

# Computation of the strong coupling in $N_f = 4$ QCD

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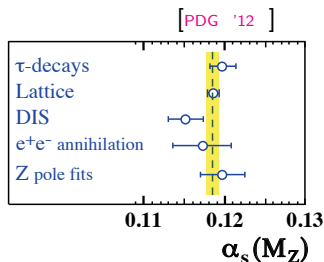
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## Motivation

- Error coming from the theory is dominating
- Lattice QCD: *ab initio* determination of QCD parameters
- Full control over systematics
  - non-perturbative renormalization
  - chiral extrapolation
  - continuum extrapolation



| $N_f$ | $\Lambda_{\overline{MS}}$ | Theory  |
|-------|---------------------------|---|
| 0     | 238(19)                   | LGT,  [Della Morte'05] |
| 2     | 310(20)                   | LGT,  [Fritzsch'12]    |
| 5     | 160(11)                   | NNLO PT, fits to PDFs [Alekhin'12]  |
| 5     | 198(16)                   | NNLO PT, fits to PDFs [Martin'09]   |
| 5     | 275(57)                   | 4-loop PT at $M_Z$  |

## Strategy for the computation of $\Lambda_{\overline{MS}}^{(5)}$ in physical units

$$\frac{\Lambda_{\overline{MS}}^{(5)}}{f_K} = \frac{1}{f_K L_{\max}} \times \frac{L_{\max}}{L_k} \Lambda_{SF}^{(4)} L_k \times \frac{\Lambda_{\overline{MS}}^{(4)}}{\Lambda_{SF}^{(4)}} \times \frac{\Lambda_{\overline{MS}}^{(5)}}{\Lambda_{\overline{MS}}^{(4)}}$$

- Scale setting: chiral and continuum extrapolation needed
- Intermediate scheme: connecting low and high energies  
1-loop PT  $\Lambda_{\overline{MS}}^{(4)}/\Lambda_{SF}^{(4)} = 2.9065$  [Sint'96, Sommer'11]
- Decoupling across the b-threshold [Chetyrkin'00]

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# The step scaling function (SSF)

$$\sigma(s, \bar{g}^2(L))$$

## Problem of the lattice formulation

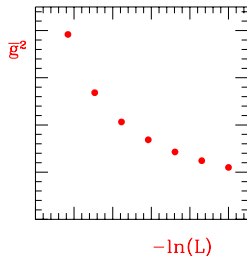
$$\begin{array}{ccccccc} L & \gg & \frac{1}{0.2\text{GeV}} & \gg & \frac{1}{\mu} \sim \frac{1}{10\text{GeV}} & \gg & a \\ \uparrow & & \uparrow & & \downarrow & & \uparrow \\ \text{box size} & & \text{conf. scale, } m_\pi & & & & \text{spacing} \\ & & & & L/a \gg 50 & & \end{array}$$

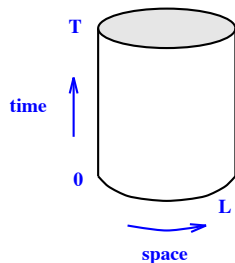
•  Solution:  $L = 1/\mu$  finite size scaling [Lüscher, Weisz, Wolff 1991]

• Integrated beta function - step scaling function  
 $\sigma(s, u) = \bar{g}^2(sL)$ , with  $u = \bar{g}^2(L)$  [ $s=2$ ]

• determines the non-perturbative running:

$$\begin{aligned} u_0 &= \bar{g}^2(L_{\max}) \\ &\downarrow \\ \sigma(2, u_{k+1}) &= u_k \\ &\downarrow \\ u_k &= \bar{g}^2(2^{-k} L_{\max}) \end{aligned}$$





[Lüscher, Narayanan, Weisz, Wolff 1992]

$$e^{-\Gamma} = \int \mathcal{D}U \mathcal{D}\bar{\psi} \mathcal{D}\psi e^{-S[U, \bar{\psi}, \psi]}$$

$$U(x, k)|_{x_0=0} = \exp\{aC_k(\eta)\},$$

$$U(x, k)|_{x_0=T} = \exp\{aC'_k(\eta)\}$$

$$\Gamma' = \left. \frac{\partial \Gamma}{\partial \eta} \right|_{\eta=0} = \frac{k}{\bar{g}^2(L)} \propto \langle F_{0k}^8 |_{\text{boundary}} \rangle$$

Advantages of SF approach:

- Infrared safe: one can simulate massless quarks ( $\lambda_{min}^2(\gamma_5 D) \approx 1/L^2$ )
- improvement coefficients can be computed non-perturbatively
- $\frac{\partial \Gamma}{\partial \eta}$  easily computed in Monte Carlo simulations  $\langle dS/d\eta \rangle$

- On the lattice:

additional dependence on the resolution  $a/L$

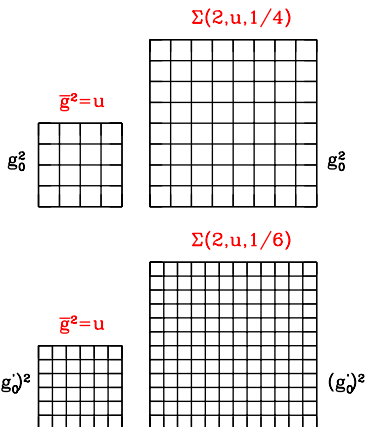
→  $g_0$  fixed,  $L/a$  doubled:

$$\begin{aligned} \bar{g}^2(L) &= u, & \bar{g}^2(sL) &= u', \\ \Sigma(s, u, a/L) &= u' \end{aligned}$$

→ continuum limit:

$$\Sigma(s, u, a/L) = \sigma(s, u) + O(a/L)$$

- In our work:  $s = 2$
- $m = 0$  (PCAC mass defined in  $(L/a)^4$  lattice)



## Lattice setup

- $N_f = 4$   $O(a)$  improved *massless Wilson fermions*

## New code for SF simulations

- $N_f = 0, 2, 4$   $\tau$ -language *APE*  $\rightarrow$  flexible *PC code*
- *SF-MP-HMC*: Hasenbusch preconditioning,  $N_{PF} = 2$
- $L \equiv T \rightarrow$  CG solver, good scalability for SF lat. sizes

## Set of $N_f = 4$ SF ensembles

- High statistics:  $O(160000)$  MDU trajectories
  - $L/a = 6 \rightarrow 2L/a = 12$  ( 9 couplings)
  - $L/a = 8 \rightarrow 2L/a = 16$  (10 couplings)
  - $L/a = 12 \rightarrow 2L/a = 24$  ( 1 coupling - *check of the cutoff effects*)
- Goal: to reduce the errors from [Tekin et. al 2010] by a factor of 2




## Lattice setup

- $N_f = 4$   $O(a)$  improved *massless* Wilson fermions
  - Achieved by tuning  $\kappa_c(\beta, a/L)$ , s.t.  $m_{PCAC}(L) = 0$
  - Tuning criteria:  $|\text{syst. err. from } m_{PCAC}| < \text{aimed precision of } \sigma(\bar{g}^2(L))$

## Lattice setup

- $N_f = 4$   $O(a)$  improved *massless Wilson fermions*

### $O(a)$ improvement

- NP determined  $c_{SW}(g_0)$  for  $N_f = 4$  theory [Tekin, Sommer, Wolff, 2009]
- Boundary  $O(a)$ -terms, e.g.  $c_t(g_0)a^4 \sum_x F_{0k}F_{0k}$  at  $x_0 = 0$  and  $x_0 = T$   
2-loop computation by [Bode, Weisz, Wolff 1999]
- Remaining cutoff effects: [De Divitiis et al. 1993,  1993-1999]

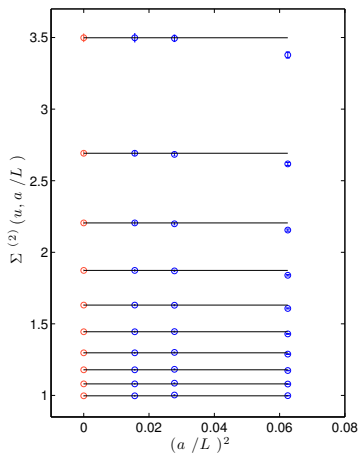
$$\delta(u, a/L) = \frac{\Sigma(u, a/L) - \sigma(u)}{\sigma(u)} = \delta_1(a/L)u + \delta_2(a/L)u^2 + \dots$$

$$\Sigma^{(2)}(u, a/L) = \frac{\Sigma(u, a/L)}{1 + \delta_1(a/L)u + \delta_2(a/L)u^2}$$

- Improved lattice SSF:  $\Sigma^{(2)}(u, a/L) = \sigma(u) + O(a^2)$

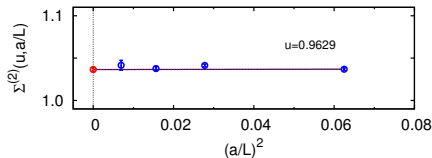
## Lattice SSF for $N_f = 4$

- $L/a = 4$  O(50000) MDU [Tekin '10]
- $L/a = 6, 8$  O(200000) MDU [this work + Tekin '10, combined]

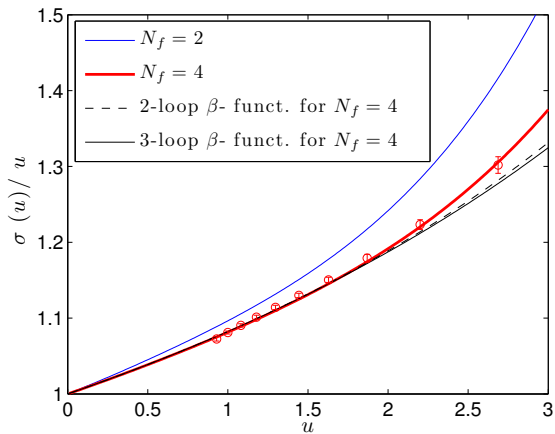


### Check of the cutoff effects:

- $L/a = 4, 6, 8, 12$ ,  
 $u = \bar{g}^2(L) = 0.9629$ .
- constant fit of  $L/a = 8$  data

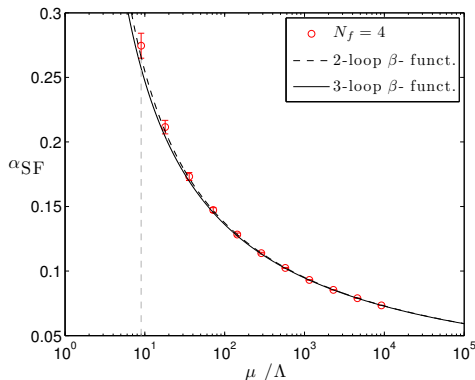


$$\Sigma^{(2)}(u, a/L) = \sigma(u), \quad L/a = 8$$



[  [Tekin '10 + this work; combined result] ]

# Non-perturbative running of $\alpha$ in the SF scheme $N_f = 4$



energies  $\mu$

0.7 GeV

.

.

.

700 GeV

$$\bar{g}^2(L_{\max}) = 3.45$$

$$\ln(\Lambda L_{\max}) = -2.195(80) \quad \Lambda \equiv \Lambda_{SF}$$

[old. res.  $\ln(\Lambda L_{\max}) = -2.294(153)$  [Tekin '10] corrected errorbars]

## Strategy for the computation of $\Lambda_{\overline{MS}}^{(5)}$ in physical units

Still missing!

$$\frac{\Lambda_{\overline{MS}}^{(5)}}{f_K} = \boxed{\frac{1}{f_K L_{\max}}} \times \frac{L_{\max}}{L_k} \Lambda_{SF}^{(4)} L_k \times \frac{\Lambda_{\overline{MS}}^{(4)}}{\Lambda_{SF}^{(4)}} \times \frac{\Lambda_{\overline{MS}}^{(5)}}{\Lambda_{\overline{MS}}^{(4)}}$$

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# Summary and outlook

## Summary

- Precise determination  $\Lambda$  from non-perturbative methods needed to resolve discrepancies of PT estimates
- Finite-size scaling + Schrödinger functional scheme
- SF-MP-HMC code: good scalability, CG solver, flexible for different PC clusters
- 2<sup>nd</sup> computation of the  $N_f = 4$  SF running coupling
- Lambda parameter in the units of  $L_{\max}$
- Ultimate goal:  $N_f = 4$  Lambda parameter in physical units

## Outlook

- Gradient flow coupling [See talks: P.Fritsch (Mon, 14.20h, 1C) & A. Ramos (Tue, 12h 6E) ]
- Implementation of advanced solvers
- Matching to the physical units  
precise setting of the scale for  $N_f = 2 + 1 + 1$  is needed

Thank you !



## SF in the lattice regularization: $O(a)$ improved Wilson fermions

- $S_W[U] = \beta \sum_P w(P) (1 - \frac{1}{3} \text{Re Tr } U_P), \quad \beta = 6/g_0^2$

$$w(P) = \begin{cases} 1 & U_P \text{ is the plaquette in the bulk,} \\ c_t(g_0) & U_P \text{ is the time-like touching } x_0 = 0 \text{ or } x_0 = T \text{ boundary,} \\ \frac{1}{2} c_s(g_0) & U_P \text{ is the space-like plaquette at } x_0 = 0, T. \end{cases}$$

- Fermionic fields also: periodic in space

$$\begin{aligned} \psi(x + L\hat{k}/a) &= e^{i\theta_k} \psi(x), \\ \bar{\psi}(x + L\hat{k}/a) &= e^{-i\theta_k} \bar{\psi}(x). \end{aligned}$$

fixed on the time boundaries

$$\begin{aligned} P_+ \psi(x)|_{x_0=0} &= \rho(x), & P_- \psi(x)|_{x_0=T} &= \rho'(x), \\ \bar{\psi}(x) P_-|_{x_0=0} &= \bar{\rho}(x), & \bar{\psi}(x) P_+|_{x_0=T} &= \bar{\rho}'(x), \end{aligned}$$

- $S_F = S_f^W + S_{SW} + S_b$

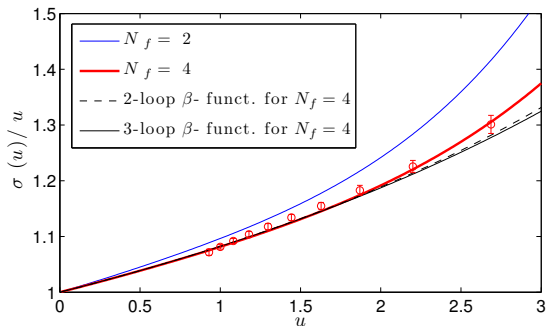
$$\begin{aligned} S_b[U, \psi, \bar{\psi}] &= \frac{a^4}{2} \sum_x \{ (\tilde{c}_s(g_0) - 1) [(\bar{\rho}(x) \gamma_k (\nabla_{\mathbf{k}}^* + \nabla_{\mathbf{k}}) \rho(x)) \\ &\quad - (\bar{\rho}'(x) \gamma_k (\nabla_{\mathbf{k}}^* + \nabla_{\mathbf{k}}) \rho'(x))] \\ &\quad + (\tilde{c}_t(g_0) - 1) [(\bar{\psi}(x) (P_- \nabla_0 + P_+ \overleftarrow{\nabla}_0^*) \psi(x))_{x_0=a} \\ &\quad - (\bar{\psi}(x) (P_- \nabla_0 + P_+ \overleftarrow{\nabla}_0^*) \psi(x))_{x_0=T-a}] \} \end{aligned}$$

- Additional improvement coefficients: on-shell improvement of the axial current [CA](#)

# $N_f = 4$ running coupling - Backup slides

| $u$     | $\sigma(u)$ | $\Delta(\sigma(u))$ |
|---------|-------------|---------------------|
| 0.93    | 0.998567    | 0.00164622          |
| 1       | 1.07943     | 0.00148896          |
| 1.08128 | 1.17405     | 0.00187277          |
| 1.1787  | 1.28861     | 0.00237175          |
| 1.29717 | 1.42987     | 0.00280331          |
| 1.44354 | 1.60788     | 0.00342647          |
| 1.62854 | 1.83945     | 0.00493976          |
| 1.87    | 2.1551      | 0.00761488          |
| 2.20033 | 2.61708     | 0.0105432           |
| 2.68704 | 3.37799     | 0.0213151           |

| $n$ | $u_n$      | $\ln(\Lambda L_{max})$ | $\ln(\Lambda L_{max})^{(2-loop)}$ |
|-----|------------|------------------------|-----------------------------------|
| 0   | 3.450      | -2.028                 | -2.117                            |
| 1   | 2.658(17)  | -2.076(21)             | -2.146                            |
| 2   | 2.177(17)  | -2.109(31)             | -2.166                            |
| 3   | 1.850(15)  | -2.132(38)             | -2.182                            |
| 4   | 1.612(13)  | -2.149(45)             | -2.193                            |
| 5   | 1.431(12)  | -2.163(52)             | -2.201                            |
| 6   | 1.287(11)  | -2.173(59)             | -2.208                            |
| 7   | 1.171(10)  | -2.182(66)             | -2.214                            |
| 8   | 1.0737(93) | -2.189(73)             | -2.218                            |
| 9   | 0.9921(86) | -2.195(80)             | -2.222                            |
| 10  | 0.9222(80) | -2.200(87)             | -2.225                            |



[ ALPHA Collaboration 2010, F.Tekin, *Corrected errorbars* ]