS(Lat05).

- Topology fixing
 - avoid/suppress appearance of near zero mode of H_w
 - by adding extra terms or by using an admissible gauge action.

$$S_{G} = \begin{cases} \beta \sum \frac{1 - \operatorname{Re} Tr[P_{\mu\nu}]/3}{1 - (1 - \operatorname{Re} Tr[P_{\mu\nu}]/3)/\varepsilon} & (\text{for } 1 - \operatorname{Re} Tr[P_{\mu\nu}]/3 < \varepsilon) \\ \infty & \text{Luscher, NPB549 (1999)} \\ \text{Izubuchi, Dawson (RBC) NPB ProcSuppl 106(2002)} \\ \text{Bietenholz, Jansen, Nagai, Necco, Scorzato, Shcheredin, JHEP 0603 (2006)} \end{cases}$$

$$\det[H_W^2 + \mu^2]$$

Vranas, PRD74(2006)

JLQCD: Fukaya et al. PRD74(2006),

Lattice 2013@Makenfrew, Blum, Christ, Mawhinney, Vranas, PoS(Lat08)

Fukaya, Hashimoto, Hirohashi, Ogawa, Onogi; PRD73(2006)

- Approximate (truncated) Overlap fermion
 - Give up exact Lattice chiral symmetry
 - Keep Better Chiral property than that of Wilson type.
 - Domain wall fermion with finite N5.
 - Truncate the approximation for the Sign function.
- Overlap fermion construction (approximation for sign function)
 - Domain wall type

$$D_{TROV} = \varepsilon^T P^{\dagger} \left[\left(D_{DWF(N_5, m_f=1)} \right)^{-1} D_{DWF(N_5, m_f)} \right] P \varepsilon \underset{N_5 \to \infty}{\longrightarrow} m_f + (1 - m_f) D_{OV}$$

Kikukawa-Noguchi'99, Borici'99,Chiu'03, Shamir,Brower'05,....

 $D_{TROV} = 1 + \gamma_5 H_W \left(\alpha_0 + \sum_{i=1}^N \frac{\alpha_i}{H_W + \beta_i} \right) \xrightarrow[N \to \infty]{} 2D_{OV}$

Partial fraction type

JLQCD, TWQCD,....

Continued fraction type

$$D_{TROV} = 1 + \gamma_5 (\beta_0 H_W + \frac{\alpha_1}{\beta_1 H_W + \frac{\alpha_2}{\beta_1 H_W + \frac{\alpha_3}{\vdots \vdots \frac{1}{2N} \alpha_{2N}}}}) \xrightarrow{N \to \infty} 2D_{OV}$$

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Eigenvalues of free Overlap fermion operator on 16⁴ with Nucberger kernel type with a constraint of the overlap fermion operator on 16⁴ 0.5 0.4 0.3 0.2 0.1 0 0 0.1 0 0 0.1 0.2 0.3 0.2 0.1 0 0 0.5 0.4 0.5

0.7 0.8 0.9

0.3 0.4 0.5 0.6

- Here I use Domain wall type realization of overlap fermion in the HMC algorithm.
- This will smoothen the sign function. => Lowering the cost of HMC?, Topology changing?
 - c.f. Nf=2 Dynamical Domain wall simulation:
 Y.Aoki et al., PRD72(2005)

S.Shaefer, PoS(Lat06)

- Lattice Chiral symmetry is explicitly broken. But it should be better than Wilson type fermions.
- Recovering Chiral symmetry at the measurement step by Reweighting method.



Trunc'd OV with DWF kernel(N5=4)



2. Reweighting

- We consider QCD (Nf=2) $Z = \int DU \det [D_{OV}]^2 e^{-S_G[U]}$ $= \int DU \operatorname{det} \left[D_{OV} / D_{TROV(1)} \right]^2 \operatorname{det} \left[D_{TROV(1)} / D_{TROV(2)} \right]^2 \cdots$ $\cdots \det \left[D_{TROV(Nstep-1)} / D_{TROV(Nstep)} \right]^2 \det \left[D_{TROV(Nstep)} \right]^2 e^{-S_G[U]}$ $= \int DUD\phi^{\dagger} D\phi \prod^{Nstep} W_{(J)} \exp\left[-S_G[U] - |D_{TROV(Nstep)}|^2\right]$ $D_{TROV(0)} \equiv D_{OV}$: exact sign function (numerically) : truncated overlap op. (at a finite N_5) $D_{TROV(J)}$ $W_{(I)} = \det \left[D_{TROV(I-1)} / D_{TROV(I)} \right]^2$: reweighting factor (evaluated via noisy method)
 - average over ensemble generated with the truncated overlap fermion.

$$\left\langle O \right\rangle = \frac{\int DUO[U] \det[D_{OV}]^2 e^{-S_G[U]}}{\int DU \det[D_{OV}]^2 e^{-S_G[U]}} = \left\langle O \prod_{j=1}^{Nstep} W_{(J)} \right\rangle_{TROV(Nstep)} / \left\langle \left\langle \prod_{j=1}^{Nstep} W_{(J)} \right\rangle_{TROV(Nstep)} \right\rangle_{TROV(Nstep)}$$

$$\left\langle X \right\rangle_{TROV(Nstep)} = \frac{\int DUX[U] \det[D_{TROV(Nstep)}]^2 e^{-S_G[U]}}{\int DU \det[D_{TROV(Nstep)}]^2 e^{-S_G[U]}}$$
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2. Reweighting

- Reweighting for overlap/Domain wall fermions
 - for Domain wall, Chirality improvement : T.Ishikawa, Y.Aoki, T.Izubuchi(RBC-UKQCD), PoS(Lat09)035.
 - for Overlap in epsilon regime: T.DeGrand, PRD78(2008)117504.
 - for Overlap strange quark: H.Ohki et al.(JLQCD), arXiv:1208.4185[hep-lat]
 - for Domain wall, change mass: Q.Liu et al (RBC-UKQCD), PRD87(2013)054503.
 - In this conference
 - 8/2(Fri) 14:20-, RoomD: H.Fukaya et al.,

"Overlap/Domain-wall reweighting"

- Other reweighting related talks:
 - (Mon), 18:30, Room C, W. Freeman et al.
 - (Fri), 15:00, Room B, C. Aubin et al.
 - (Mon), 16:30, RoomD, *B.Leder et al.*
 - (Mon), 16:50, Room D, J. Finkenrath et al.
 - Poster, S. Schaefer

- HMC with truncated overlap fermion (Nf=2)
 - Domain wall operator realization

A.Borici, hep-lat/0402035 A.Allkoci, A.Borici, PoS(Lat05)099

$$D_{TROV} = \varepsilon^T P^{\dagger} \left[\left(D_{DWF(N_5, m_f=1)} \right)^{-1} D_{DWF(N_5, m_f)} \right] P \varepsilon$$

 $P\varepsilon = (P_L, 0, 0, \dots, 0, P_R)^T$: restrictor from 5D to 4D

- Final step in the reweighing factor requires the exact overlap operator.
 - Optimal Domain wall operator via Zolotarev approximation coefficients T.-W. Chiu, PRL 90,071601(2003)
 - + Pulling up Low-modes (low-mode improvements)

P.Hernandez, K.Jansen, M.Luscher, [hep-lat/0007015] K.Jansen, K.Nagai, JHEP12(2003)038

=> Extension to the even/odd 4d-site precoditioning (New?).

- Even/odd site preconditioning with low-mode improvements $D_{TROV}w = \varepsilon^{T} P^{\dagger} \left[\left(D_{DWF(N_{5},m_{f}=1)} \right)^{-1} D_{DWF(N_{5},m_{f})} \right] P \varepsilon w$ $\left(D_{TROV} \right)^{-1} w = \varepsilon^{T} P^{\dagger} \left[\left(D_{DWF(N_{5},m_{f})} \right)^{-1} D_{DWF(N_{5},m_{f}=1)} \right] P \varepsilon w$
- need the inversion of D_{DWF}.

 $\widetilde{D}_{DWF(N_5,m_f)} = D_{DWF(N_5,m_f)} \left(B + CM_{W5} \right)^{-1} = K - \frac{1}{2} M_{hop}$ $K = 4 - m_{dw} + (1 - M_{W5})(B + CM_{W5})^{-1}$ M_{w_5} : hopping in 5th direction. (extra - flavor mass matrix) $B = \text{diag}(b_1, b_2, \dots, b_{N_s}), \quad C = \text{diag}(c_1, c_2, \dots, c_{N_s}),$ $(b_i, c_i: \text{improvement coefficients}, j:5D - \text{index})$ $H_{W} = \gamma_{5} (D_{W} - m_{dw}) (a_{5} (D_{W} - m_{dw}) / 2 + 1)^{-1}$ Shamier standard Domain wall: $b_i = a_5, c_i = 0$ Chiu Optimal (Shamier kernel) Domain wall: $b_j = \frac{w_j + a_5}{2}, c_j = \frac{w_j - a_5}{2}$ Borici Domain wall (Neuberger kernel): $b_i = a_5, c_i = a_5$ $H_{W} = \gamma_{5} D_{W}$ Chiu Optimal (Neuberger kernel) Domain wall: $b_i = w_i, c_i = w_i$ Gneral, Moebius Domain wall: b_i, c_i R.Brower, H.Neff, K.Orginos, [arXiv:1206.5214]; Nucl.Phys. ProcSuppl.153(2006)191

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• Even/odd site preconditioning with low-mode improvements

Low-modes of sign kernel Hw

$$H_W V_k = V_k \Lambda_k, \quad V_k = (v_1, v_2, \dots, v_k),$$
$$\Lambda_k = \operatorname{diag}(\lambda_1, \lambda_2, \dots, \lambda_k)$$

Low-modes improvement for sign kernel Hw

$$H_W \to H_W^{imp} = H_W + V_k S_k V_k^{\dagger}$$

$$\widetilde{D}_{DWF(N_5,m_f)} = K - \frac{1}{2} M_{hop} \Longrightarrow \widetilde{D}_{DWF(N_5,m_f)}^{imp} = K - \frac{1}{2} M_{hop} + W_k S_k^{\dagger}$$
Low-modes improved Domain

Zolotarev coefficients in
$$B$$
 and C and N_5 are optimized with respect to $|Hw|$

$$W_{k} = ((a_{5} / 2)(D_{W} - m_{dw}) + 1)\gamma_{5}V_{k}$$

$$S_{k} = \operatorname{diag}(\hat{\lambda}_{1} - \lambda_{1}, \hat{\lambda}_{2} - \lambda_{2}, \cdots, \hat{\lambda}_{k} - \lambda_{k})$$

$$\hat{\lambda}_{j} \equiv 2\operatorname{sign}(\lambda_{j}) \max_{s=1,\cdots,k} |\lambda_{s}|$$

$$D_{DWF(N_5 \approx \infty, m_f)} x = b \Leftrightarrow \widetilde{D}_{DWF(N_5, m_f)}^{imp} \widetilde{x} = b, x = (B + CM_{W5})^{-1} \widetilde{x}$$

wal

• Even/odd site preconditioning with low-mode improvements

$$\widetilde{D}_{DWF(N_5,m_f)}^{imp}\widetilde{x} = b \qquad \widetilde{D}_{DWF(N_5,m_f)}^{imp} = \widetilde{D}_{DWF(N_5,m_f)} + W_k S_k^{\dagger}$$

Even/odd (4D) ordered form of the linear equation with improved DWF operator.

$$\begin{pmatrix} \left(\widetilde{D}_{DWF(N_{5},m_{f})} \right)_{ee} & \left(\widetilde{D}_{DWF(N_{5},m_{f})} \right)_{eo} & \left(W_{k} \right)_{e} \\ \left(\widetilde{D}_{DWF(N_{5},m_{f})} \right)_{oe} & \left(\widetilde{D}_{DWF(N_{5},m_{f})} \right)_{oo} & \left(W_{k} \right)_{o} \\ \left(S_{k}^{\dagger} \right)_{e} & \left(S_{k}^{\dagger} \right)_{o} & -1 \end{pmatrix} \begin{pmatrix} \widetilde{x}_{e} \\ \widetilde{x}_{o} \\ \zeta \end{pmatrix} = \begin{pmatrix} b_{e} \\ b_{o} \\ 0 \end{pmatrix}$$
 climinating x_{o} and ζ

Even/odd (4D) preconditioned form of the linear equation with improved DWF operator.

$$\left(\hat{\widetilde{D}}_{DWF(N_5,m_f)}^{imp}\right)_{ee}\tilde{x}_e = \hat{b}_e$$

$$\begin{aligned} G_{oo} &= \left(\mathbf{1} + S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} W_{o} \right)^{-1} \\ &\left(\hat{\widetilde{D}}_{DWF} \right)_{ee} = \left(\hat{\widetilde{D}}_{DWF} \right)_{ee}^{-1} \left\{ W_{e} - \left(\widetilde{D}_{DWF} \right)_{eo} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} W_{o} \right\} G_{oo} \left[S_{e}^{\dagger} + S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} \left(\widetilde{D}_{DWF} \right)_{oe} \right] \\ &\hat{b}_{e} &= \left(\widetilde{D}_{DWF} \right)_{ee}^{-1} \left[b_{e} - \left(\widetilde{D}_{DWF} \right)_{eo} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ &- \left(W_{e} - \left(\widetilde{D}_{DWF} \right)_{eo} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} W_{o} \right) G_{oo} S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ & \tilde{x}_{o} &= \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} \left[b_{o} - \left(\widetilde{D}_{DWF} \right)_{oe} \widetilde{x}_{e} \\ &- \left(W_{e} - \left(\widetilde{D}_{DWF} \right)_{eo} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} W_{o} \right) G_{oo} S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ & \tilde{x}_{o} &= \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} \left(\widetilde{D}_{DWF} \right)_{oo} \left(\widetilde{D}_{DWF} \right)_{oe} W_{o} \right) \widetilde{x}_{e} + S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} \left(\widetilde{D}_{DWF} \right)_{oe} W_{o} \right] \\ & \tilde{x}_{o} &= \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} \left(\widetilde{D}_{DWF} \right)_{oo} \left(\widetilde{D}_{DWF} \right)_{oe} W_{o} \right) \widetilde{x}_{e} + S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} \left(\widetilde{D}_{DWF} \right)_{oe} W_{o} \right) \widetilde{x}_{e} + S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo}^{-1} b_{o} \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} \left(\widetilde{D}_{DWF} \right)_{oe} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{oo} W_{o} \right] \\ & \left[S_{e}^{\dagger} - S_{o}^{\dagger} \left(\widetilde{D}_{DWF} \right)_{o} W_{o} \right$$

- Very preliminary
 - Test on Quench configurations
 (10 configs, DBW2 gauge action beta=0.87, 8³x32)

[explored by Y.Aoki et al.(RBC), PRD69(2004)074504]

[DBW2: Takaishi, PRD54(1996)1050 de Forcrand et al.(QCD-TARO), NPB577(2000)263]

- monitoring
 - 12 interior/exterior eigen values of $H_W = \gamma_5 (D_W m_{dw}) / ((D_W m_{dw}) / 2 + 1)$
 - finite mas G.-W. relation violation:

 $\frac{\eta^{\dagger} \Delta_{_{GW}} \eta}{\eta^{\dagger} \eta} \qquad \eta: \text{ Gaussian noise ve}$

 $\frac{1}{\eta^{\dagger}\eta} \qquad \eta: \text{ Gaussian noise vector}$

$$\Delta_{GW} = \left\{ \gamma_5, \left(D_{OV/TROV} \right)^{-1} \right\} - \frac{2}{1 + m_f} \left[m_f \left(D_{OV/TROV} \right)^{-1} \gamma_5 \left(D_{OV/TROV} \right)^{-1} + \gamma_5 \right]$$

- exponent in reweighting factor: $dS \equiv \left| \left(D_{TROV(j-1)} \right)^{-1} D_{TROV(j)} \eta \right|^2 \left| \eta \right|^2$
- reweighting factor : exp(-dS) η : Gaussian noise vector

No meson spectrum measurement with reweighing....

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• meson masses N₅=12



- M_{PS}= 0.274(21), M_V=0.611(65)
- M_{PS}/M_V=0.45

RBC,PRD69(2004), 16^3x32 , N₅=16,Domain wall M_{PS}= 0.2997(19), M_V=0.640(17)



Small |Hw| in conf # 1,2,3,9 10 \Rightarrow large G.-W. relation violation is expected with a fixed N5.



Small |Hw| in conf # 1,2,3, 9, 10 \Rightarrow large G.-W. relation violation.

Optimal operator keeps the violation at the tolerance level as it should be.

We used:

m=0.02

optimal Zolotarev tolerance 10⁻¹² eigen solver tolerance 10⁻¹²

Exponent dS

N5=30->32 and 32->Optimal

• Exponent dS

N5=12->14 and 14->Optimal



Reweighing factor exp(-dS)
 N5=12->14 and 14 -> optimal

Reweighing factor exp(-dS)
 N5=30->32 and 32 -> optimal



Reweighting between N5=32 and optimal/ N5=14 and optimal: small eigen values in Conf #1,2,9,10.

=> large fluctuation in exp(-dS).

The factor exp(-dS) seems to be sensitive to the appearance of small eigen values.

5. Summary

- Needs methods for Lattice chiral symmetry with lower computational cost
- One possibility: Truncated/approximate overlap Dirac operator with reweighting method
- We tested:
 - Truncated overlap fermions on top of the Domain wall fermion.
 - the G.-W. relation violation and reweighting factor are investigated on quenched configurations.
 - the violation and the reweighting factor are sensitive to the small eigen values of |*Hw*|
- Further studies:
- *dS* must proportional to the size of lattice. How about for larger lattice?
- How such configurations with small eigen values affect the expectation value of observables?
- Does the HMC with the truncated Overlap fermion generate such configurations?

Backups

G.-W. relation violation



• Quench, m=0.01

2013/7/29

Backups

• Quench, m=0.01

• Exponent dS



• Reweighing factor exp(-*dS*)