

*Eigenvalue spectrum of  
lattice  $\mathcal{N} = 4$  super Yang-Mills*

*(work in progress)*

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# Outline

1. Introduction
2.  $\mathcal{N} = 4$  SYM on the lattice
3. Techniques used
4. Pfaffian
5. Eigenvalue spectrum – general results
6. Mass anomalous dimension
7. Crude attempt at Patella's calculation
8. Conclusions

# Introduction

- $\mathcal{N} = 4$  SYM is of great theoretical interest
  - Very elegant theory; maximal supersymmetry
  - AdS/CFT correspondence
  - BSM physics, phenomenology
  - Connections to string theory
- Heavily studied analytically, simulations are much more recent
- Putting supersymmetry on the lattice is a fascinating problem
- Need to find a way to deal with breaking of translational invariance
  - Twisting procedure, Dirac-Kähler fermions, simulate twisted theory
  - Orbifolding procedure; produces similar action
- This is, however, not a technical talk about how to formulate SUSY on the lattice

- Difficult to implement (SUSY algebra is an extension of Poincaré algebra)
- Technique used here is to preserve only one of the supercharges, and recover the rest in the continuum limit
  - Anticommutator of two supercharges generates spacetime translations
  - Preserving only one supercharge sidesteps this issue
  - Resulting algebra not broken by the discretisation
- Corresponds to a topologically twisted form of the original theory
- Restriction that we work with a target theory having  $2^D$  supercharges
- This formulation of  $\mathcal{N} = 4$  SYM on the lattice was examined in [1209.5285](#)
- Dirac-Kähler fermion action that one obtains is similar to the standard staggered action (in free field theory at least)
- We vary two parameters, the 't Hooft Coupling  $\lambda$  and a mass parameter  $\mu_L$  to control fluctuations (take  $\mu_L \rightarrow 0$  later)

# Questions

1. Double-check: does this theory have a sign problem?
  - Properties of the Pfaffian we know it is 'small enough', but. . .
  - Can we infer what will happen for larger volumes?
2. What does the eigenspectrum look like?
  - Zero modes, trace modes
  - Vary 't Hooft coupling  $\lambda$  and breaking parameter  $\mu_L$
3. Can we measure the mass anomalous dimension?
  - Scaling of the individual eigenvalues with  $L$
  - Crudely fitting to the eigenmode density [1204.4432](#)

## Reminder – the Arnoldi/Lanczos algorithm

- Krylov subspace methods – standard way to calculate  $N$  extremal eigenvalues of  $M \times M$  matrix,  $M \gg N$
- The Arnoldi algorithm depends on the factorisation

$$AV = VH + fe^T; \quad V^\dagger f = 0$$

where  $A$  is our operator,  $V$  is an  $M \times N$  matrix of Ritz values and  $H$  is a Hessenberg matrix;  $f$  is a  $M$ -component vector

- Lanczos:  $A$  is Hermitian,  $H$  is tridiagonal
- Given the factorisation, the eigenvalues of  $H$  are calculated; these Ritz values approximate some eigenvalues of  $A$
- Choose the ‘best’ eigenvalues in  $H$ , throw away the rest, repeat
- How reliable is the algorithm?
  - Difficult to define convergence for non-Hermitian operators
  - Gershgorin circle theorem underpins Hermitian results theoretically

## Technical details

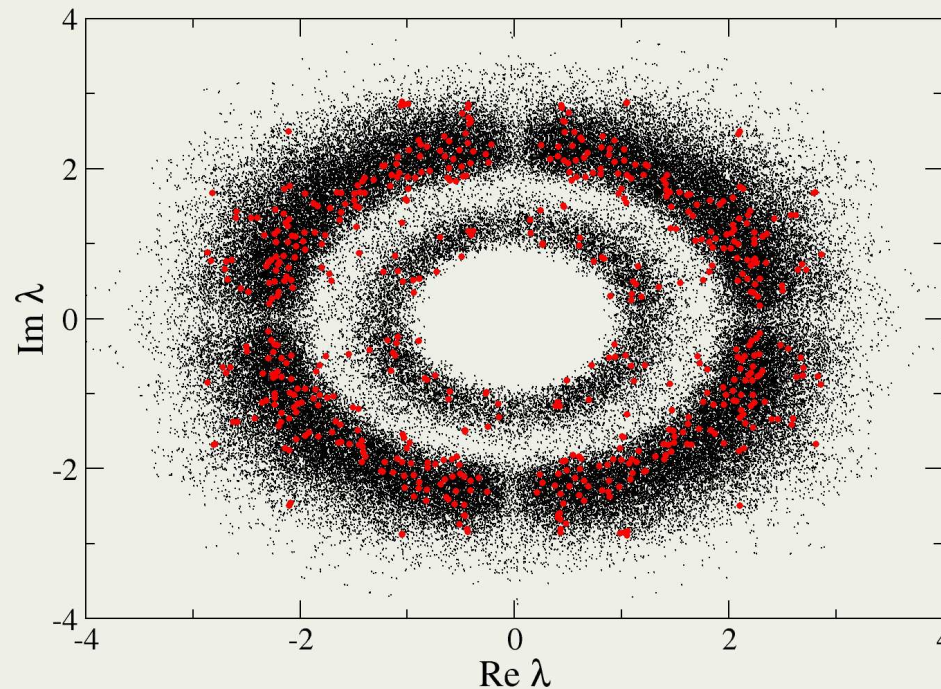
- The simulation code is not (yet) parallelised (coming soon!) which restricts the volumes we can reach
- There is support for GPUs in the conjugate gradient algorithm but we instead make use of grid resources
  - Problem trivially parallelises; not time-critical
  - Therefore use FGI (Finnish Grid Infrastructure) rather than a traditional cluster
- Limited volumes mean we can get away with rudimentary eigenvalue calculations
  - We use ARPACK for the eigenvalue calculations<sup>1</sup>
  - In general, calculate 200 smallest eigenvalues, taking about 2 hours

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<sup>1</sup>See Joni Suorsa's talk, parallel D (at 15:00, i.e. next talk) for a fully parallel, GPU-accelerated non-Hermitian eigenvalue calculation

## Briefly: measuring the Pfaffian

- We cannot take for granted that the theory does not have a sign problem!
- Can measure exactly in small volumes (see [1209.5285](#))

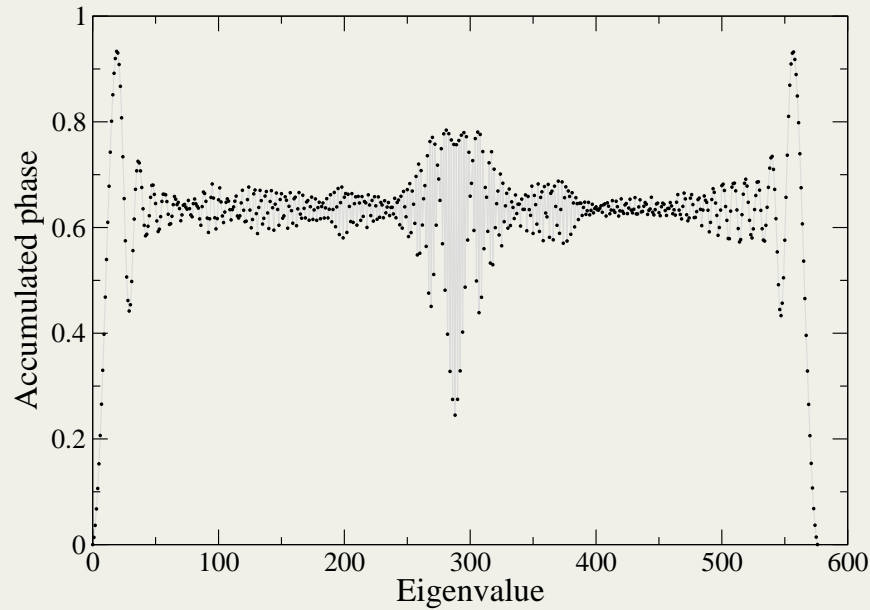


apbc,  $3^4$ , 500 configurations, 1 configuration in red

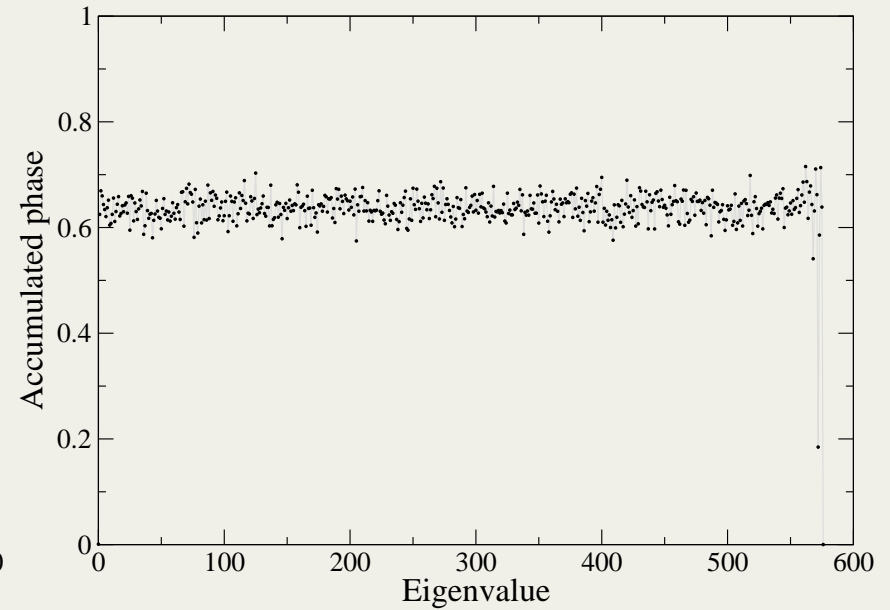
- Looks nice, but phase only measurable if we know *full* eigenspectrum?
- Any structure at extremal eigenvalues that tells us about the sign, that we could examine at larger volumes?



# Pfaffian analysis



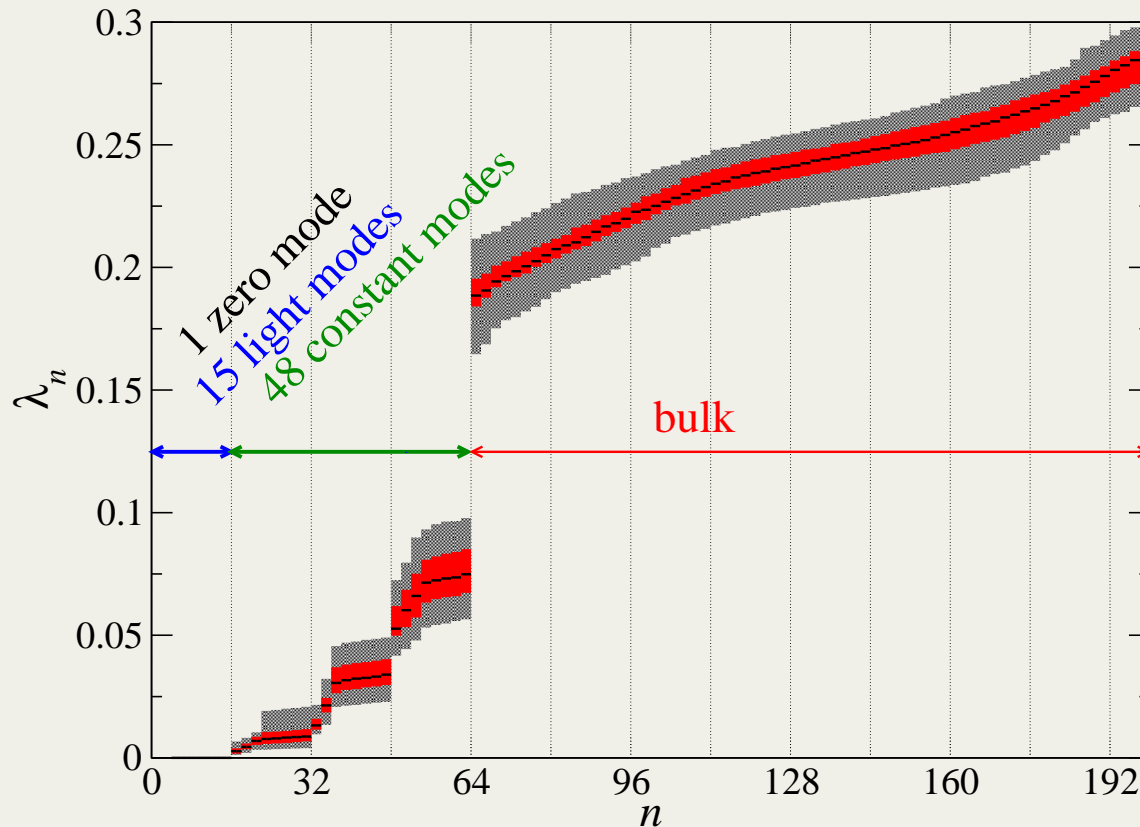
(a) by phase



(b) by magnitude

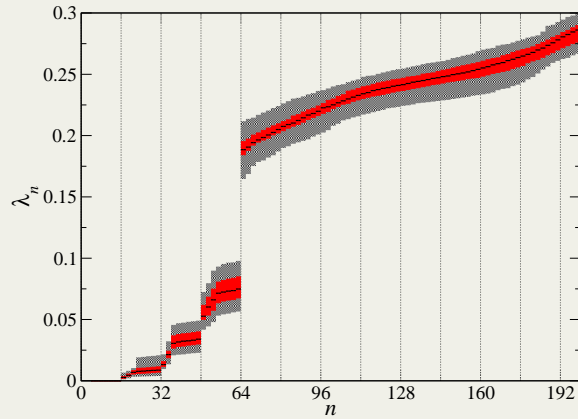
No helpful extremal structure, but at small volumes, the phase is unimportant.

# Eigenvalues: what do we see?

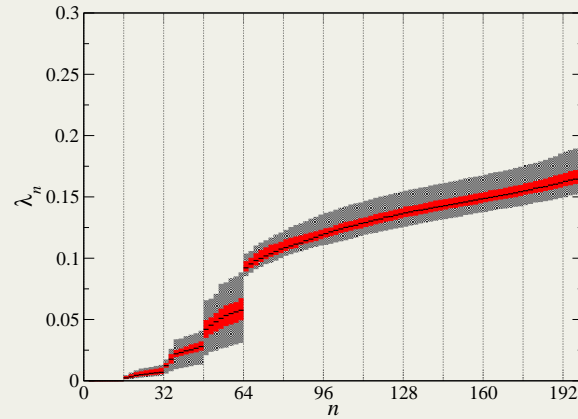


- $U(2) \rightarrow 2 \times 2 \times 16 = 48$  total would-be zero (trace) modes ( $U(3)$  theory yields 144 such modes)
- We therefore discard these modes when carrying out further analysis

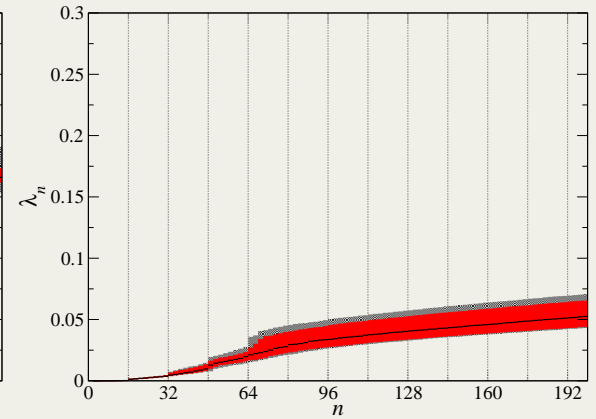
# Eigenspectra



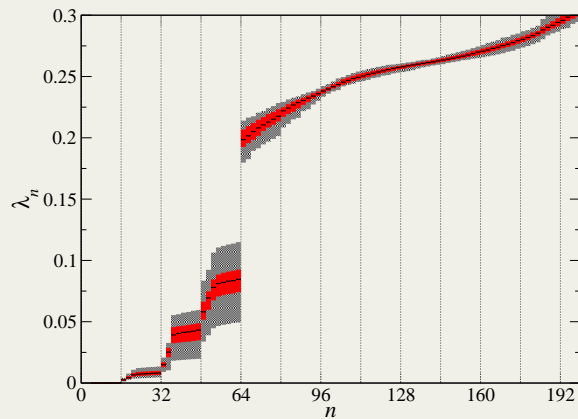
(c)  $\lambda = 0.5, \mu_L = 0.5$



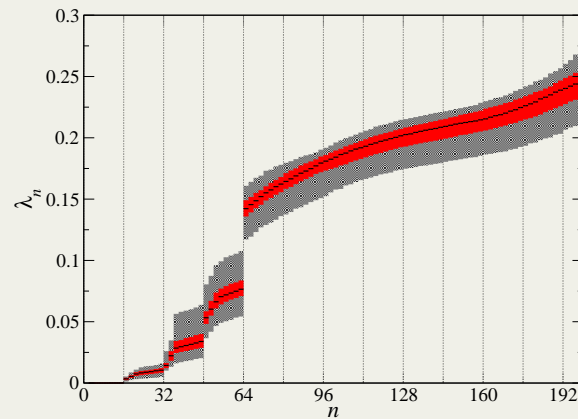
(d)  $\lambda = 1.0, \mu_L = 0.5$



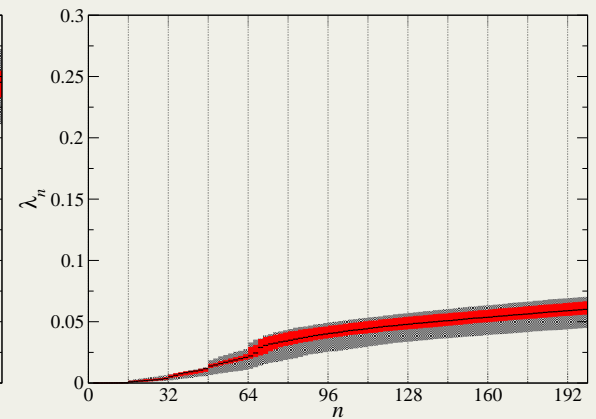
(e)  $\lambda = 2.0, \mu_L = 0.5$



(f)  $\lambda = 0.5, \mu_L = 1.0$



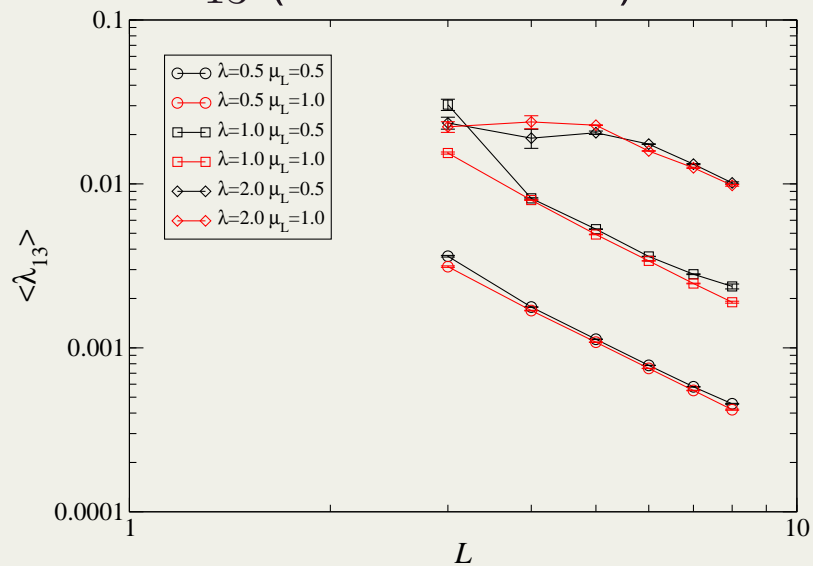
(g)  $\lambda = 1.0, \mu_L = 1.0$



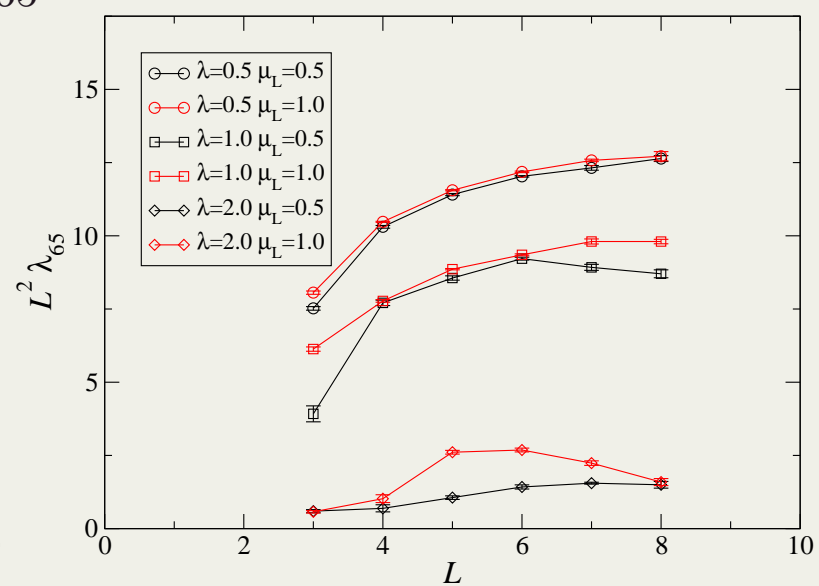
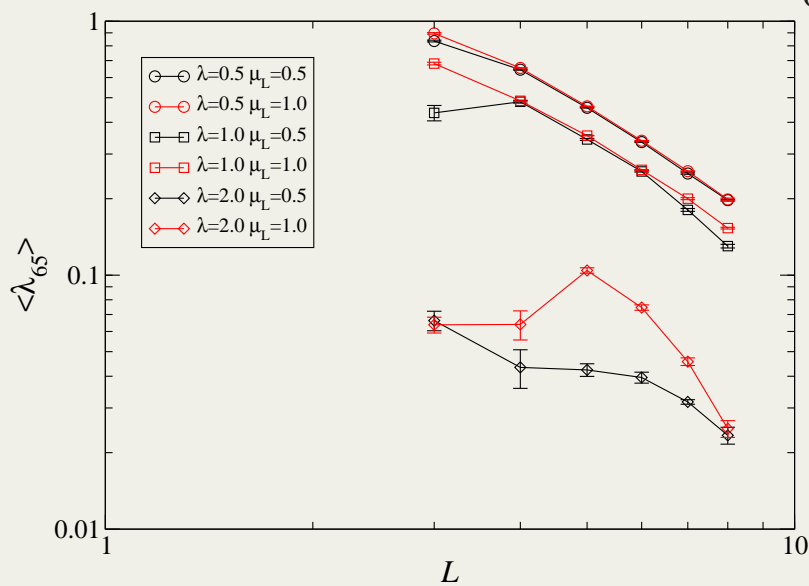
(h)  $\lambda = 2.0, \mu_L = 1.0$

# Scaling with L

$\lambda_{13}$  (scales  $\sim L^{-4}$ ):



$\lambda_{65}$ :



# Mass anomalous dimension

- Only observable in CFTs
- Important for phenomenology – can we reproduce the walking by perturbing the theory?
- Comparisons with continuum predictions
- Dependence of  $\gamma_*$  on  $\lambda$ ?

# Finite volume scaling of eigenvalues [0910.3072](#)

- Fit  $\langle \lambda_n(L) \rangle$  to  $C_n L^{-y_n}$ , with jackknifed errors
- Distinct change in behaviour at  $\lambda_{65}$

$n$	$\lambda = 0.5$		$\lambda = 1.0$	
	$\mu_L = 0.5$	$\mu_L = 1.0$	$\mu_L = 0.5$	$\mu_L = 1.0$
9	$3.97 \pm 0.02$	$4.03 \pm 0.01$	$3.94 \pm 0.14$	$4.09 \pm 0.03$
17	$1.57 \pm 0.01$	$1.83 \pm 0.03$	$1.75 \pm 0.52$	$1.76 \pm 0.24$
25	$1.64 \pm 0.02$	$1.86 \pm 0.03$	$1.83 \pm 0.46$	$1.84 \pm 0.21$
33	$2.33 \pm 0.15$	$2.17 \pm 0.04$	$2.29 \pm 0.32$	$2.20 \pm 0.06$
41	$2.41 \pm 0.20$	$2.25 \pm 0.03$	$2.36 \pm 0.42$	$2.29 \pm 0.07$
49	$2.27 \pm 0.13$	$2.16 \pm 0.06$	$2.21 \pm 0.21$	$2.19 \pm 0.02$
57	$2.18 \pm 0.11$	$2.14 \pm 0.05$	$2.13 \pm 0.16$	$2.13 \pm 0.02$
65	$1.74 \pm 0.04$	$1.74 \pm 0.02$	$1.82 \pm 0.18$	$1.70 \pm 0.01$
73	$1.75 \pm 0.04$	$1.74 \pm 0.03$	$1.84 \pm 0.15$	$1.69 \pm 0.02$
81	$1.75 \pm 0.03$	$1.74 \pm 0.03$	$1.83 \pm 0.12$	$1.69 \pm 0.03$

# Mass anomalous dimension from eigenmode density

- Working with the ansatz ...

$$a^{-4}\bar{\nu}(\Omega) \approx a^{-4}\bar{\nu}_0 + A [(a\Omega)^2 - (am)^2]^{2/(1+\gamma_*)}$$

- ... where  $\bar{\nu}(\Omega)$  is the integrated spectral density ...

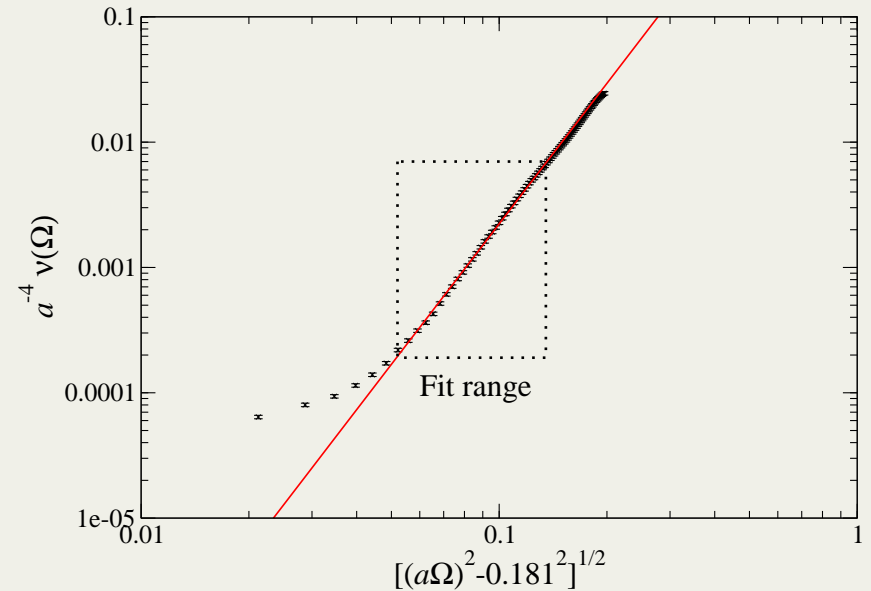
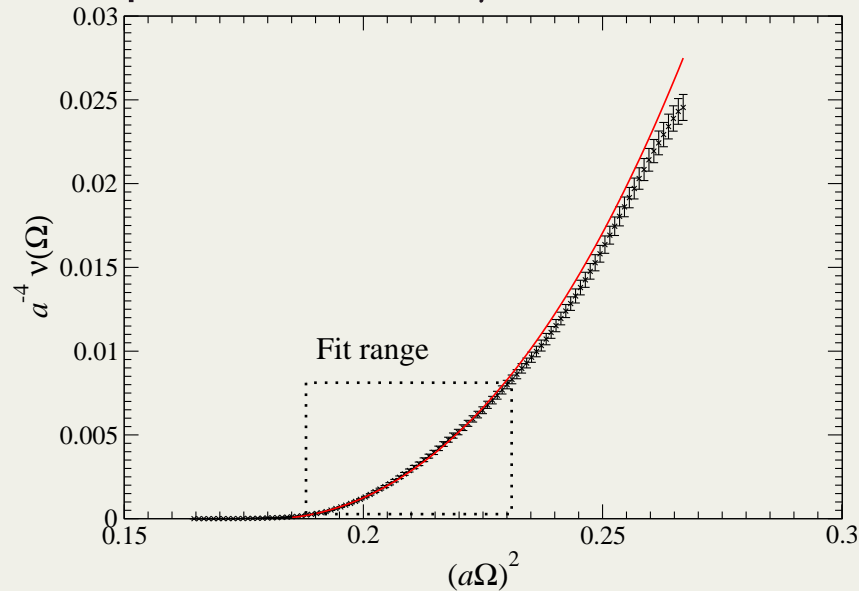
$$\bar{\nu}(\Omega) = 2 \int_0^{\sqrt{\Omega^2 - m^2}} \rho(\omega) d\omega; \quad \rho(\omega) = \lim_{V \rightarrow \infty} \frac{1}{V} \sum_k \langle \delta(\omega - \omega_k) \rangle$$

- ... we optimistically fit our largest lattice data ( $L = 8$  – albeit ‘staggered’) to the ansatz and see if anything meaningful results
- Discard first 64 eigenvalues (zero modes, trace, constant)
- The term  $a^{-4}\bar{\nu}_0$  is (therefore) consistent with zero in all our fits
- $\gamma_*$  can also be obtained by fitting to our data for  $\langle \lambda_n(L) \rangle$

# Eigenmode density fitting – results

1. Fit to all possible ranges
2. Find the fitting range which yields  $\chi^2/\text{dof}$  closest to unity
3. Examine the effect on  $\gamma_*$  of varying the lower and upper bounds
4. If the results lie on a ‘plateau’, we use that value of  $\gamma_*$

Example –  $\lambda = 0.5$ ,  $\mu_L = 0.5$ :



- For  $\lambda = 0.5$ ,  $\mu_L = 0.5$ , we get  $\gamma_* \approx 0.07$
- ... which is consistent with fitting to  $\langle \lambda_{65}(L) \rangle$ , which yields  $\gamma_* = 0.08 \pm 0.01$



# Conclusions

- In this talk we have looked briefly at
  - The Pfaffian
  - Basic properties of the eigenspectrum; scaling with  $L$
  - Fitting the spectral density
- As far as we can tell the phase of the Pfaffian is close to zero. . .
- . . . unfortunately, to say for sure, we would need to measure it exactly (hard in large volumes)
- We see zero modes, trace modes (that would disappear in continuum) and then the actual physics
- We have (optimistically) attempted to measure the anomalous mass dimension by fitting the spectral density
  - Although the volume is still relatively small
  - And our understanding of the trace modes is still very rudimentary
- typical values are around  $\gamma_* \approx 0.1$ .

## Still to do

- Parallelised code
  - We have reached the limit of what can be done with single-core CPU code, even with the trivial parallelism used here (roughly 10k hours per parameter choice at  $L = 8$ , total use  $> 100k$  hours)
  - Larger volumes
  - Better performance without needing GPU acceleration?
- Better control of fits to spectral density
- Stronger coupling
  - Anomalous dimensions seem to be still quite small for  $\lambda = 1$
  - Results not of sufficiently high quality to draw conclusions for  $\lambda = 2$