

Can a light Higgs
impostor
hide in composite
gauge models?

Chik Him Wong

Can a light Higgs impostor hide in composite gauge models?

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Fitting Strategies

Preliminary Results

Conclusion

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Lattice Higgs Collaboration (LHC):

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LATTICE 2013

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- Introduction: Sextet model as Composite Higgs candidate
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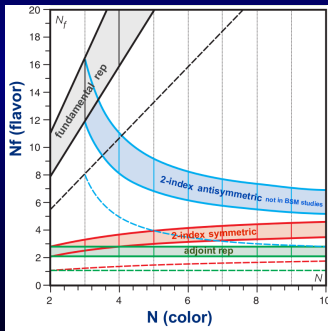
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- Goal: Look for a Composite Higgs model:
An infrared fixed point almost exists + Confining below Electroweak scale \Rightarrow models at the edge of conformal window
- After Higgs boson discovery : Light 0^{++} Higgs + reproduce detected phenomenology
- Parameter Space: N_C, N_f , Representations of $SU(N_C)$



- Focus of this talk: $SU(3) N_f = 2$ Sextet (Two-index symmetric) Model

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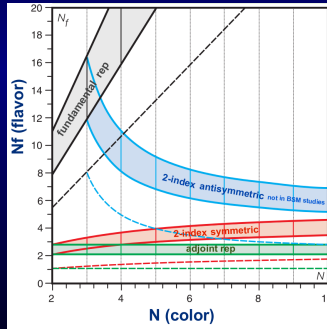
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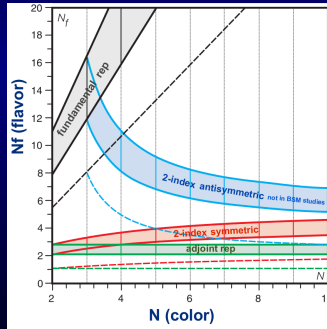
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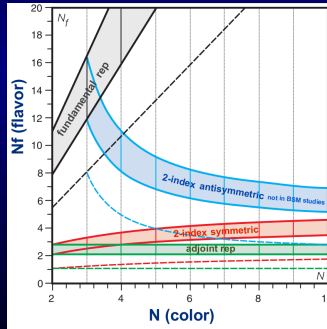
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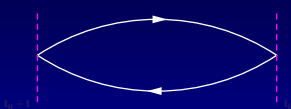
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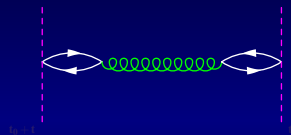
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- 0^{++} : Most computationally demanding and trickiest channel in spectroscopy, since

- For fermionic operators (f_0), two diagrams are involved:



Connected Diagram



Annihilation Diagram

- Annihilation diagram requires Same-time Quark Propagator
⇒ Cost of Exact Inversion is prohibitive → Stochastic calculation
- For gluonic operators (G , 0^{++} glueball), they are typically very noisy. Near Conformal Window, they can be light and coupled to the ground state
⇒ a very long trajectory is needed
- The above, possibly together with multi-hadron operators, are expected to mix in the ground state
⇒ Correlator Matrix may be needed

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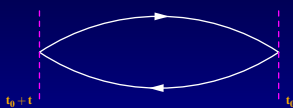
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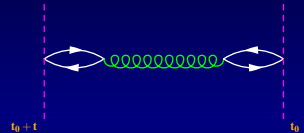
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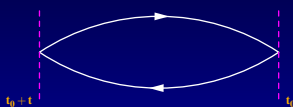
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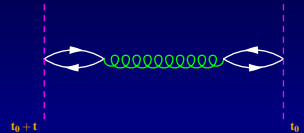
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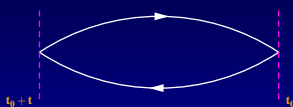
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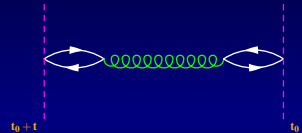
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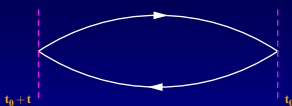
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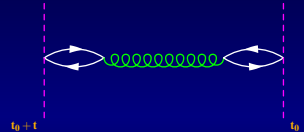
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- Staggered formalism:

$\langle \text{Tr} M^{-1}(t, t) \rangle = m \langle \text{Tr} (M^{-1}(t, t') [M^{-1}(t, t')]^\dagger) \rangle$, but they fluctuate differently:

- $\Delta[\text{Tr} M^{-1}(t, t)] \sim \langle \bar{\psi} \psi \rangle / m$, $m \Delta[\text{Tr} (M^{-1}(t, t') [M^{-1}(t, t')]^\dagger)] \sim \langle \bar{\psi} \psi \rangle^2$
- $\langle \bar{\psi} \psi \rangle < 1 \Rightarrow m \langle \text{Tr} (M^{-1}(t, t') [M^{-1}(t, t')]^\dagger) \rangle$ more preferred

- Same-time Quark Propagators needed \Rightarrow Stochastic Method

- Basic Idea: $M^{-1} \approx M^{-1}(\eta \eta^\dagger)_\eta = \langle \phi \eta^\dagger \rangle_\eta$, η : $Z(2)$ random noise
- “Dilution” employed: η projected to individual colors, timeslices and Even/Odd spatial partitions: $\eta_{[E]}(t)$ and $\eta_{[O]}(t)$

- $\phi_{[E/O]}(t, t_0) = \phi_{[E/O]}(t_0 + t, \eta_{[E/O]}(t_0))$

- Connected Diagram:

$$C(t) = -(-1)^t \text{Tr}(\phi_{[E]}(t, t_0) \phi_{[E]}(t, t_0)^\dagger - \phi_{[O]}(t, t_0) \phi_{[O]}(t, t_0)^\dagger) U, \eta, t_0$$

- Annihilation Diagram:

$$D(t) = \frac{N_c}{4} \langle \text{Tr}[\phi_{[E]}(0, t_0 + t) \phi_{[E]}(0, t_0 + t)^\dagger + \phi_{[O]}(0, t_0 + t) \phi_{[O]}(0, t_0 + t)^\dagger] \text{Tr}[\phi_{[E]}(0, t_0) \phi_{[E]}(0, t_0)^\dagger + \phi_{[O]}(0, t_0) \phi_{[O]}(0, t_0)^\dagger] \rangle U, \eta, t_0$$

- In case of finite momenta,

$$\phi_{[E/O]}(t, t_0) \rightarrow e^{-i\vec{y} \cdot \vec{p}} \phi_{[E/O]}(t_0 + t, e^{i\vec{x} \cdot \vec{p}} \eta_{[E/O]}(t_0))$$

$(\phi_{[E/O]}(t, t_0)^\dagger)$ unchanged

- 1 set of noise vectors per gauge configuration

- different implementations and dilution schemes are possible

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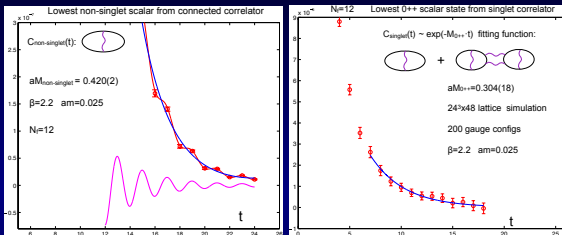
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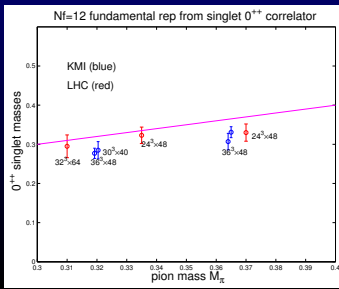
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• Test on $N_f = 12$ Fundamental SU(3) Model

- Known to be also close to Conformal Window
- Runs faster and more statistics available



• Comparison with KMI result [LHC: Fodor et al, KMI: Aoki et al (more details in Enrico Rinaldi's talk)]



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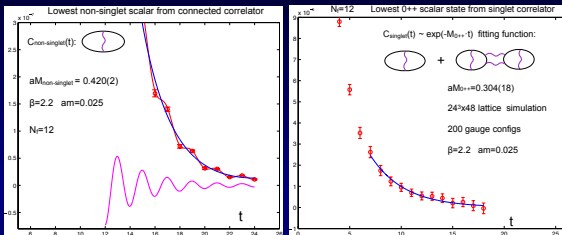
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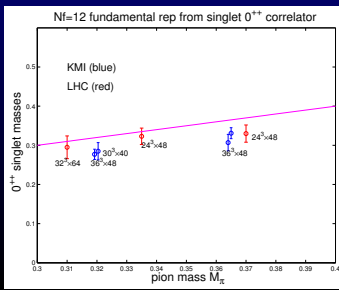
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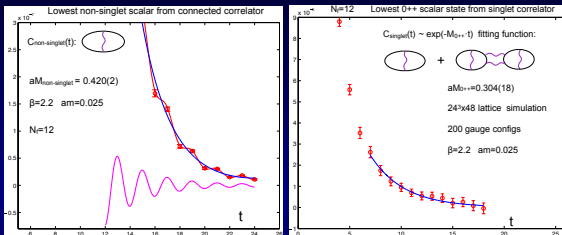
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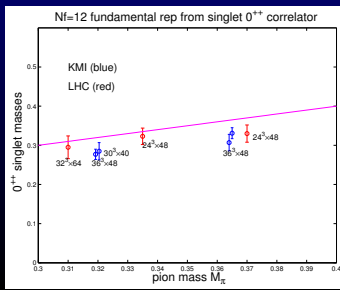
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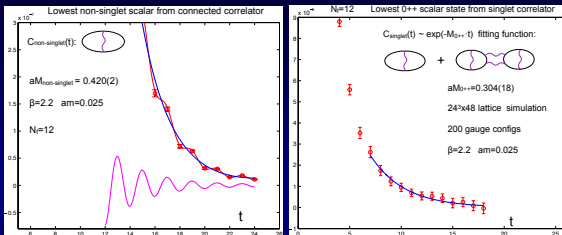
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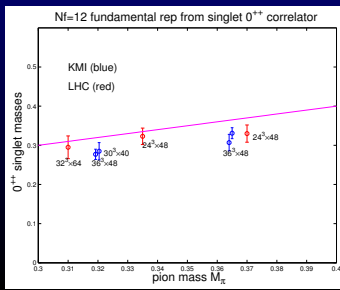
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- **Action: Tree-level Symanzik-Improved gauge action with Staggered $N_f = 2$ Sextet SU(3) fermions**
- RHC algorithm with multiple time scales and Omelyan integrator
- Autocorrelations monitored by time histories of effective masses and correlators
- $\beta \equiv 6/g^2 = 3.20$ and 3.25 , which is in the weak coupling regime
- **Lattices available: ($\sim 2000 - 4000$ Trajectories each)**

β	L	T	m_q
3.20	48	96	0.003
	32	64	0.003, 0.004, 0.005, 0.006, 0.007, 0.008
	28	56	0.003, 0.004, 0.005, 0.006, 0.007, 0.008
	24	48	0.003, 0.004, 0.005, 0.006, 0.007, 0.008, 0.009, 0.010, 0.012, 0.014
3.25	32	64	0.004, 0.005, 0.006, 0.007, 0.008
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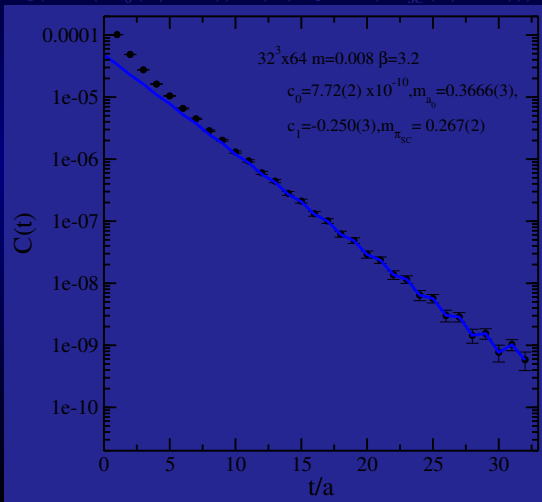
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- Observations of typical data:

- $C(t)$, also correlator of a_0 , is quiet and can be fitted well with the following ansatz:

$$C(t) = c_0(\cosh(m_{a_0}(T/2 - t)) + (-)^t c_1 \cosh(m_{\pi_{\text{SC}}}(T/2 - t)))$$



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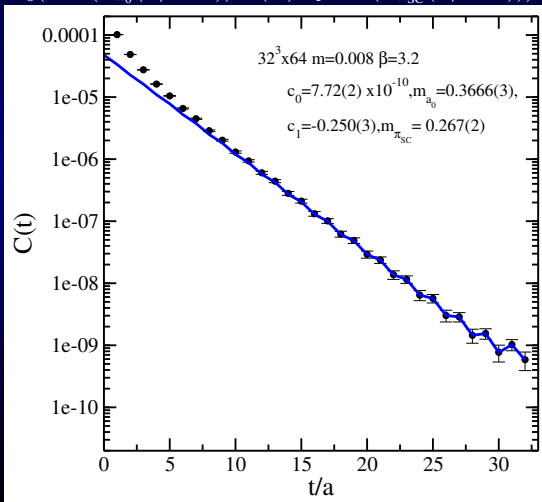
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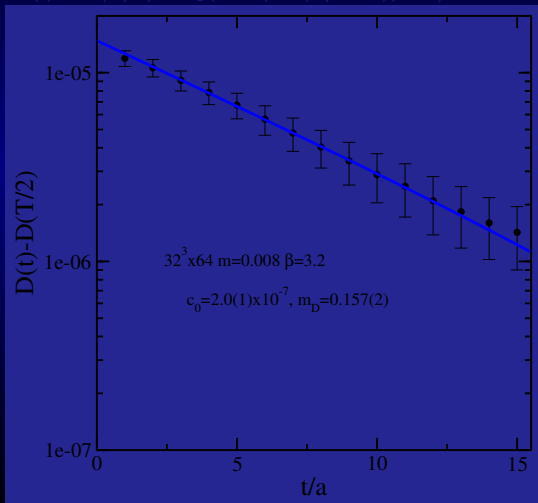
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- Observations of typical data:

- Difference between $D(t)$ and $D(T/2)$, $\tilde{D}(t)$ behaves exponential without detectable oscillation, with smaller exponent than $C(t)$
 $\tilde{D}(t) \equiv D(t) - D(T/2) = c_0(\cosh(m_D(T/2 - t)) - 1)$



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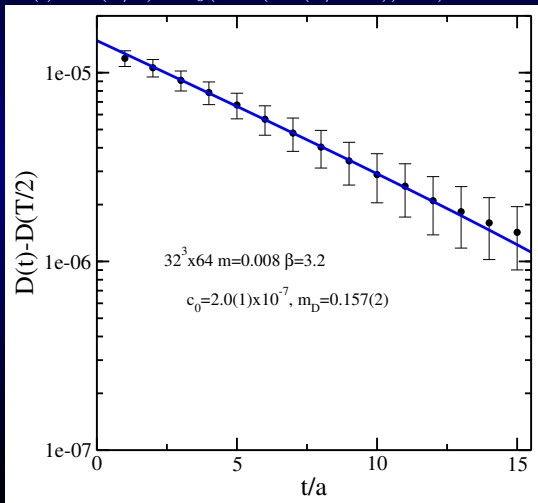
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- Observations of typical data:

- \Rightarrow Full subtracted correlator can be fitted well with the ansatz:

$$\begin{aligned}\tilde{D}(t) + C(t) = & c_0(\cosh(m_{f_0}(T/2 - t)) - 1) \\ & + c_1(\cosh(m_1(T/2 - t)) + (-)^t c_2 \cosh(m_{\eta_{SC}}(T/2 - t))),\end{aligned}$$

where $m_{f_0} \approx m_D$, $m_1 \approx m_{a_0}$ and $m_{\eta_{SC}} \approx m_{\pi_{SC}}$

- \Rightarrow Fitting $\tilde{D}(t)$ alone gives f_0 mass

- Effective mass definition:

$$\begin{aligned}& \frac{\tilde{D}(t) + 2\tilde{D}(t+1) + \tilde{D}(t+2)}{\tilde{D}(t-1) + 2\tilde{D}(t) + \tilde{D}(t+1)} \\ & \equiv [\cosh(m_{\text{eff}}(T/2 - t)) + 2\cosh(m_{\text{eff}}(T/2 - (t+1))) \\ & \quad + \cosh(m_{\text{eff}}(T/2 - (t+2))) - 4] \\ & \quad / [\cosh(m_{\text{eff}}(T/2 - (t-1))) + 2\cosh(m_{\text{eff}}(T/2 - t)) \\ & \quad + \cosh(m_{\text{eff}}(T/2 - (t+1))) - 4]\end{aligned}$$

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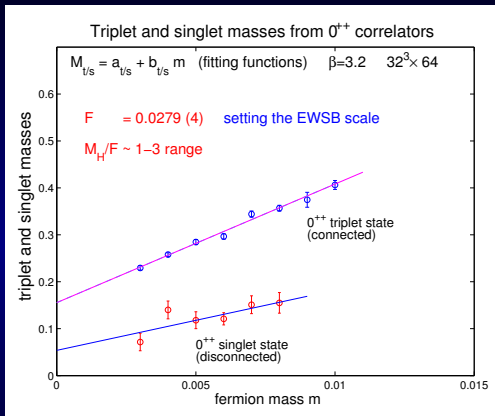
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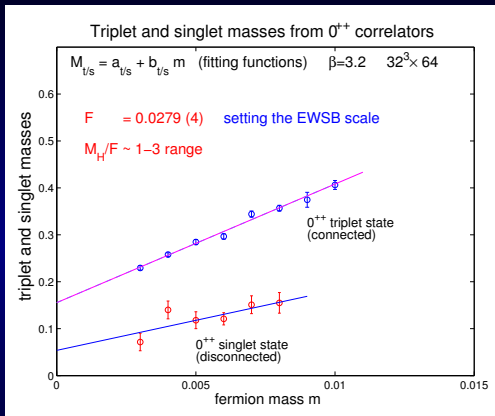
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 $[E_{f_0}(\vec{p} = (0, 0, 1))]^2 - m_{f_0}^2 - 4 \sin^2(\pi/L) = 0.019(17)$
- m_{f_0} can be as light as $250 - 750 \text{ GeV}$
- Radiative corrections due to top quarks can turn it into a Higgs impostor (R. Foadi, M Frandsen, F Sannino hep-ph: 1211.1083)

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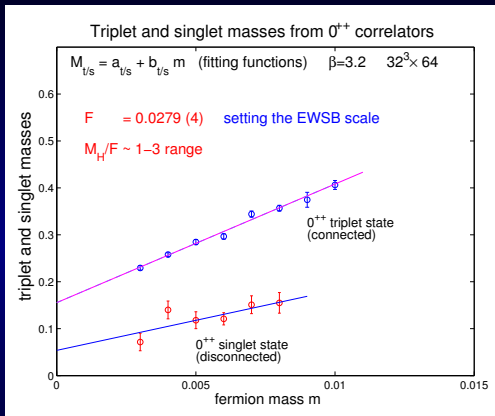
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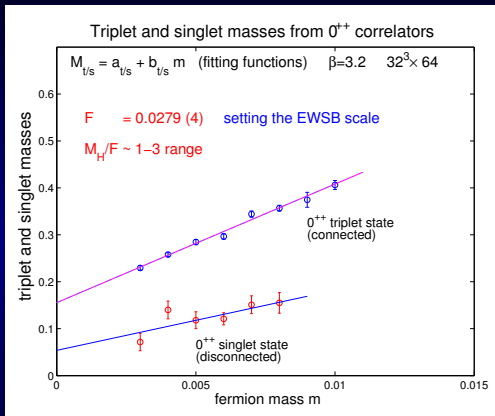
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- Future Plans
 - Investigate Finite Volume Effects on larger lattices
 - Investigate glueball and multi-hadron contributions
 - Investigate possible relation with dilatons
 - Investigate behavior in other weaker couplings
 - Improve efficiency by optimizing the choice of dilution schemes
 - Compare behavior with other models (e.g. $N_f = 8$ Fundamental $SU(3)$)

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