

# Moments of structure functions for $N_f = 2$ near the physical point.

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# Outline

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- ▶  $\langle x \rangle_{u-d}$
- ▶  $g_T$  and  $g_S$
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# Introduction

Nucleon structure investigated using matrix elements of the form  
 $\langle N, p' | \bar{q} \Gamma q | N, p \rangle$ :  $q = p' - p$ , (generalised) form factors → A. Sternbeck.

$p' = p = 0$ : lowest moments of pol. and unpol. structure functions.

$$2 \int_0^1 dx x F_1(x, Q^2) = c_{11}^{(u)} \langle x \rangle_u(\mu) + c_{11}^{(d)} \langle x \rangle_d(\mu) + c_{11}^{(s)} \langle x \rangle_s(\mu)$$

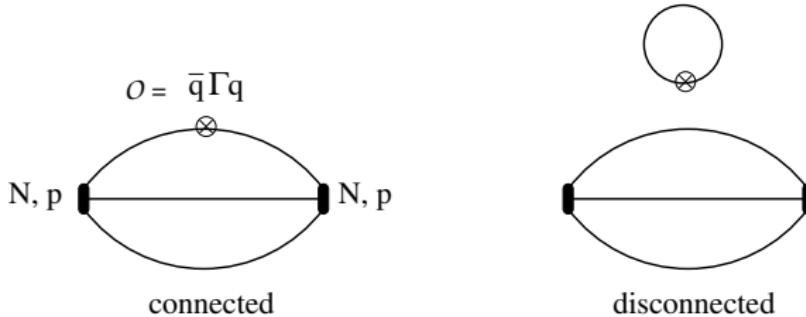
$$\langle x \rangle_u(\mu) = \sum_{q=u,d,s} Z_q(\mu) \langle x \rangle_q^{latt} \quad \langle x \rangle_q^{latt} \propto \langle N, p | \bar{q} [\gamma_4 \overset{\leftrightarrow}{D}_4 - \frac{1}{3} \gamma \cdot \overset{\leftrightarrow}{\mathbf{D}}] q | N, p \rangle^{latt}$$

$$2 \int_0^1 dx \mathbb{1} g_1(x, Q^2) = \frac{1}{2} [e_{10}^{(u)} \langle \mathbb{1} \rangle_{\Delta u}(\mu) + e_{10}^{(d)} \langle \mathbb{1} \rangle_{\Delta d}(\mu) + e_{10}^{(s)} \langle \mathbb{1} \rangle_{\Delta s}(\mu)]$$

$$\langle \mathbb{1} \rangle_{\Delta u}(\mu) = \sum_{q=u,d,s} Z'_q(\mu) \langle \mathbb{1} \rangle_{\Delta q}^{latt}, \quad \langle \mathbb{1} \rangle_{\Delta q}^{latt} \propto \langle N, p, s | \bar{q} \gamma_i \gamma_5 q | N, p, s \rangle^{latt}$$

Also non-Standard Model couplings  $\bar{q} q \rightarrow g_S$  and  $\bar{q} \sigma_{\mu\nu} q \rightarrow g_T$ .

Evaluate  $\langle N, p | \bar{q} \Gamma q | N, p \rangle$



$q = u/d$ : both connected and disconnected terms.

$q = s$ : only the disconnected term.

Mixing under renormalisation possible.

Disconnected results for  $\Gamma = 1$ , [talk by A. Sternbeck](#).

Here use isovector combinations  $\langle N, p | \bar{u} \Gamma u - \bar{d} \Gamma d | N, p \rangle \rightarrow g_A = \Delta u - \Delta d$ ,  $\langle x \rangle_{u-d}$ ,  $g_S^{u-d}$ ,  $g_T^{u-d}$ .

## Simulation details

$N_f = 2$ , NP-improved clover fermions, Regensburg+QCDSF configs.

$\beta$	Vol	#	$a$ fm	$m_\pi$ MeV	$Lm_\pi$	$t_{sink}$
5.20	$32 \times 64$	$823 \times 2$	0.08	280	3.7	13
5.29	$24 \times 48$	$1124 \times 2$	0.07	430	3.7	15
	$32 \times 64$	$2027 \times 2$ (1)		294	3.4	7,9,11,13,15,17
	$40 \times 64$	$2028 \times 2$		289	4.2	15
	$48 \times 64$	$3400 \times 2$		157	2.7	15
5.40	$64 \times 64$	$940 \times 3$		150	3.5	15
	$32 \times 64$	$1170 \times 2$	0.06	491	4.8	17
	$48 \times 64$	$2178 \times 2$		260	3.8	17

This analysis using sequential propagators: [using stochastic estimates see talk J. Najjar on Thursday](#).

Gauge invariant Wuppertal smearing - optimised for two-point function.

**Systematics:** in general - finite  $a$ , finite  $V$ , non-physical  $m_q$ , excited states.

## Excited state analysis

$$\begin{aligned} C_{2pt}(t_{sink}) &= \sum_{\vec{x}} T^{\alpha\beta} \langle \mathcal{N}_\alpha(\vec{x}, t_{sink}) \overline{\mathcal{N}}_\beta(\vec{0}, 0) \rangle = |Z_0|^2 e^{-m_0 t_{sink}} + |Z_1|^2 e^{-m_1 t_{sink}} + \dots \\ &= A_0 e^{-m_0 t_{sink}} [1 + e^{-\Delta m t_{sink}} + \dots] \end{aligned}$$

$$\begin{aligned} C_{3pt}(t_{sink}, t_{ins}) &= \sum_{\vec{x}, \vec{y}} T^{\alpha\beta} \langle \mathcal{N}_\alpha(\vec{x}, t_{sink}) O(\vec{y}, t_{ins}) \overline{\mathcal{N}}_\beta(\vec{0}, 0) \rangle \\ &= |Z_0|^2 \langle N_0 | O | N_0 \rangle e^{-m_0 t_{sink}} + Z_1 Z_0^* \langle N_1 | O | N_0 \rangle e^{-m_0 t_{ins}} e^{-m_1 (t_{sink} - t_{ins})} \\ &\quad + Z_0 Z_1^* \langle N_0 | O | N_1 \rangle e^{-m_1 t_{ins}} e^{-m_0 (t_{sink} - t_{ins})} + |Z_1|^2 \langle N_1 | O | N_1 \rangle e^{-m_1 t_{sink}} + \dots \\ &= A_0 e^{-m_0 t_{sink}} (B_0 + B_1 [e^{-\Delta m (t_{sink} - t_{ins})} + e^{-\Delta m t_{ins}}] + B_3 e^{-\Delta m t_{sink}}) + \dots \end{aligned}$$

where  $B_0 = \langle N_0 | O | N_0 \rangle$ ,  $B_1 \propto \langle N_1 | O | N_0 \rangle$ ,  $B_2 \propto \langle N_1 | O | N_1 \rangle$ ,

Fit  $C_{2pt}$  and  $C_{3pt}$  simultaneously for different  $t_{sink}$  and compare with constant fit to

$$\frac{C_{3pt}(t_{sink}, t_{ins})}{C_{2pt}(t_{sink})} = B_0 + \dots$$

A lot of work on excited state contamination recently. For example,

Dinter et al. (2011) arXiv:1108.1076 - multiple  $t_{sink}$ .

Owen et al. (2012) arXiv:1212.4668 - variational approach.

Capitani et al. (2012) arXiv:1205.0180 - summation method.

Green et al. (2012) arXiv:1209.1687 - summation method.

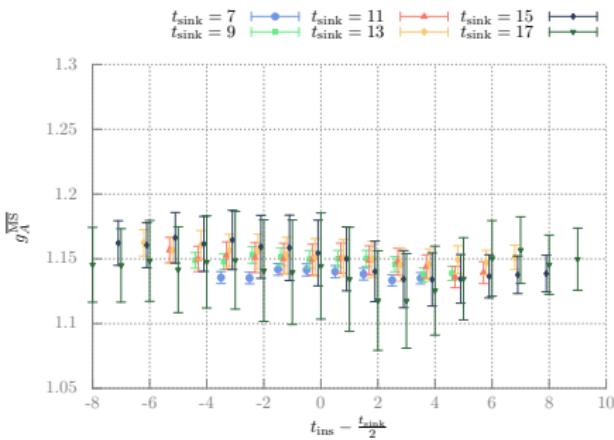
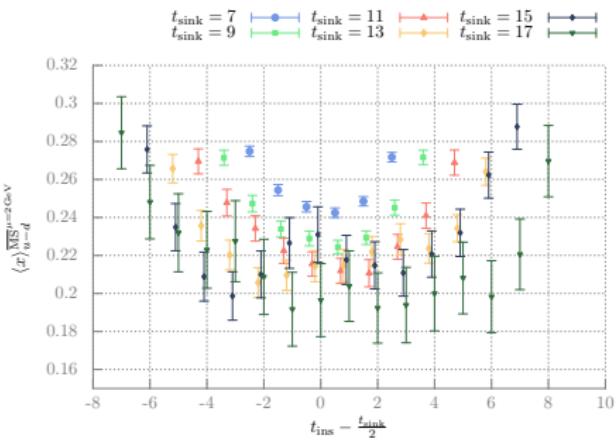
Bhattacharya et al. (2013) arXiv:1306.5435 - fits to multiple  $t_{sink}$ .

Non-perturbative renormalisation: Constantinou et al. arXiv:1303.6776,  
Göckeler et al. (QCDSF/UKQCD) arXiv:1003.5756.

All results preliminary! Work in progress.

$\beta = 5.29$ ,  $m_\pi = 290$  MeV,  $V = 32 \times 64$  lattice,  $t_{sink} = 7, 9, 11, 13, 15, 17$ .

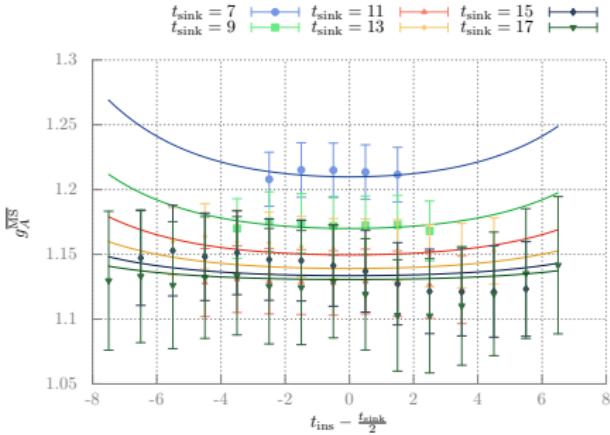
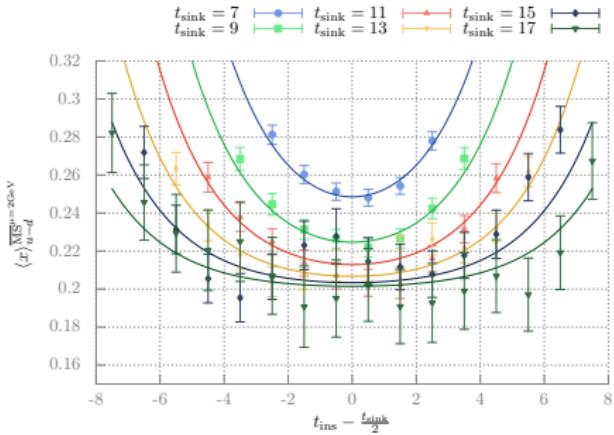
Example  $x_{u-d}$  and  $g_A$ :  $C_{3pt}(t_{ins}, t_{sink})/C_{2pt}(t_{sink}) + \text{additional factors}$ .



Fit compared to data:  $C_{3pt}(t_{ins}, t_{sink})/(A_0 e^{-m_0 t_{sink}})$  compared to

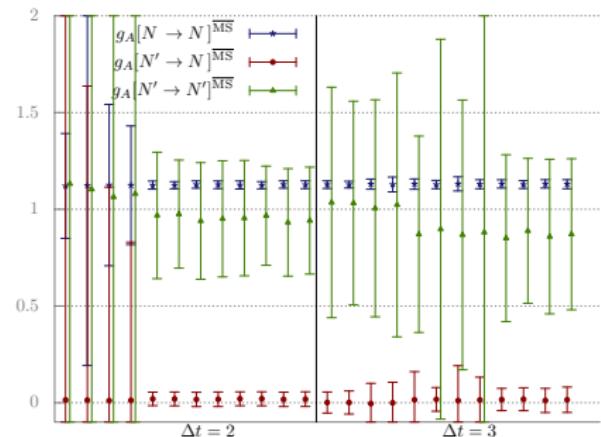
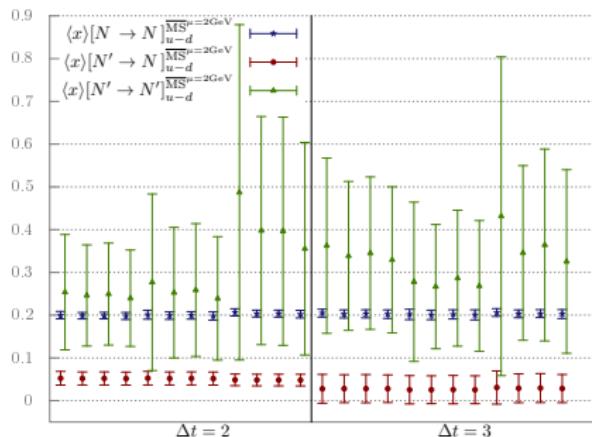
$$F = B_0 + B_1 [e^{-\Delta m(t_{sink} - t_{ins})} + e^{-\Delta m t_{ins}}] + B_3 e^{-\Delta m t_{sink}}$$

+ additional factors.



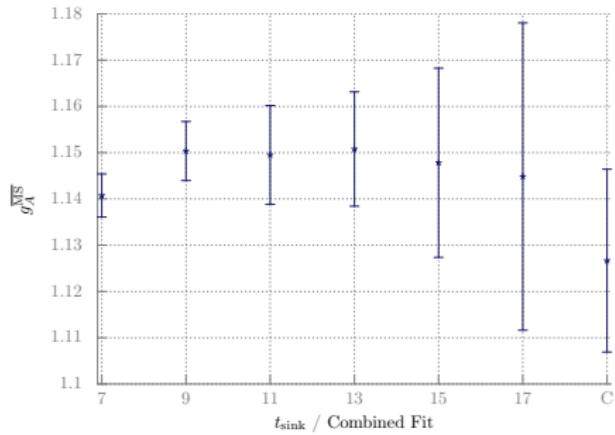
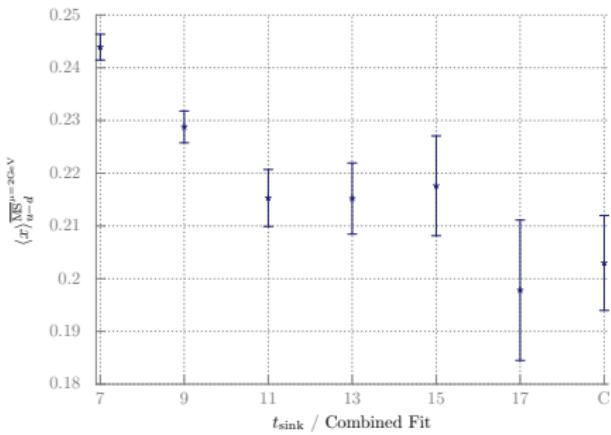
Excited state contribution to  $g_A$  cancels in  $C_{3pt}/C_{2pt}$ .

Variation in fitting range:  $\Delta t = t_{min} = t_{sink} - t_{max}$  for  $C_{3pt}$ .

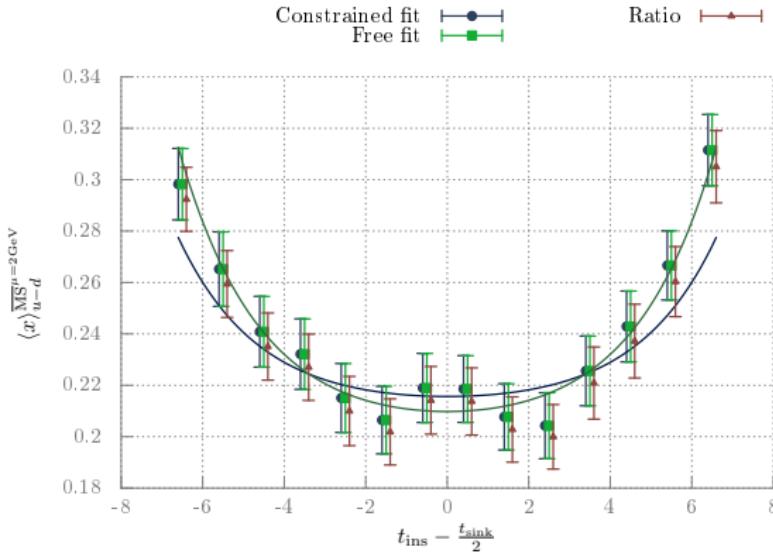


Variation for  $C_{2pt}$ ,  $t_{max} = 13, 17, 20, 23$  and  $t_{min} = 2, 3, 4$ .

Comparison with constant fits to  $C_{3pt}/C_{2pt}$  for different  $t_{sink}$ .



Apply to ensembles with single  $t_{sink}$ :  $\beta = 5.29$ ,  $m_\pi = 157$  MeV,  
 $V = 48^3 \times 64$ .

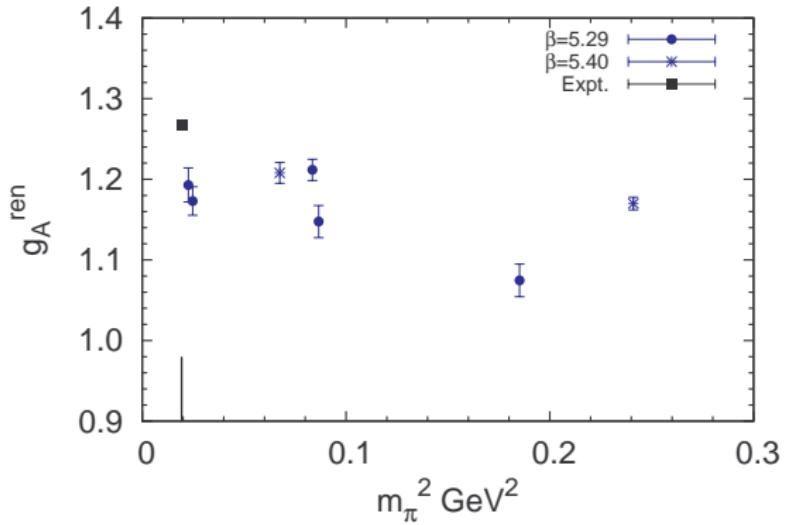


$$C_{3pt}(t_{sink}, t_{ins}) = A_0 e^{-m_0 t_{sink}} (B_0 + B_1 [e^{-\Delta m(t_{sink}-t_{ins})} + e^{-\Delta m t_{ins}}] + B_3 e^{-\Delta m t_{sink}})$$

Free fit:  $B_3 = 0$ .

Constrained fit: using excited state parameters from multi- $t_{sink}$  fits.  
 $\langle N_1 | O | N_0 \rangle$  and  $\langle N_1 | O | N_1 \rangle$

$g_A$

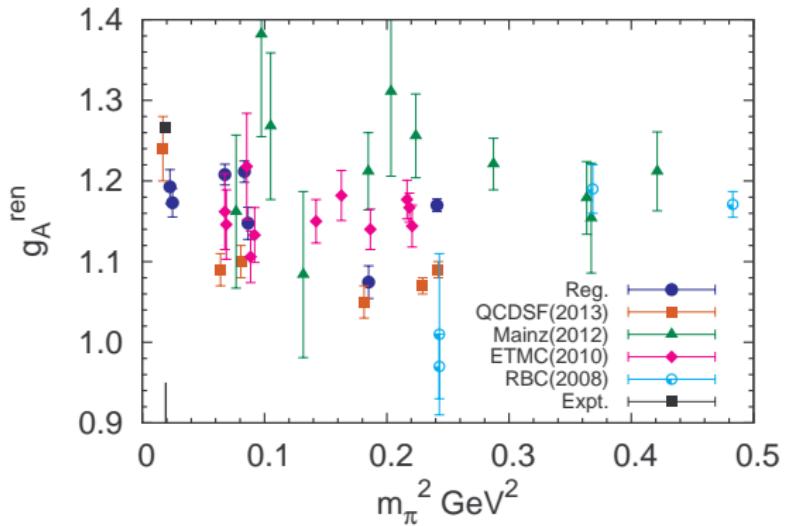


Results from a constant fit to  $C_{3pt}/C_{2pt}$ .

Finite volume effects:  $m_\pi \sim 290$ ,  $Lm_\pi = 3.4$  and 4.2  
 $m_\pi \sim 150$ ,  $Lm_\pi = 2.7$  and 3.5

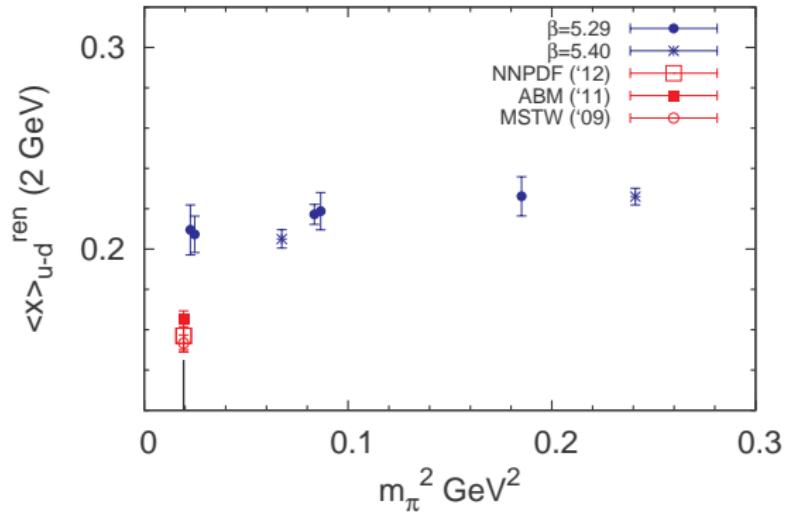
Finite  $a$  effects possible.

Comparison of  $N_f = 2$  results.



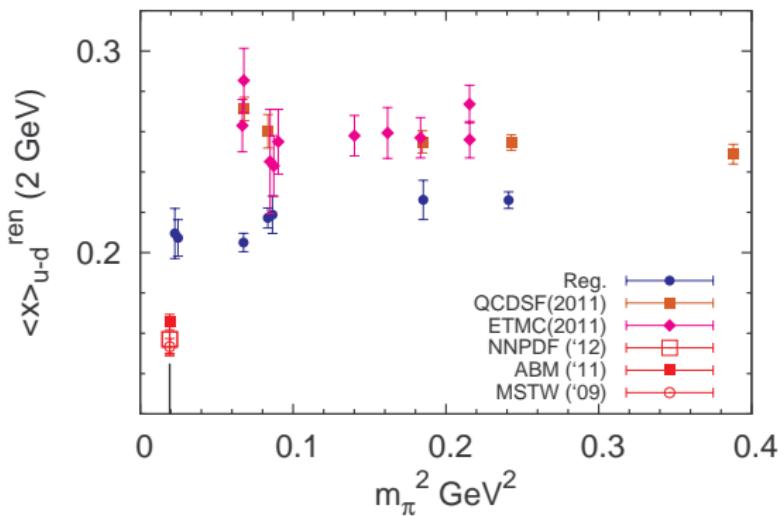
$N_f = 2 + 1$ : RBC/UKQCD (2010), QCDSF (2010), HSC (2011),  
CSSM (2012), LHPC (2012).  $N_f = 2 + 1 + 1$ : ETMC (2013), PNDME (2013).

$\langle x \rangle_{u-d}$



- ▶ No significant finite volume effects
- ▶ Possible  $a$  effect?

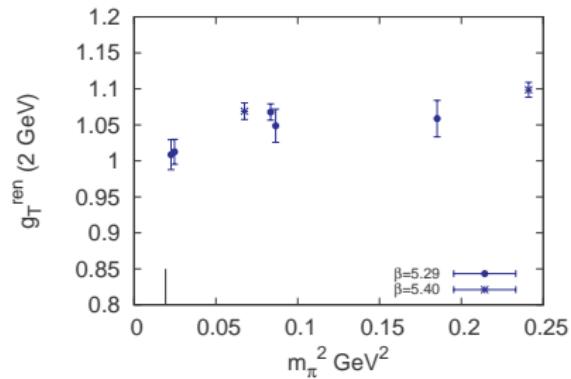
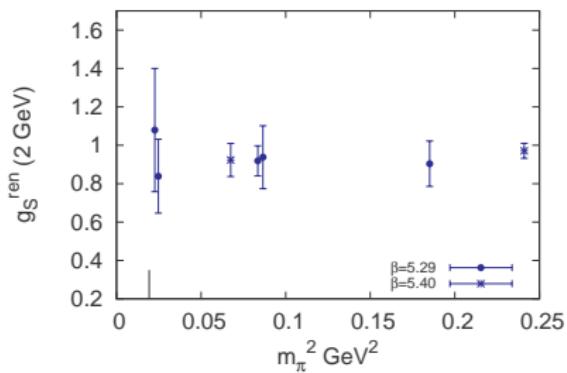
## Comparison of NP renormalised $N_f = 2$ results.



$N_f = 2 + 1$ : RBC/UKQCD (2010), LHPC (2010), LHPC (2012).

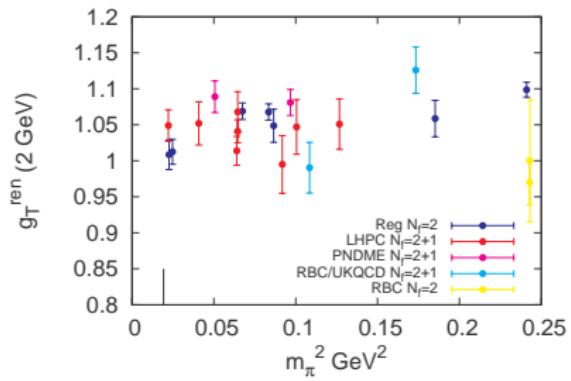
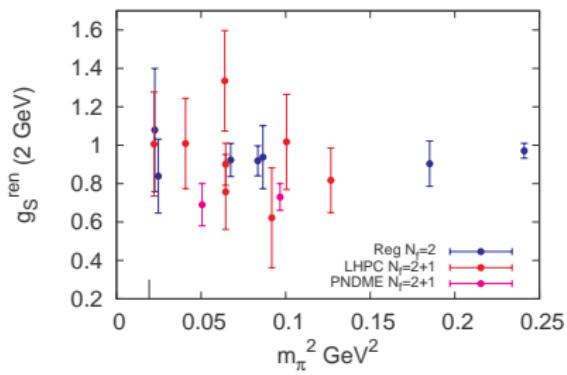
$N_f = 2 + 1 + 1$ : ETMC (2013)

$g_S^{u-d}$  and  $g_T^{u-d}$



- ▶ No significant dependence on  $\beta$ ,  $V$  or  $m_q$ .

# $g_S^{u-d}$ and $g_T^{u-d}$ : comparison with other groups



## Outlook

Work in progress to investigate the systematic uncertainties in calculating  $\langle x \rangle_{u-d}$  and  $g_A$ .

- ▶ Analysis needs to be finalised - increased statistics at  $m_\pi = 150$  MeV,  $Lm_\pi = 3.5$ .
- ▶ Excited states contamination can be brought under control.
- ▶ Volume effects for  $g_A$  - generating  $Lm_\pi = 6.7$  at  $m_\pi = 290$  MeV, finite  $V$  extrapolation.
- ▶  $g_A$  and  $\langle x \rangle_{u-d}$ , possible finite  $a$  effects - generating  $\beta = 5.20$  ensemble.

Results for  $g_S$  and  $g_T$  consistent with other groups.

Calculation of disconnected contributions in progress.