Overlap/Domain-wall reweighting

SAKA UNIVERSITY Live Locally, Grow Globally

Hidenori Fukaya (Osaka Univ.), S. Aoki, G. Cossu, S. Hashimoto, T. Kaneko, J. Noaki [JLQCD collaboration]



JLQCD collaboration

JLQCD (+TWQCD) collaboration 2006-2012

We have been simulating QCD with overlap quarks.





New project launched.

Simulations on bigger & finer lattices started.

Computers @KEK: SR11000 (2 TFLOPS) + BG/L (57 TFLOPS) \rightarrow SR16000 (55 TFLOPS) + BG/Q (1.2 PFLOPS) Lattice cut-off: 1.8 GeV \rightarrow 2.4, 3.6, 4.2 GeV Lattice size : 16³ x48 \rightarrow 32³ x64, 48³x96, 64³x128 (Physical size : 1.8 fm \rightarrow 2.6 fm \sim 4 fm)

Fermion action : overlap fermion \rightarrow DomainWall(Mobius) fermion







New project launched.

Simulations on bigger & finer lattices started.

Computers @KEK: SR11000 (2 TFLOPS) + BG/L (57 TFLOPS) \rightarrow SR16000 (55 TFLOPS) + BG/Q (1.2 PFLOPS) Lattice cut-off: 1.8 GeV \rightarrow 2.4, 3.6, 4.2 GeV Lattice size: 16³ x48 \rightarrow 32³ x64, 48³x96, 64³x128 (Physical size: 1.8 fm \rightarrow 2.6 fm \sim 4 fm)

Fermion action : overlap fermion → DomainWall(Mobius) fermion

Our goal = high precision of (B)SM calculations (in particular, D & B mesons)



THE THE PARTY





Theoretically, they are just different expressions (approximations) of the same action:

$$D_{ov} = \frac{m_0 + m}{2} + \frac{m_0 - m}{2} \gamma_5 \operatorname{sgn} H_T$$

$$D_{DW} = \frac{m_0 + m}{2} + \frac{m_0 - m}{2} \gamma_5 \frac{(1 + H_T)^n - (1 - H_T)^n}{(1 + H_T)^n - (1 - H_T)^n}$$

In this talk, let me *define*

Overlap : 10⁻⁸ precision of chiral symmetry (m_{res} ~ 10 eV) Domain-wall : 10⁻³ or less. (m_{res} ~ 1 MeV) [cf. non-chiral D : m_{res} ~ 1 GeV.]

Numerical differences

Domain-wall : good for HMC.

- + cheaper numerical cost, topology tunnelings.
- chiral sym. violation, eigenvalues need
 5D eigen vectors
- Overlap : good for Measurements.
- + clean analysis with <u>exact chiral symmetry</u>. recycling eigen values/vectors (for different m)
- high numerical cost, topology tunnelings difficult (we fixed the topology in our previous works.)

Overlap/Domain-wall reweighting

Let's use both fermions with reweighting. HMC with domain-wall fermions + $\langle O \rangle_{\text{overlap}} = \left\langle O \frac{\det D_{ov}^{N_f}}{\det D_{DW}^{N_f}} \right\rangle_{\text{DW}}$ Measurements with overlap fermions Ishikawa, Algorithms & machines, Monday

Exact chiral symmetry, topology changes with reasonable numericial cost.

Overlap/Domain-wall reweighting

Our goal = to find the optimal implementations for $D_{DW} \& D_{ov}$.

$$D_{ov/DW} = \frac{m_0}{2} \left(1 + \gamma_5 \mathrm{sgn}H\right)$$

Approximation for sgn function

Kernel operator

- 1. D_{DW} with a reasonable numerical HMC cost, and marginal chiral condition.
- 2. $D_{ov}\,$ with a good overlap in configurations generated by D_{DW} and good chirality

Contents

1. Introduction
2. Overlap vs. Domain-wall
3. Preliminary lattice results
4. Summary



Domain-wall fermion [Kaplan 1992, Shamir 1993 …] and its variations, [Borici, Chiu, Brower …] summarized by Edwards-Heller (2000) :

$$D_{GDW}^{5} = \begin{pmatrix} (D_{-}^{1})^{-1}D_{+}^{1} & -P_{-} & 0 & \cdots & 0 & mP_{+} \\ -P_{+} & (D_{-}^{2})^{-1}D_{+}^{2} & -P_{-} & 0 & \cdots & 0 \\ 0 & -P_{+} & (D_{-}^{3})^{-1}D_{+}^{3} & -P_{-} & \ddots & \vdots \\ \vdots & 0 & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & \ddots & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & \ddots & -P_{+} & (D_{-}^{L_{s}-1})^{-1}D_{+}^{L_{s-1}} & -P_{-} \\ mP_{-} & 0 & \cdots & 0 & -P_{+} & (D_{-}^{L_{s}})^{-1}D_{+}^{L_{s}} \end{pmatrix}$$
$$D_{+}^{s} = 1 + b_{s}D_{W}(-M_{0}), D_{-}^{s} = 1 - c_{s}D_{W}(-M_{0}); P_{\pm} = (1 \pm \gamma_{5})/2$$



$4D \text{ expression (with Pauli-Villars)} \rightarrow \begin{pmatrix} D & C \\ B & A \end{pmatrix} = \begin{pmatrix} 1 & CA^{-1} \\ 0 & 1 \end{pmatrix} \begin{pmatrix} D - CA^{-1}B(\equiv S_{\chi}) & 0 \\ 0 & A \end{pmatrix} \begin{pmatrix} 1 & 0 \\ A^{-1}B & 1 \end{pmatrix}$

Note: det A = 1

$$D_{DW}{}^{4D} = S_{\chi}^{-1}(m=1)S_{\chi}(m) = \frac{1+m}{2} - \frac{1-m}{2}\gamma_5 \frac{T_1^{-1}T_2^{-1}\cdots T_{L_s}^{-1} - 1}{T_1^{-1}T_2^{-1}\cdots T_{L_s}^{-1} + 1}$$

Approximation for sgn function Kernel operator





OSAKA UNIVERSIT

Overlap fermion [Neuberger 1998]

Exact treatment of the low-modes of the kernel:

allows us 10^{-8} chirality with Ls=O(10).

→ Exact chiral symmetry [L["]uscher 1998]

through the Ginsparg-Wilson relation [1982] :

$$D_{ov}(0)\gamma_5 + \gamma_5 D_{ov}(0) = D_{ov}(0)\gamma_5 D_{ov}(0)/m_0.$$



Overlap/Domain-wall reweighting

$$(D_{DW}^{4D})^{-1}D_{ov} = 1 + \sum_{\lambda_i < \lambda_{th}} ((D_{DW}^{4D})^{-1}D_{ov} - 1)|\lambda_i\rangle\langle\lambda_i|$$

Only sub-volume of matrix contributes.



Overlap/Domain-wall reweighting

$$(D_{DW}^{4D})^{-1}D_{ov} = 1 + \sum_{\lambda_i < \lambda_{th}} ((D_{DW}^{4D})^{-1}D_{ov} - 1)|\lambda_i\rangle\langle\lambda_i|$$

Only sub-volume of matrix contributes.

Let's reweight them !

$$\langle O \rangle_{\text{overlap}} = \left\langle O \frac{\det D_{ov}^{N_f}}{\det D_{DW}^{N_f}} \right\rangle_{\text{DW}}$$

Domain-wall fermion for sea quarks Our choice [T. Kaneko Tue (Chiral), S. Hashimoto, poster] Kernel = scaled Shamir Kernel: $2H_T = \gamma_5 \frac{2D_W}{2 + D_W}$

Sgn function = Tanh:

sgn_{tanh}
$$(2H_T) = \frac{(1+2H_T)^{L_s} - (1-2H_T)^{L_s}}{(1+2H_T)^{L_s} + (1-2H_T)^{L_s}} = \tanh(L_s \tanh^{-1}(2H_T))$$



Ls =

 m_{res} <0.5MeV, Chiral symmetry ~ 10⁻³

Simulation parameters Lattice size : 16³x32(x12) [test runs] Symanzik gauge action with $\beta = 4.17$ 3 steps of stout smearing 2+1 DWF m_{ud} = 0.7 m_s , $m_s \sim physical point$ 1/a ~ 2.4 GeV, L ~ 1.3 fm (L=32(2.6fm), L=48(3.9fm) lattices running.) [Simulated w/ Iroiro++ code, G. Cossu, poster] Faster than conventional DW (x4)

while reducing m_{res} to 1/10.

Domain-wall fermion for sea quarks

Topology tunnelings are active. mud0.012 ms0.030



Overlap fermion for valence quarks

$$D_{ov} = \frac{1}{2} \underbrace{\sum_{\lambda_i < \lambda_{th}} \left(1 + \gamma_5 \operatorname{sgn} \lambda_i\right) |\lambda_i\rangle \langle \lambda_i| + D_{DW}^{4D} \left(1 - \sum_{\lambda_i < \lambda_{th}} |\lambda_i\rangle \langle \lambda_i|\right)}_{\text{Exact low modes}}, \quad \text{High modes}$$

We try different

- 1. sgn functions : Zolotarev or Tanh for D^{4D}_{DW}
- 2. value of threshold λ_{th} for exact treatment of low-modes of H_T,
- 3. Ls (5th direction).

Chiral symmetry violation + & - eigenpairs of $H_{\alpha\nu} = \gamma_5 D_{\alpha\nu}$: $H_{ov}|\lambda\rangle = \lambda|\lambda\rangle$ $\rightarrow H_{ov}\Gamma_5|\lambda\rangle = -\lambda\Gamma_5|\lambda\rangle \quad {}^{(\Gamma_5 = \gamma_5(1 - aD_{ov}/2))}$ $\left| \frac{\lambda_+ + \lambda_-}{\lambda_+} \right| \simeq the precision of the GW relation.$ $(|\lambda_{ m zero}| - m \text{ as well})$





Reweighting factor

 $D_{ov} = \text{Zolotarev}, L_s = 12$, exact lowmodes $(\lambda_{th} = 0.3),$ $D_{DW} = \text{Tanh with } L_s = 12, \text{ w/o lowmode precond.},$

$$R = \left(\frac{\det \gamma_5 D_{ov}(m_{ud})}{\det \gamma_5 D_{DW}(m_{ud})}\right)^2 \frac{\det \gamma_5 D_{ov}(m_s)}{\det \gamma_5 D_{DW}(m_s)}$$
$$= \prod_{i=1}^{100} \left(\frac{\lambda_i^{ov}(m_{ud})}{\lambda_i^{DW}(m_{ud})}\right)^2 \left(\frac{\lambda_i^{ov}(m_s)}{\lambda_i^{DW}(m_s)}\right) \times \text{high-mode part}$$

High mode part is stochastically estimated with 60 - 300 Gaussian noises.



Results (with 46 confs separated by 50 trj)



L=1.3fm, 1/a=2.4 GeV





Results (with 46 confs separated by 50 trj)



Mismatch in topological charge ?

Conf1400 [R=1.54(6)]



Mismatch in topological charge?

Conf1400 [R=1.54(6)]





Mismatch in topological charge ?

Conf1400 [R=1.54(6)]



Conf2000 [R=0.00009(8)]





 D_{ov} and D_{DW} are just different expressions (approximations) of the same operator:

$$D_{ov} = \frac{m_0 + m}{2} + \frac{m_0 - m}{2} \gamma_5 \text{sgn} H_T$$
$$D_{DW} = \frac{m_0 + m}{2} + \frac{m_0 - m}{2} \gamma_5 \frac{(1 + H_T)^n - (1 - H_T)^n}{(1 + H_T)^n - (1 - H_T)^n}$$
1. Reweighting using

 $D_{ov} = \text{Zolotarev}, L_s = 12, \text{ exact lowmodes } (\lambda_{th} = 0.3),$

 D_{DW} = Tanh with $L_s = 12$, w/o lowmode precond.,

gives R~O(1) in 80 % configurations.

2. But R is very small when D_{DW} misses topological zero-modes.



To do

- 1. Any way to avoid mismatches in topology ?
- 2. Larger volume
- 3. Smaller quark mass

Other directions :

- 4. Loosen chirality of sea quarks
- 5. Isospin breaking effects
- 6. mixed action

Backup slide 1

of the kernel low-modes

below 0.3, we need

$$Ls = 16 \rightarrow \sim 30$$
$$Ls = 32 \rightarrow \sim 300$$
$$Ls = 48 \rightarrow \sim 2400$$

Backup slide 2

Reweighting factor of Zolo/Tanh

$$\frac{D_{ov}^{Zolo}(\lambda_{th}=0.3)}{D_{ov}^{Tanh}(\lambda_{th}=0.3)} = 1 + ((D_{DW}^{Tanh})^{-1} D_{DW}^{Zolo} - 1)(1 - \sum_{\lambda_i < \lambda_{th}} |\lambda_i\rangle \langle)\lambda_i|$$

$$L=1.3 \text{fm}, 1/a=2.4 \text{ GeV}$$



Backup slide 3

Correlation with kernel eigenvalues ?

Eigenvalues of $2H_T$



33

Generalized 5d implementation

– After an unitary transformation,

$$D_{\chi}^{5} \equiv \begin{pmatrix} P_{-} - mP_{+} & -T_{1}^{-1} & 0 & \cdots & 0 \\ 0 & 1 & T_{2}^{-1} & 0 & \cdots & 0 \\ \vdots & 0 & 1 & T_{3}^{-1} & 0 & 0 \\ \vdots & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & \cdots & 0 & 1 & T_{L_{s}^{-1}}^{-1} \\ -T_{L_{s}}^{-1}(P_{+} - mP_{-}) & 0 & \cdots & \cdots & 0 & 1 \\ T_{s}^{-1} \equiv -(Q_{-}^{s})^{-1}Q_{+}^{-1}, & Q_{\pm}^{s} = (D_{-}^{s})^{-1}D_{+}^{s}P_{\mp} - P_{\pm} \end{pmatrix}$$

- Then, Schur complement $\begin{pmatrix}
D & C \\
B & A
\end{pmatrix} = \begin{pmatrix}
1 & CA^{-1} \\
0 & 1
\end{pmatrix}
\begin{pmatrix}
D - CA^{-1}B(\equiv S_{\chi}) & 0 \\
0 & A
\end{pmatrix}
\begin{pmatrix}
1 & 0 \\
A^{-1}B & 1
\end{pmatrix}$ Note: $\overset{34}{\det} A = 1$

Generalized 5d implementation

$$S_{GDW} = \sum_{x} \bar{\psi} D_{GDW}^{5} \psi$$

Edwards-Heller (2000)

$$D_{GDW}^{5} \equiv \begin{pmatrix} (D_{-}^{1})^{-1}D_{+}^{1} & -P_{-} & 0 & \cdots & 0 & mP_{+} \\ -P_{+} & (D_{-}^{2})^{-1}D_{+}^{2} & -P_{-} & 0 & \cdots & 0 \\ 0 & -P_{+} & (D_{-}^{3})^{-1}D_{+}^{3} & -P_{-} & \ddots & \vdots \\ \vdots & 0 & \ddots & \ddots & \ddots & \vdots \\ 0 & \cdots & \ddots & -P_{+} & (D_{-}^{L_{s}-1})^{-1}D_{+}^{L_{s-1}} & -P_{-} \\ mP_{-} & 0 & \cdots & 0 & -P_{+} & (D_{-}^{L_{s}})^{-1}D_{+}^{L_{s}} \end{pmatrix}$$

 $D_{\pm}^{s} = 1 + b_{s}D_{W}(-M_{0}), D_{\pm}^{s} = 1 - c_{s}D_{W}(-M_{0}); P_{\pm} = (1 \pm \gamma_{5})/2$

4D effective operator

 $-\det A = 1$ doesn't contribute to path integ.

$$S_{\chi}(m) = -(1 + T_1^{-1}T_2^{-1} \cdots T_{L_s}^{-1})\gamma_5 \left[\frac{1+m}{2} - \frac{1-m}{2}\gamma_5 \frac{T_1^{-1}T_2^{-1} \cdots T_{L_s}^{-1} - 1}{T_1^{-1}T_2^{-1} \cdots T_{L_s}^{-1} + 1}\right]$$

- Combining with Pauli-Villars (m=1), $D^{(4)} \equiv S_{\chi}^{-1}(m=1)S_{\chi}(m) = \frac{1+m}{2} - \frac{1-m}{2}\gamma_5 \frac{T_1^{-1}T_2^{-1}\cdots T_{L_s}^{-1} - 1}{T_1^{-1}T_2^{-1}\cdots T_r^{-1} + 1}$
 - Looks similar to overlap

$$\operatorname{sgn}^{(approx)} = \frac{1 - \prod_{s} T_{s}}{1 + \prod_{s} T_{s}}$$

may approximate the sign function