

Spectroscopy of doubly and triply-charmed baryons from lattice QCD

Padmanath M.

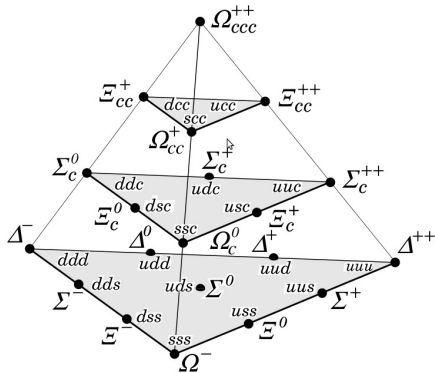
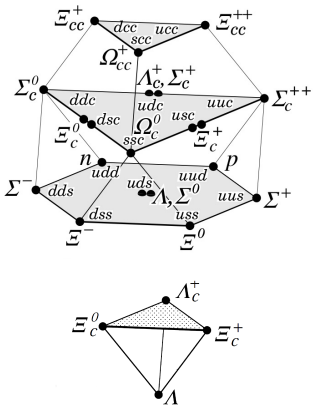


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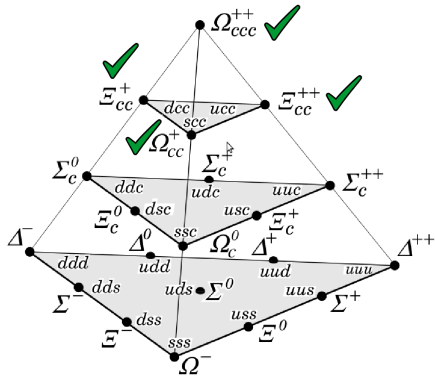
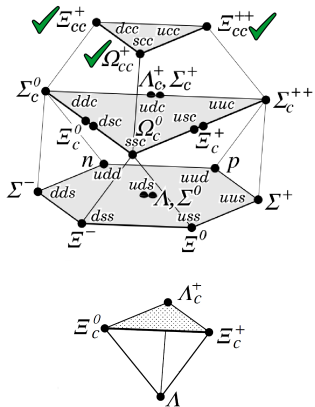
- arXiv:1307.7022 [hep-lat] .
- In collaboration with R. G. Edwards, N. Mathur and M. Peardon.
- Computations performed on computational facilities at DTP, TIFR, Mumbai, Jefferson Laboratory and TCHPC, Trinity College, Dublin.

4 (u, d, s, c) degenerate flavors



We have one heavy and 2+1 light flavor quarks.

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Ensemble details

Calculations performed on lattices generated by **Hadron Spectrum Collaboration**.

- Dynamical configurations ($N_f = 2 + 1$).
- Anisotropic lattices with $\xi = a_s/a_t \sim 3.5$.
- Scale set via m_Ω : $a_s = 0.12$ fm
- Lattice size : $16^3 \times 128$.
- Statistics : 96 cfgs and 4 time sources.
- **Clover** fermions : Non-perturbative $O(a)$ improvement.
- Spatial links are **stout smeared**.
- Quark fields are **distilled**.

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- Spatial links are stout smeared.
- Quark fields are distilled.
- Caveat : Pion mass ~ 391 MeV.

Interpolating operators

 Ω_{ccc} **Non-Rel:** $SU(6) \otimes O(3)$

D \ J	1/2	3/2	5/2	7/2
0	0	1	0	0
1	1	1	0	0
2_{hybrid}	1	1	0	0
2	2	3	2	1

 Ω_{cc} and Ξ_{cc} **Non-Rel:** $SU(6) \otimes O(3)$

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Whole operator set

Ω_{ccc}	G_1		H		G_2	
	g	u	g	u	g	u
Total	20	20	33	33	12	12
Hybrid	4	4	5	5	1	1
NR	4	1	8	1	3	0

$g \rightarrow +$ $u \rightarrow -$

Whole operator set

Ξ_{cc}	G_1		H		G_2	
	g	u	g	u	g	u
Total	55	55	90	90	35	35
Hybrid	12	12	16	16	4	4
NR	11	3	19	4	8	1

Generalized eigenvalue problem

Using this large operator basis, with definite J^P in the continuum limit, to build the correlation matrix

$$C_{ij}(t) = \langle 0 | \mathcal{O}_i(t) \mathcal{O}_j^\dagger(0) | 0 \rangle = \sum_n \frac{Z_i^n Z_j^{n\dagger}}{2E_n} \exp^{-E_n t}$$

Solving the generalized eigenvalue problem for this correlation matrix

$$C_{ij}(t) v_j^{(n)}(t, t_0) = \lambda^{(n)}(t, t_0) C_{ij}(t_0) v_j^{(n)}(t, t_0)$$

- Principal correlators given by eigenvalues
 $\lambda_n(t, t_0) \sim (1 - A_n) \exp^{-m_n(t-t_0)} + A_n \exp^{-m'_n(t-t_0)}$
- Eigenvectors related to the overlap factors
 $Z_i^{(n)} = \langle 0 | \mathcal{O}_i | n \rangle = \sqrt{2E_n} \exp^{E_n t_0/2} v_j^{(n)\dagger} C_{ji}(t_0)$

Spin identification

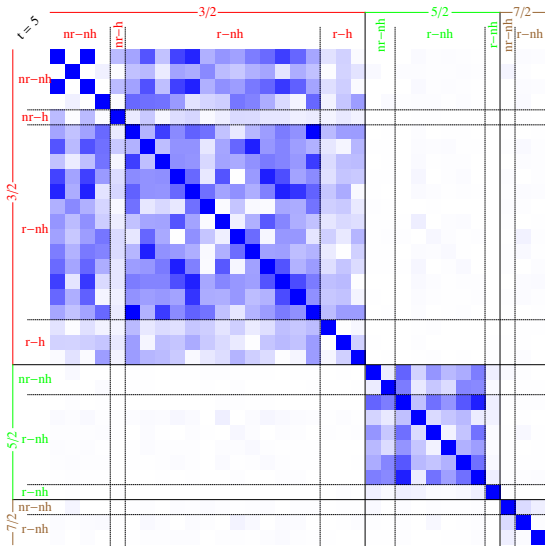
Discretized space-time breaks rotational symmetry down to octahedral symmetry.

- Continuum spin operators subduced to lattice irreps.
- G_1 , H and G_2 : O_h irreps representing half spin.

Λ	d_Λ	J
G_1	2	1/2, 7/2, 9/2, ...
H	4	3/2, 5/2, 7/2, ...
G_2	2	5/2, 7/2, 9/2, ...

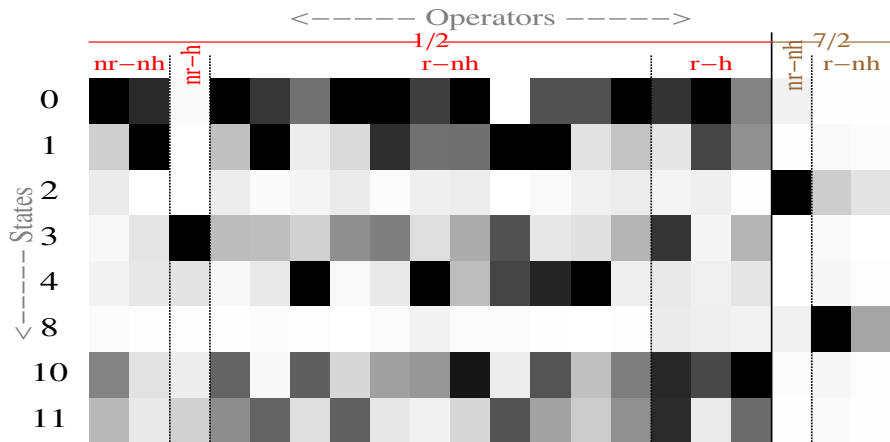
- Subduced operators carry a memory of the continuum spin J .
- An operator of spin J overlaps mainly with states of spin J .
Overlap factors to identify spin of states.

ccc correlation matrix plot (H^g ; at $t=5$) : $C_{ij}/\sqrt{C_{ii}C_{jj}}$



- $nr - nh$ = non - relativistic & non - hybrid
- $nr - h$ = non - relativistic & hybrid
- $r - nh$ = relativistic & non - hybrid
- $r - h$ = relativistic & hybrid

Spin identification using overlap factors : (ccc, G_1^g)



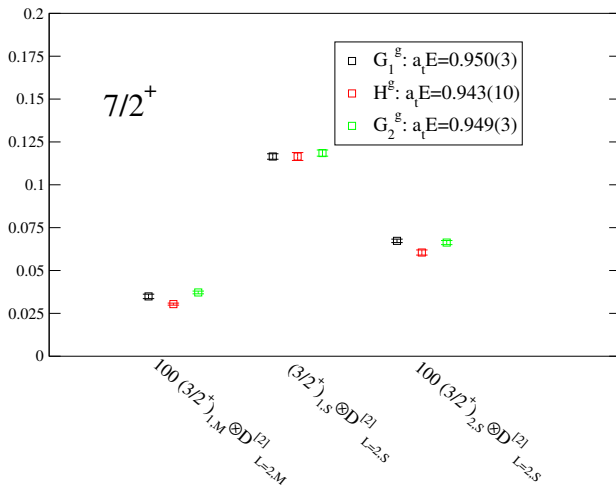
$nr-nh = non-relativistic \& non-hybrid$

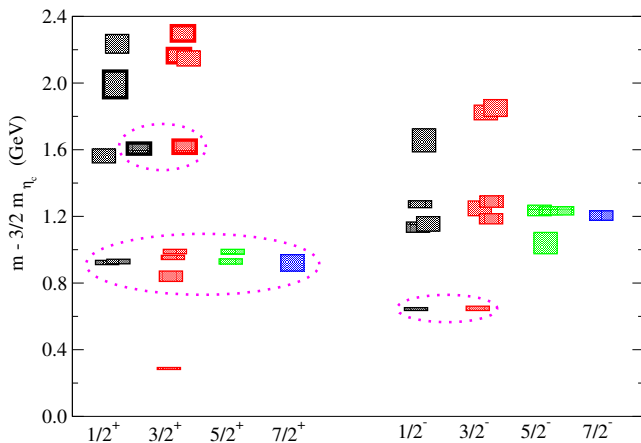
$nr-h = non-relativistic \& hybrid$

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$r-h = relativistic \& hybrid$

Spin identification across multiple irreps : $7/2^+$





arXiv:1307.7022 [hep-lat]

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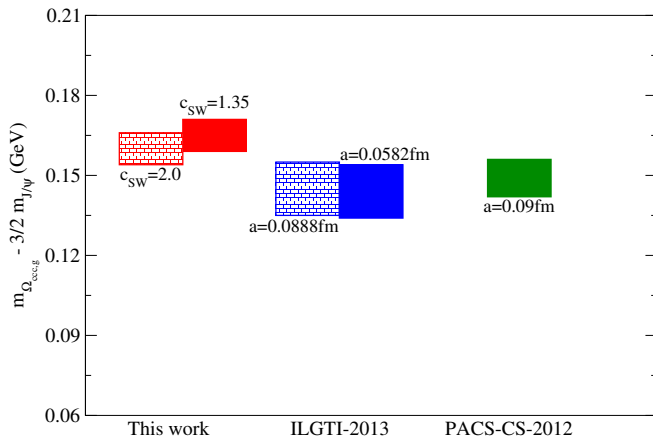
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Ω_{cc} and Ξ_{cc}

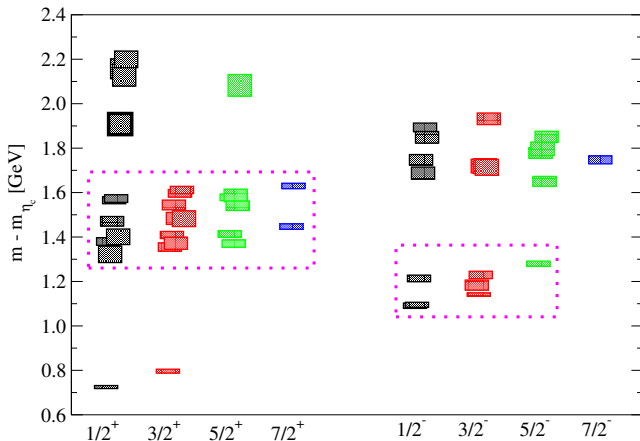
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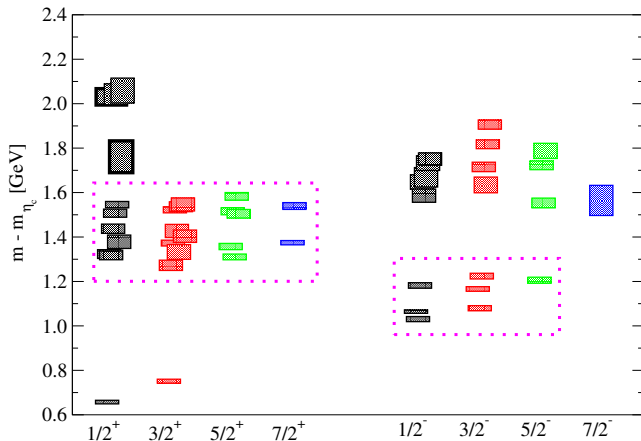
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$\Omega_{ccc} (3/2^+)$ ground state : discretization errors



Ω_{cc} spectrum





Interpolating operators

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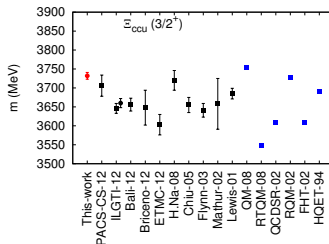
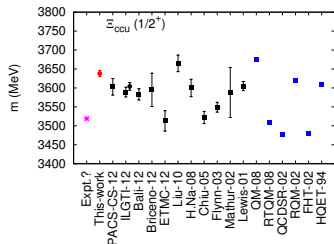
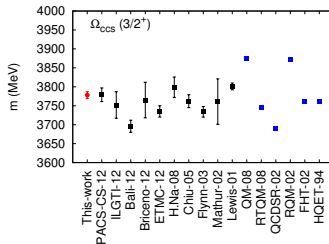
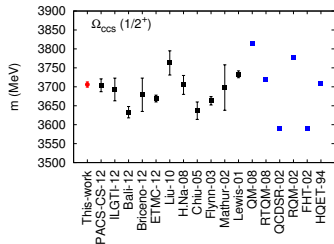
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cc(q) ground states



● → This work

■ → Other lattice results

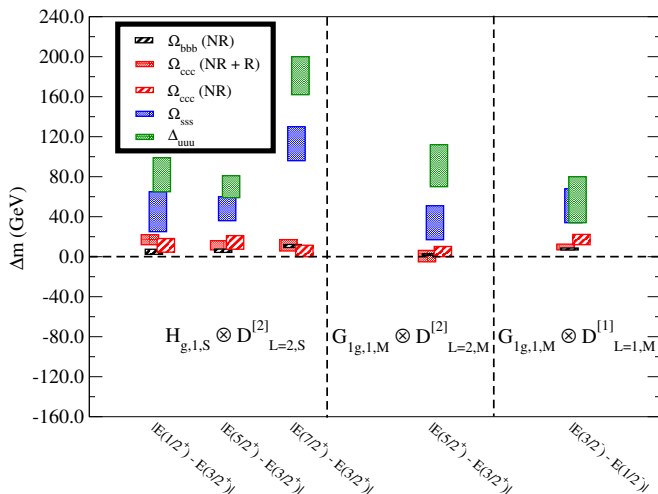
■ → EFT and model calculations

★ → Experiment

m_q dependence of energy splittings

- Spin-Orbit interactions inversely proportional to m_q^2 .
Vanishes in the heavy quark limit.
Degeneracy lifts : a measure of heavyness of the quark mass.
- Binding energy quark mass dependence.
Mass of a hadron with n heavy quarks: $M_{H_{nq}} = nm_Q + A + B/m_Q + \mathcal{O}(1/m_Q^2)$.
Energy splittings : $a + b/m_Q + \mathcal{O}(1/m_Q^2)$.
Fits with heavy quark inspired functional forms.
- From energy splittings ($\Xi_{cc}^* - D_c$, $\Omega_{cc}^* - D_s$ and $\Omega_{ccc} - \eta_c$) and ($\Xi_{cc}^* - D_c^*$, $\Omega_{cc}^* - D_s^*$ and $\Omega_{ccc} - J/\psi$), we extrapolate to bottom mass and get $B_c^* - B_c = 80 \pm 8$ MeV and $\Omega_{ccb}^* = 8050 \pm 10$ MeV.

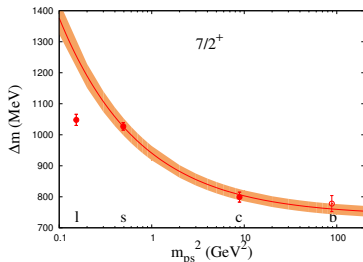
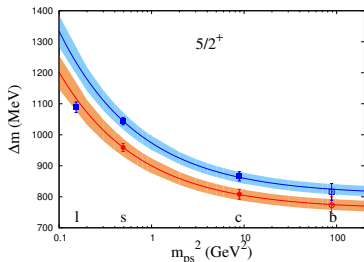
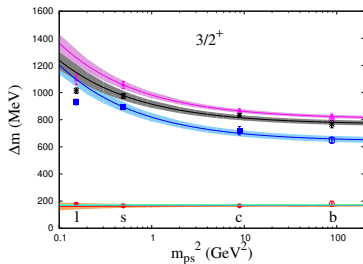
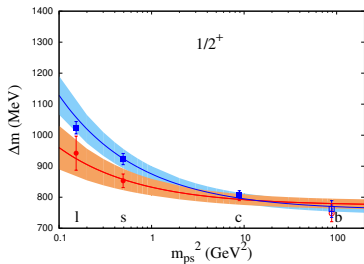
Spin-Orbit splittings in Ω like baryons



u and s \rightarrow Edwards, et. al., Phys. Rev. D **87**, 054506 (2013)

b \rightarrow S. Meinel, Phys. Rev. D **85**, 114510 (2012)

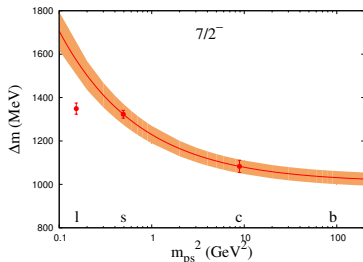
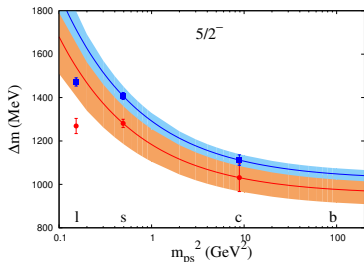
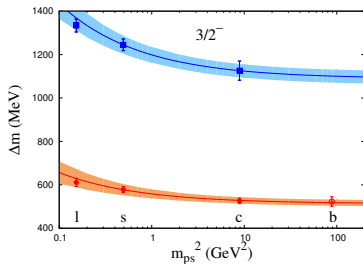
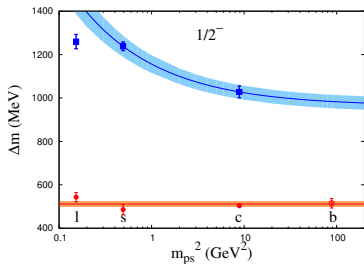
Quark mass dependence of Ω like baryons



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Summary and conclusions

- Non-perturbative calculation for excited state spectroscopy of Ω_{ccc} , Ω_{cc} and Ξ_{cc} .
- Non-relativistic spectrum pattern observed up to the second energy band.
- Identification of the spin and spatial structure of the states using the overlap factors.
- SO splittings : The degeneracy more or less satisfied for m_c .
- Energy splittings : Heavy quark inspired form gives good fit with m_b , m_c as well as m_s . For some, the fits even pass through m_l also.
- Extrapolations to bottom sector : $B_c^* - B_c = 80 \pm 8$ MeV and $\Omega_{ccb}^* = 8050 \pm 10$ MeV.
- No multi hadron operators being used : Further works required to see their effects.
- Singly charm baryons under investigation.