Study of the $a_0(980)$ on the lattice 31st International Symposium on Lattice Field Theory – Mainz, Germany

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Introduction, motivation (1)

- The nonet of light scalar mesons $(J^P = 0^+)$
 - $\sigma \equiv f_0(500), I = 0, 400...550 \,\mathrm{MeV},$
 - $-~\kappa \equiv K_0^*(800)$, I=1/2 , $682\pm29~{\rm MeV}$,
 - $a_0(980)$, $f_0(980)$, I = 1, $980 \pm 20 \text{ MeV}$, $990 \pm 20 \text{ MeV}$

is poorly understood:

- All nine states are unexpectedly light (should rather be close to the corresponding $J^P = 1^+, 2^+$ states around $1200 \dots 1500$ MeV).
- The ordering of states is inverted compared to expectation:
 - * E.g. in a $q\bar{q}$ picture the I = 1 states $a_0(980)$, $f_0(980)$ must necessarily be formed by two u/d quarks, while the $I = 1/2 \kappa$ states are made from an s and a u/d quark; since $m_s > m_{u/d}$ one would expect $m(\kappa) > m(a_0(980)), m(f_0(980))$.

Introduction, motivation (2)

- * In a tetraquark picture the quark content could be the following: $\kappa \equiv \bar{s}l\bar{l}l$, while $a_0(980), f_0(980) \equiv \bar{s}l\bar{l}s$; this would naturally explain the observed ordering.
- Certain decays also support a tetraquark interpretation: e.g. $a_0(980)$ readily decays to $K + \bar{K}$, which indicates that besides the two light quarks required by I = 1 also an $s\bar{s}$ pair is present.
- \rightarrow Study these states by means of lattice QCD to confirm or to rule out their interpretation in terms of tetraquarks.
- Examples of heavy mesons, which are tetraquark candidates:
 - $D_{s0}^*(2317)^{\pm} (I(J^P) = 0(0^+)), D_{s1}(2460)^{\pm} (I(J^P) = 0(1^+)),$
 - charmonium states X(3872), $Z(4430)^{\pm}$, $Z(4050)^{\pm}$, $Z(4250)^{\pm}$, ...

Outline

(1) Wilson twisted mass study of $a_0(980)$.

[C. Alexandrou et al. [ETM Collaboration], JHEP 1304, 137 (2013) [arXiv:1212.1418 [hep-lat]]]

- (2) Recent advances:
 - Lattice discretization changed, now Wilson + clover fermions (generated by the PACS-CS Collaboration).

[S. Aoki et al. [PACS-CS Collaboration], Phys. Rev. D 79, 034503 (2009) [arXiv:0807.1661 [hep-lat]]]

- Inclusion of disconnected diagrams.

Part 1: Wilson twisted mass study of $a_0(980)$

Tetraquark creation operators

- $a_0(980)$:
 - Quantum numbers $I(J^{PC}) = 1(0^{++})$.
 - Mass 980 ± 20 MeV.
- Tetraquark creation operators:
 - Need two light quarks due to I = 1, e.g. $u\bar{d}$.
 - $a_0(980)$ decays to $KK \dots$ suggests an $s\bar{s}$ component.
 - Molecule type (models a bound $K\bar{K}$ state):

$$\mathcal{O}_{a_0(980)}^{Kar{K} ext{ molecule }} = \sum_{\mathbf{x}} \Big(ar{s}(\mathbf{x}) \gamma_5 u(\mathbf{x}) \Big) \Big(ar{d}(\mathbf{x}) \gamma_5 s(\mathbf{x}) \Big).$$

- Diquark type (models a bound diquark-antidiquark):

$$\mathcal{O}_{a_0(980)}^{\mathsf{diquark}} = \sum_{\mathbf{x}} \left(\epsilon^{abc} \bar{s}^b(\mathbf{x}) C \gamma_5 \bar{d}^{c,T}(\mathbf{x}) \right) \left(\epsilon^{ade} u^{d,T}(\mathbf{x}) C \gamma_5 s^e(\mathbf{x}) \right).$$

Two-particle creation operators (1)

• There are two-particle states, which have the same quantum numbers as $a_0(980)$, $I(J^{PC}) = 1(0^{++})$,

$$-K + \overline{K} (m(K) \approx 500 \,\mathrm{MeV}),$$

 $-\eta_s + \pi \ (m(\eta_s \equiv \bar{s}\gamma_5 s) \approx 700 \text{ MeV}, \ m(\pi) \approx 300 \text{ MeV}$ in our lattice setup),

which are both around the expected $a_0(980)$ mass 980 ± 20 MeV.

• To determine, whether there is a bound $a_0(980)$ tetraquark state, we need to resolve the above listed two-particle states $K + \bar{K}$ and $\eta_s + \pi$ and check, whether there is an additional 3rd state in the mass region around 980 ± 20 MeV; to this end we need operators of two-particle type.

Two-particle creation operators (2)

- Two-particle operators:
 - $\text{Two-particle } K + \bar{K} \text{ type:} \\ \mathcal{O}_{a_0(980)}^{K+\bar{K} \text{ two-particle}} = \Big(\sum_{\mathbf{x}} \bar{s}(\mathbf{x})\gamma_5 u(\mathbf{x})\Big)\Big(\sum_{\mathbf{y}} \bar{d}(\mathbf{y})\gamma_5 s(\mathbf{y})\Big).$
 - Two-particle $\eta_s + \pi$ type:

$$\mathcal{O}_{a_0(980)}^{\eta_s+\pi ext{ two-particle }} = \Big(\sum_{\mathbf{x}} ar{s}(\mathbf{x}) \gamma_5 s(\mathbf{x}) \Big) \Big(\sum_{\mathbf{y}} ar{d}(\mathbf{y}) \gamma_5 u(\mathbf{y}) \Big).$$

Wilson twisted mass lattice setup

- Gauge link configurations generated by the ETM Collaboration. [R. Baron *et al.*, JHEP 1006, 111 (2010) [arXiv:1004.5284 [hep-lat]]]
- 2+1+1 dynamical Wilson twisted mass quark flavors, i.e. u, d, s and c sea quarks.
- Iwasaki gauge action.
- Lattice spacing $a \approx 0.086$ fm.
- Various lattice volumes.
- Various light u/d quark masses corresponding pion masses $m_{\pi} \approx 280 \dots 460 \text{ MeV}.$
- APE smeared links and Gaussian smeared quark fields to improve the overlap of trial states to the low lying states.
- Singly disconnected contributions neglected, i.e. no *s* quark propagation within the same timeslice.

Numerical results $a_0(980)$ (1)

- Study all four operators ($K\bar{K}$ molecule, diquark, $K + \bar{K}$ two-particle, $\eta_s + \pi$ two-particle) at the same time, extract the four lowest energy eigenstates by diagonalizing a 4×4 correlation matrix.
 - Only two low-lying states around $980\pm20\,{\rm MeV},$ the 2nd and 3rd excitation are $\approx750\,{\rm MeV}$ heavier.
 - The signal of the low-lying states is of much better quality than when only considering tetraquark operators
 - ightarrow suggests that the observed low-lying states have much better overlap to the two-particle operators and are most likely of two-particle type.



Numerical results $a_0(980)$ (2)

- When determining low-lying eigenstates from a correlation matrix, one does not only obtain their mass, but also information about their operator content, i.e. which percentage of which operator is present in an extracted state:
 - \rightarrow The ground state is a $\eta_s + \pi$ state ($\gtrsim 95\%$ two-particle $\eta_s + \pi$ content).
 - \rightarrow The first excitation is a $K+\bar{K}$ state ($\gtrsim\!95\%$ two-particle $K+\bar{K}$ content).



Numerical results $a_0(980)$ (3)

- What about the 2nd and 3rd excitation? ... Are these tetraquark states? ... What is their nature?
- Two-particle states with one relative quantum of momentum (one particle has momentum $+p_{\min} = +2\pi/L$ the other $-p_{\min}$) also have quantum numbers $I(J^{PC}) = 1(0^{++})$; their masses can easily be estimated:
 - $p_{\min} = 2\pi/L \approx 715 \text{ MeV}$ (the results presented correspond to the small lattice with spatial extension L = 1.73 fm);
 - $m(K(+p_{\min}) + \bar{K}(-p_{\min})) \approx 2\sqrt{m(K)^2 + p_{\min}^2} \approx 1750 \text{ MeV};$ $- m(\eta(+p_{\min}) + \pi(-p_{\min})) \approx \sqrt{m(\eta)^2 + p_{\min}^2} + \sqrt{m(\pi)^2 + p_{\min}^2} \approx 1780 \text{ MeV};$

these estimated mass values are consistent with the observed mass values of the 2nd and 3rd excitation

 \rightarrow suggests to interpret these states as two-particle states.



Numerical results $a_0(980)$ (4)

• Summary:

- In the $a_0(980)$ sector (quantum numbers $I(J^{PC}) = 1(0^{++})$) we do not observe any low-lying (mass $\leq 1750 \text{ MeV}$) tetraquark state, even though we employed operators of tetraquark structure ($K\bar{K}$ molecule, diquark).
- The experimentally measured mass for $a_0(980)$ is 980 ± 20 MeV.
- Conclusion: $a_0(980)$ does not seem to be a strongly bound tetraquark state (either of molecule or of diquark type) ... maybe an ordinary quark-antiquark state or a rather unstable resonance.

Part 2: Recent advances

Wilson + clover lattice setup

- Gauge link configurations generated by the PACS-CS Collaboration.
 [S. Aoki et al. [PACS-CS Collaboration], Phys. Rev. D 79, 034503 (2009) [arXiv:0807.1661 [hep-lat]]]
- 2+1 dynamical Wilson + clover quark flavors, i.e. u, d and s sea quarks.
 → In contrast to twisted mass parity and isospin are exact symmetries, i.e. no pion and kaon mass splitting, easy separation of P = +, states, ...
- Iwasaki gauge action.
- Lattice spacing $a \approx 0.09$ fm, lattice volume $(L/a)^3 \times T/a = 32^3 \times 64$.
- Light u/d quark masses corresponding to pion masses $m_{\pi} \approx 150 \text{ MeV}$ and $m_{\pi} \approx 300 \text{ MeV}$.

 \rightarrow Computations close to physically light u/d quark masses possible.

- APE smeared links and Gaussian smeared quark fields to improve the overlap of trial states to the low lying states.
- Singly disconnected contributions included.
 - \rightarrow s quark propagation within the same timeslice.

Singly disconnected diagrams (1)

- In our previous Wilson twisted mass study of $a_0(980)$ we neglected singly disconnected contributions:
 - \rightarrow We could not consider a $q\bar{q}$ operator,

$$\mathcal{O}^{qar{q}}_{a_0(980)} \;\;=\;\; \sum_{\mathbf{x}} \Big(ar{d}(\mathbf{x}) u(\mathbf{x}) \Big),$$

because cross correlations between this operator and any of the four-quark operators $\mathcal{O}_{a_0(980)}^{K\bar{K}}$ molecule, $\mathcal{O}_{a_0(980)}^{\text{diquark}}$, $\mathcal{O}_{a_0(980)}^{K+\bar{K}}$ two-particle or $\mathcal{O}_{a_0(980)}^{\eta_s+\pi}$ two-particle correspond to singly disconnected diagrams.

 \rightarrow Also correlations between the four-quark operators include singly disconnected diagrams; therefore, we introduced a source of systematic error, which is difficult to estimate or to control.



Singly disconnected diagrams (2)

- Technical aspects of computing singly disconnected diagrams:
 - Blue: point-to-all propagators applicable.
 - Red: due to $\sum_{\mathbf{x}}$, timeslice-to-all propagators needed.
 - Timeslice-to-all propagators can be estimated stochastically.
 - Using several stochastic timeslice-to-all propagators results in a poor signal-to-noise ratio.
 - \rightarrow Combine three point-to-all (blue) and one stochastic timeslice-to-all (red) propagator.



Singly disconnected diagrams (3)

- Effective masses from a 2×2 correlation matrix $(\mathcal{O}_{a_0(980)}^{q\bar{q}})$ and $\mathcal{O}_{a_0(980)}^{K\bar{K}}$ molecule):
 - Lowest (two) energy level(s) consistent with K + K, $\eta + \pi$ and a possibly existing additional $a_0(980)$ state.
 - For physically interesting statements we also need to include $\mathcal{O}_{a_0(980)}^{\text{diquark}}$, $\mathcal{O}_{a_0(980)}^{K+\bar{K} \text{ two-particle}}$ and $\mathcal{O}_{a_0(980)}^{\eta_s+\pi \text{ two-particle}}$ (work in progress).



Outlook

- Enlarge correlation matrices such that
 - $-~q\bar{q}$ operators,
 - tetraquark operators (mesonic molecules, diquark-antidiquark pairs),
 - two-meson operators

are included.

- Perform computations at pion mass $m_{\pi} \approx 150 \text{ MeV}$.
- Address various physical questions/systems (tetraquark candidates with different flavor structure, search for additional bound states, ...).