Neutron and proton EDM in $N_f = 2+1$ domain-wall fermion

Eigo Shintani (RIKEN-BNL)

T. Blum and T Izubuchi and A. Soni for RBC/UKQCD collaboration

Outline

• Motivation of lattice calculation of EDM in θ term

- Strategy and method
- (Preliminary) results
- Summary and future work

1. Introduction Motivation

- CP violation (CPV) in QCD and BSM
 - EDM is sensitive to CPV in BSM
 - ▶ SM contribution is extremely small, 10⁻³³ e⁻cm
 - Strong CP problem
 - There are many experimental plans of EDM pEDM experiment @ BNL, nEDM experiment @ ORNL, ILL, FRM-2, FNAL, PSI/KEK/TRIUMF, ... Charged particle (d, p)EDM @ COSY Lepton EDM @ J-PARC, FNAL aim to 10⁻²⁹ e[•] cm (present nEDM limit : |d_N^{exp}| < 2.9 × 10⁻²⁶ e · cm)
 - ▶ BSM (like SUSY) predicts ~10⁻²⁷ e·cm
 - Large uncertainty is hadron effect in low-energy scale.
 - Lattice QCD provides important information of BSM prediction.

1. Introduction Nucleon EDM in lattice QCD

- θ term contribution
 - Renormalizable
 - Feasible study toward the BSM (quark EDM and chromo EDM) calculation

Bhattacharya et al, Lattice 2012, Lattice 2013

- Nice check of uncertainties in quark model, BChPT
- Some difficulties
 - Statistical noise
 - Gauge background (topological charge, sea quark) which are intrinsically noisy.
 - Disconnected diagram (flavor singlet) should include (In SU(3) limit this is vanishing. Study including disc. diagram is future work).
 - Systematic study
 - Finite volume effect may be significant. (e.g. BChPT discussion)

O'Connell, Savage, PLB633, 319(2006), Guo, Meissner, 1210.5887

• Chiral behavior($d_N \sim O(m)$) is also important check.

\blacktriangleright Precise calculation of EDM in θ term

• Precise calculation of EDM in θ term

Use of domain-wall fermion (DWF)

Chiral symmetry on the lattice, and control the O(a) lattice artifact. Good control of chiral behavior without counterterm (cf.Wilson-clover)

- Precise calculation of EDM in θ term
 - Use of domain-wall fermion (DWF)

Chiral symmetry on the lattice, and control the O(a) lattice artifact. Good control of chiral behavior without counterterm (cf.Wilson-clover)

- Error reduction techniques
 - All-mode-averaging (AMA)
 - Efficient way to reduce the statistical error of correlator

Blum, Izubuchi, ES, arXiv:1208.4349 [hep-lat], ES (lattice 2012)

- Precise calculation of EDM in θ term
 - Use of domain-wall fermion (DWF)

Chiral symmetry on the lattice, and control the O(a) lattice artifact. Good control of chiral behavior without counterterm (cf.Wilson-clover)

- Error reduction techniques
 - All-mode-averaging (AMA)
 - Efficient way to reduce the statistical error of correlator

Blum, Izubuchi, ES, arXiv:1208.4349 [hep-lat], ES (lattice 2012)

- EDM form factor
 - Extraction from matrix element including CPV of θ angle
 - Extrapolation into physical kinematics, $-q^2 = 0$ and m = m_{phy}

2. EDM calculation on the lattice **Domain-wall fermion**

- Chiral symmetry
 - L, R fermion is localized on boundaries
 - \Rightarrow Chiral symmetry is realized (if $L_s \rightarrow \infty$).
 - Remaining good chiral symmetry m_{res} ~ exp(-L_s)
 - Reasonable computational cost
- RBC/UKQCD collaboration
 - Generation of N_f = 2+1 configurations.
 - Various lattice size, lattice cut-off, quark mass are available.
 - Many studies of Keon physics, nucleon physics, finite temperature , ...



[Blum Soni, (97), CP-PACS(99), RBC(00), RBC/UKQCD. (05--)]

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator O, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

2. EDM calculation on the lattice Blum ES (la ES (la Chulve)

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator O, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

O> = <O(imp)> if approximation is covariant under lattice symmetry g

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator *O*, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

O> = <O^(imp)> if approximation is covariant under lattice symmetry g

• Improved error $\operatorname{err}^{\operatorname{imp}} \simeq \operatorname{err}/\sqrt{N_G}$

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator *O*, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

O> = <O^(imp)> if approximation is covariant under lattice symmetry g

- Improved error $\operatorname{err}^{\operatorname{imp}} \simeq \operatorname{err}/\sqrt{N_G}$
- Computational cost of O^(imp) is cheap.

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator O, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

- O> = <O^(imp)> if approximation is covariant under lattice symmetry g
- Improved error $\operatorname{err}^{\operatorname{imp}} \simeq \operatorname{err}/\sqrt{N_G}$
- Computational cost of O^(imp) is cheap.
- All-mode-averaging (AMA)
 - Relaxed CG solution for approximation $\mathcal{O}^{(\text{appx})} = \mathcal{O}[S_l], S_l = \sum_{\lambda=1}^{N_{\lambda}} v_{\lambda} v_{\lambda}^{\dagger} \frac{1}{\lambda} + P_n(\lambda)|_{|\lambda| > N_{\lambda}}$



Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
 - For original correlator O, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

- O> = <O^(imp)> if approximation is covariant under lattice symmetry g
- Improved error $\operatorname{err}^{\operatorname{imp}} \simeq \operatorname{err}/\sqrt{N_G}$
- Computational cost of O^(imp) is cheap.
- All-mode-averaging (AMA)
 - Relaxed CG solution for approximation $\mathcal{O}^{(\text{appx})} = \mathcal{O}[S_l], S_l = \sum_{\lambda=1}^{N_{\lambda}} v_{\lambda} v_{\lambda}^{\dagger} \frac{1}{\lambda} + P_n(\lambda)|_{|\lambda| > N_{\lambda}}$
 - $P_n(\lambda)$ is polynomial approximation of I/λ
 - Low mode part :# of eigen mode
 - Mid-high mode : degree of poly.



ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\langle n(P_1)|J_{\mu}^{\text{EM}}|n(P_2)\rangle_{\theta} = \bar{u}_N^{\theta} \Big[\underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N}\gamma_5\sigma_{\mu\nu}Q_{\nu}}_{\text{P,T-odd}} + \underbrace{F_1\gamma_{\mu} + \frac{F_2}{2m_N}\sigma_{\mu\nu}Q_{\nu}}_{\text{P,T-even}} + \cdots \Big]u_N^{\theta}$$

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\begin{split} \langle n(P_1)|J_{\mu}^{\mathrm{EM}}|n(P_2)\rangle_{\theta} &= \bar{u}_{N}^{\theta} \Big[\underbrace{\frac{F_{3}^{\theta}(Q^2)}{2m_{N}}\gamma_{5}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-odd}} + \underbrace{F_{1}\gamma_{\mu} + \frac{F_{2}}{2m_{N}}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-even}} + \cdots \Big]u_{N}^{\theta} \\ &\underbrace{\sum_{s} u_{N}^{\theta}(s)\bar{u}_{N}^{\theta}(s) = \frac{ip\cdot\gamma + m_{N}e^{i\alpha_{N}^{\theta}\gamma_{5}}}{2E_{N}}}_{\mathrm{2}E_{N}} \Big] \\ \end{split}$$

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\begin{split} \langle n(P_1) | J_{\mu}^{\text{EM}} | n(P_2) \rangle_{\theta} &= \bar{u}_N^{\theta} \Big[\underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N} \gamma_5 \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-odd}} + \underbrace{F_1 \gamma_{\mu} + \frac{F_2}{2m_N} \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-even}} + \cdots \Big] u_N^{\theta} \\ &\sum_s u_N^{\theta}(s) \bar{u}_N^{\theta}(s) = \frac{ip \cdot \gamma + m_N e^{i\alpha_N^{\theta} \gamma_5}}{2E_N} \Big] \quad \text{CPV phase } \alpha_{\text{N}} \text{ in nucleon propagator} \\ \langle \theta | \eta_N J_{\mu}^{\text{EM}} \bar{\eta}_N | \theta \rangle &= \langle 0 | \eta_N J_{\mu}^{\text{EM}} \bar{\eta}_N | 0 \rangle + i\theta \langle 0 | \eta_N J_{\mu}^{\text{EM}} Q \bar{\eta}_N | 0 \rangle \end{split}$$

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\begin{split} \langle n(P_{1})|J_{\mu}^{\mathrm{EM}}|n(P_{2})\rangle_{\theta} &= \bar{u}_{N}^{\theta} \Big[\underbrace{\frac{F_{3}^{\theta}(Q^{2})}{2m_{N}}\gamma_{5}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-odd}} + \underbrace{F_{1}\gamma_{\mu} + \frac{F_{2}}{2m_{N}}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-even}} + \cdots \Big] u_{N}^{\theta} \\ &\sum_{s} u_{N}^{\theta}(s)\bar{u}_{N}^{\theta}(s) = \frac{ip\cdot\gamma + m_{N}e^{i\alpha_{N}^{\theta}\gamma_{5}}}{2E_{N}} \Big] \\ & \mathsf{CPV} \text{ phase } \alpha_{\mathsf{N}} \text{ in nucleon propagator} \\ \langle \theta|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|\theta\rangle &= \langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|0\rangle + i\theta\langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}Q\bar{\eta}_{N}|0\rangle \\ \langle 0|\eta_{N}(t_{1})J_{\mu}^{\mathrm{EM}}(t)Q\bar{\eta}_{N}(t_{0})|0\rangle & \\ &= \frac{\alpha_{N}}{2}\gamma_{5}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} + \frac{1+\gamma_{4}}{2}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{\alpha_{N}}{2}\gamma_{5} & \Big] \\ \text{ Subtraction} \\ &+ \frac{1+\gamma_{4}}{2}\Big[F_{3}\frac{q_{\nu}\gamma_{5}\sigma_{\mu\nu}}{2m_{N}} + F_{A}(iq^{2}\gamma_{\mu}\gamma_{5} - 2m_{N}q_{\mu}\gamma_{5})\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} & \Big] \\ \begin{array}{c} ip\cdot\gamma + m_{N} \\ ip\cdot\gamma + m_{N} \\ 2E_{N} \\ \end{array} \end{split}$$

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

EM Matrix element

$$\langle n(P_{1})|J_{\mu}^{\mathrm{EM}}|n(P_{2})\rangle_{\theta} = \bar{u}_{N}^{\theta} \left[\underbrace{\frac{F_{3}^{\theta}(Q^{2})}{2m_{N}}\gamma_{5}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-odd}} + \underbrace{F_{1}\gamma_{\mu} + \frac{F_{2}}{2m_{N}}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-even}} + \cdots \right] u_{N}^{\theta}$$

$$\sum_{s} u_{N}^{\theta}(s)\bar{u}_{N}^{\theta}(s) = \frac{ip\cdot\gamma + m_{N}e^{i\alpha_{N}^{\theta}\gamma_{5}}}{2E_{N}} \quad \left] \quad \text{CPV phase } \alpha_{N} \text{ in nucleon propagator}$$

$$\langle \theta|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|\theta\rangle = \langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|0\rangle + i\theta\langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}Q\bar{\eta}_{N}|0\rangle$$

$$\langle 0|\eta_{N}(t_{1})J_{\mu}^{\mathrm{EM}}(t)Q\bar{\eta}_{N}(t_{0})|0\rangle \qquad \qquad \\ \left\{ \frac{\alpha_{N}}{2}\gamma_{5}\left[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\right]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} + \frac{1+\gamma_{4}}{2}\left[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\right]\frac{\alpha_{N}}{2}\gamma_{5} \right] \quad \text{Subtraction}$$

$$+ \frac{1+\gamma_{4}}{2}\left[F_{3}\frac{q_{\nu}\gamma_{5}\sigma_{\mu\nu}}{2m_{N}} + F_{A}(iq^{2}\gamma_{\mu}\gamma_{5} - 2m_{N}q_{\mu}\gamma_{5})\right]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} \quad \left\{ \frac{ip\cdot\gamma + m_{N}}{2E_{N}} + \frac{1+\gamma_{4}}{2}\left[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\right]\frac{\alpha_{N}}{2}\gamma_{5} \right] \quad \text{Extraction}$$

- Subtraction of CP-odd phase, $\alpha_{\rm N}$, in n propagator and CP-even part ${\sf F}_{\rm I,2}$

$$d_N = \lim_{Q^2 \to 0} F_3(Q^2) / 2m_N$$

3. (Preliminary) Results **Parameters**

DWF

- > $24^3 \times 64$ lattice, $a^{-1} = 1.73$ GeV (~3 fm³ lattice)
- $L_{\rm s} = 16$ and $am_{\rm res} = 0.003$
- m = 0.005, 0.01 corresponding to $m_{\pi} = 0.33, 0.42$ GeV
- Two temporal separation of N sink and source in 3 pt. function

$$t_{sep} = 12 (t_{source} = 0, t_{sink} = 12), t_{sep} = 8 (t_{source} = 0, t_{sink} = 8)$$

configs = 751 (m=0.005), 700 (m=0.01) [t_{sep} = 12] # configs = 180 (m=0.005) [t_{sep} = 8]

AMA

- # of low-mode : $N_{\lambda} = 400 \text{ (m=0.005)}, 180 \text{ (m=0.01)}$
- Stopping condition, |r| < 0.003</p>
- ▶ $N_G = 32$ (2 separation for spatial, 4 separation for temporal direction of source localtion) → effectively O(10⁴) statistics

3. (Preliminary) Results

α_N

- Projection with γ_5 for 2 pt with Q charge, global fitting with $\operatorname{tr}\left[\gamma_5 \langle N(t)\bar{N}(0)Q \rangle\right] = Z_N \frac{2m_N}{E_N} \alpha_N e^{-E_N t} + O(e^{-E_N *})$
- By using AMA, this factor is determined within 15 % error.
- It does not depend on smearing and momentum, but mass dependence is not so clear.



3. (Preliminary) Results Subtraction term and 3pt function

- Splitting EDM form factor into two parts: $F_3 = F_Q + F_{\alpha}, F_Q = C(m_N) \langle N J_t^{\text{EM}} \bar{N} Q \rangle, F_{\alpha} = F(\alpha_N, F_{1,2})$
- F_{α} is good precision, and fluctuation of F_{O} is large.



3. (Preliminary) Results Comparison with $\mu = t, z$

- EDM form factor is given from two directions of EM current
- Two signals are consistent, and data in t direction is much stable.



3. (Preliminary) Results Comparison with different t_{sep}

- The sink and source separation in 3pt function enables us to control the statistical noise and excited state contamination
 - Short: statistical fluctuation < excited state contamination</p>
 - Long: statistical fluctuation > excited state contamination
- Comparison
 - $t_{sep} = 12 \text{ (blue),}$ [$N_{conf} = 751$] $t_{sep} = 8 \text{ (green)}$ [$N_{conf} = 180$]
 - Good consistency between them.
 - Precision in t_{sep}=8 is much better.



3. (Preliminary) Results $-q^2 = 0$ extrapolation

- Fitting data of EDM form factor at each momenta.
- Open(t_{sep} = 8),
 filled (t_{sep} = 12)
- Fitting function
 - 3 point linear :
 - $-q^2 < 0.55 \text{ GeV}^2$
 - 2 point linear: -q² < 0.4 GeV²
- Estimate of systematic error of extrapolation



3. (Preliminary) Results Mass dependence

Comparison with full QCD results

- DWF results are in the lightest quark mass.
- Statistical error is dominant rather than systematic one.
- Central value is 10 times larger than models.



3. (Preliminary) Results Statistical error

- Comparison between AMA error reduction and number of configurations.
- Number of configurations : reduce stat. error and relating to Q distribution
 AMA error reduction : reduce stat. error



4. Summary Summary and future plan

- Nucleon EDM in $N_f = 2+1$ DWF in θ vacuum
 - Signal of EDM within 40% statistical error using AMA techniques.
 - ▶ 3-pt function is still noisy.
 - Short t_{sep} allows us to reduce the statistical error without large excited state contamination effect.

4. Summary Summary and future plan

- Nucleon EDM in $N_f = 2+1$ DWF in θ vacuum
 - Signal of EDM within 40% statistical error using AMA techniques.
 - ▶ 3-pt function is still noisy.
 - Short t_{sep} allows us to reduce the statistical error without large excited state contamination effect.
- (Near) physical point of DWF configurations
 - Ensembles near physical points and large volume are available.
 - > AMA with Möbius-DWF approximation is helpful.
 - Remove chiral extrapolation \rightarrow less than 10% precision

Lattice size	Physical size	а	L _s	Gauge action	Pion mass
$32^3 \times 64$	4.6 fm ³	0.135 fm	32	DSDR	171 241 MeV
48 ³ × 96	5.5 fm ³	0.115 fm	16	Iwasaki	135 MeV

Thank you for your attention !

Backup

Nucleon EDM in the SM

- CKM phase in EDM
 - No CPV phase in I-loop (|V_{dq}|² : no phase) and 2-loop diagram (unitarity).



• Three-loop order(short) or pion loop correction (long) $d_N^{\text{KM short}} \sim -10^{-34} \,\text{e} \cdot \text{cm}, \quad d_N^{\text{KM long}} \sim 10^{-30} - 10^{-32} \,\text{e} \cdot \text{cm}$

 $\Rightarrow d_N^{\rm KM} = d_N^{\rm KM\,short} + d_N^{\rm KM\,long} \simeq 10^{-30} - 10^{-32}\,{\rm e\cdot cm}$

6-order magnitude below the experimental upper limit.

• θ term in the QCD Lagrangian

Renormalizable and CP-violating from in topological charge density.

$$\mathcal{L}_{\theta} = \bar{\theta} \frac{1}{64\pi^2} G \widetilde{G}, \quad \bar{\theta} = \theta + \arg \det M, \quad \arg \det M \sim \eta \frac{m_u m_d m_s}{m_u m_d + m_d m_s + m_s m_u}$$

- m heta term has been estimated as $ar heta < 10^{-9\pm 1}$ Crewther, et al. (1979), Ellis, Gaillard (1979)
- Unnatural cancellation (strong CP problem)

Nucleon EDM in the BSM

- Possible higher dimension operators
 - In supersymmery (SUSY) model there is CPV phase from 1-loop (Im(g_L g_R^{*})≠0)



 \blacktriangleright CPV effective Hamiltonian with higher dimension than θ term

$$H_{CP} = \sum_{k} C_{k}(\mu) \mathcal{O}_{k} \qquad \qquad \mathcal{O}_{qEDM} = d_{q} \bar{q} (\sigma \cdot F) \gamma_{5} q \quad : \text{Quark-photon (5-dim)} \\ \mathcal{O}_{cEDM} = d_{q}^{c} \bar{q} (\sigma \cdot G) \gamma_{5} q \quad : \text{Quark-gluon (5-dim)} \\ \mathcal{O}_{Weinberg} = d^{G} G G \tilde{G} \quad : \text{Pure gluonic (6-dim)} \end{cases}$$

Contribution to nEDM in low energy model

BChPT:

QCD sum rule:



Lattice methods on EDM

Spectrum method

- I. S.Aoki and A. Gocksch, Phys. Rev. Lett. 63, 1125 (1989).
- 2. S.Aoki, A. Gocksch, A.V. Manohar, S. R. Sharpe, Phys. Rev. Lett. 65, 1092 (1990), in which they discussed about the possible lattice artifact in ref.1 results
- 3. ES, et al., for CP-PACS collaboration, Phys. Rev. D75, 034507 (2007)
- 4. ES, S. Aoki, Y. Kuramashi, Phys. Rev. D78, 014503 (2008)

Form factor

- I. ES, et al., for CP-PACS collaboration, Phys. Rev. D72, 014504 (2005).
- 2. Berruto, et al. for RBC collaboration, Phys. Rev. D73, 05409 (2006).
- 3. ES et al., Lattice 2008.

Imaginary θ

- I. T. Izubuchi, Lattice 2007.
- 2. Horsley et al., arXiv:0808.1428 [hep-lat]

EM form factor



36

