

Finite size effects in lattice RI-MOM

F. Di Renzo, M. Brambilla

University of Parma and INFN

LATTICE 2013

Mainz, 29-07-2013

RI-MOM (or its RI'-MOM variant) is one of the most popular renormalization schemes for Lattice QCD; being regulator independent, it can be effectively adopted in a lattice regularization. RI-MOM is defined in **infinite volume**. This is in principle a fundamental problem for the lattice, since any simulation is performed in a finite volume. From a practical point of view, one most often verifies a posteriori (by performing computations on different physical volumes) the expectation that renormalization constants, determined in the RI-MOM scheme at large momenta, should not be affected by significant finite size effects.

In the context of **Numerical Stochastic Perturbation Theory**, we have in recent years devised a novel method to explicitly **look and correct for finite size effects** (in a convenient window). This is an account of what we learnt, trying to put it in a perspective. We review the method, discussing how it can be applied in a non-perturbative formulation as well.

Agenda

- RI'-MOM: definitions and standard best practice
- NSPT for RI'-MOM renormalization constants and FINITE SIZE EFFECTS
A little history of what we understood
- Taming FINITE SIZE EFFECTS in NSPT
An account of our NSPT best practice
- Can we go the SAME WAY in the NON-PERTURATIVE approach?
A proposal for a “better best practice”

RI'-MOM: definitions and standard best practice

The celebrated RI'-MOM

$$\int dx \langle p | \psi(x) \Gamma \psi(x) | p \rangle = G_\Gamma(pa)$$

$$G_\Gamma(pa) \rightarrow \Gamma_\Gamma(pa) = S^{-1}(pa) G_\Gamma(pa) S^{-1}(pa)$$

$$O_\Gamma(pa) = \text{Tr} \left(\hat{P}_{O_\Gamma} \Gamma_\Gamma(pa) \right)$$

$$Z_{O_\Gamma}(\mu a, g(a)) Z_q^{-1}(\mu a, g(a)) O_\Gamma(pa) \Big|_{p^2=\mu^2} = 1$$

$$Z_q(\mu a, g(a)) = -i \frac{1}{12} \frac{\text{Tr}(\not{p} S^{-1}(pa))}{p^2} \Big|_{p^2=\mu^2}$$

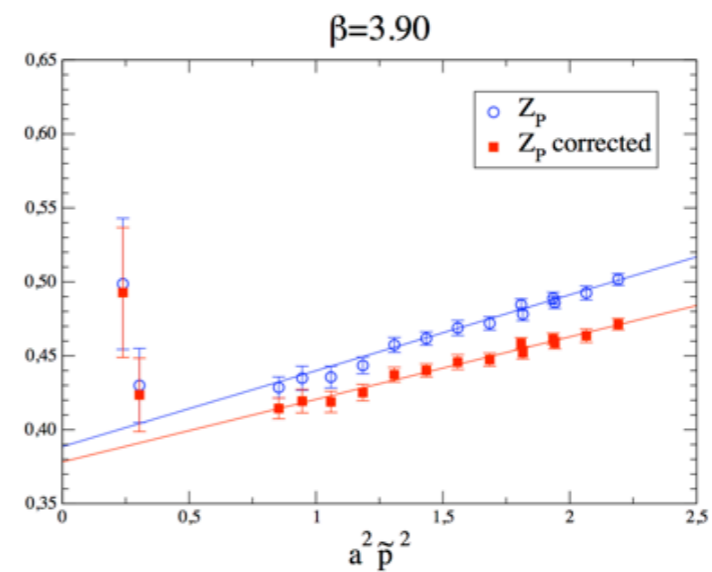
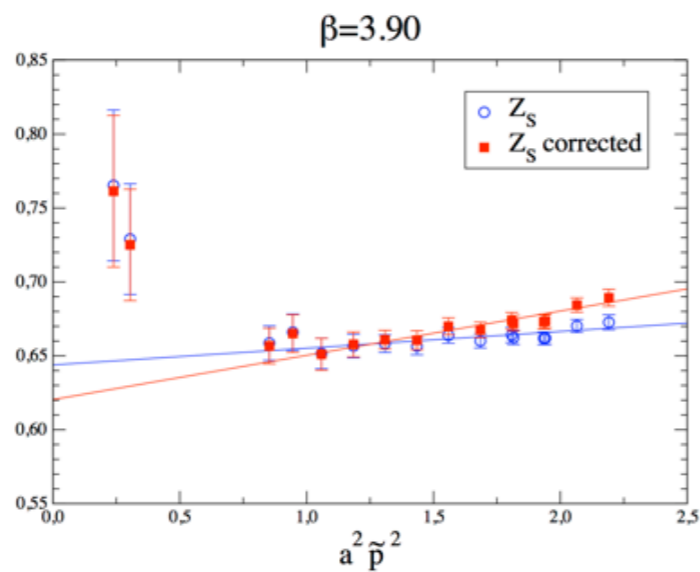
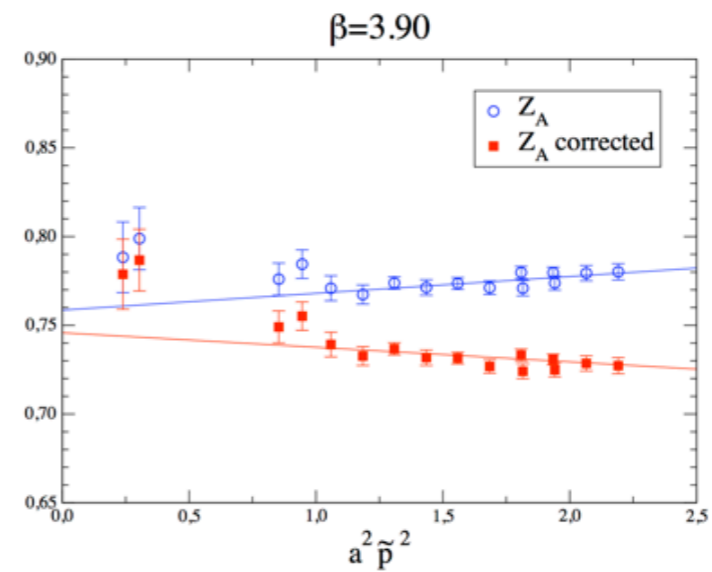
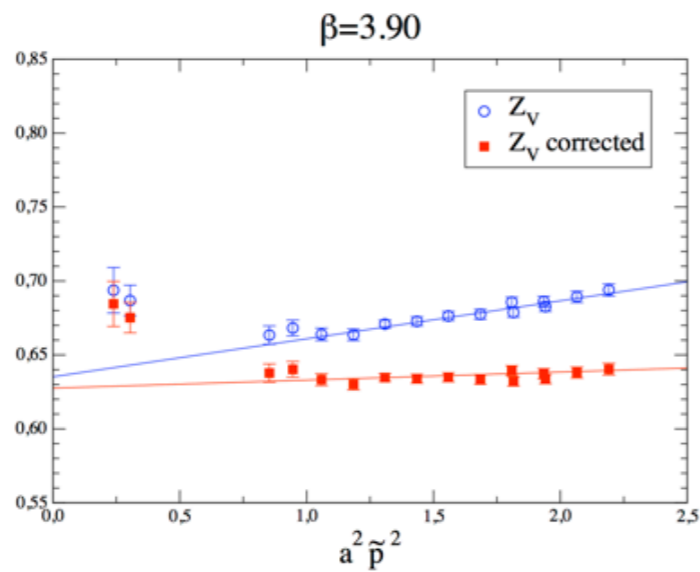
A few remarks:

- defined in the **chiral limit** (needs extrapolation)
- natural in **momentum space**

(still sources for inverting propagators are not always taken diagonal in momentum)

- defined in **INFINITE VOLUME**

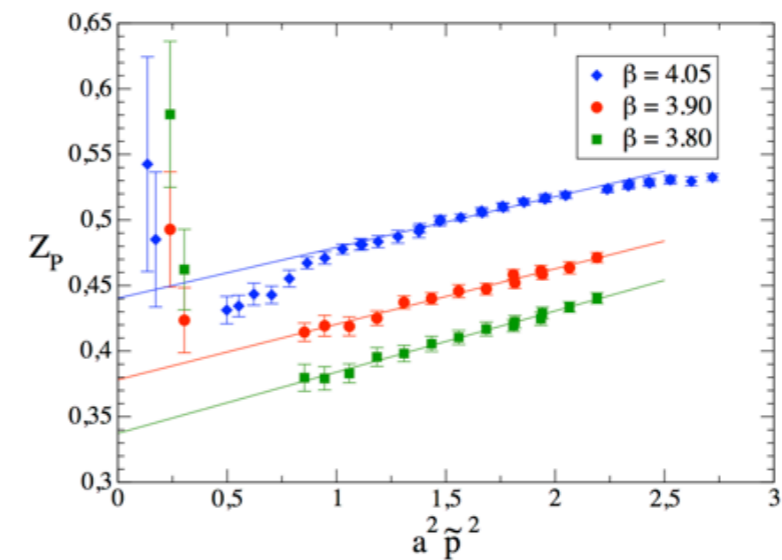
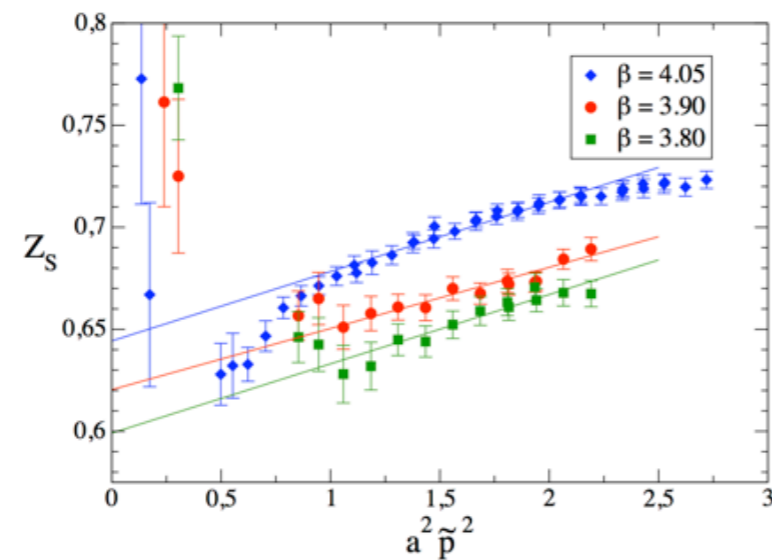
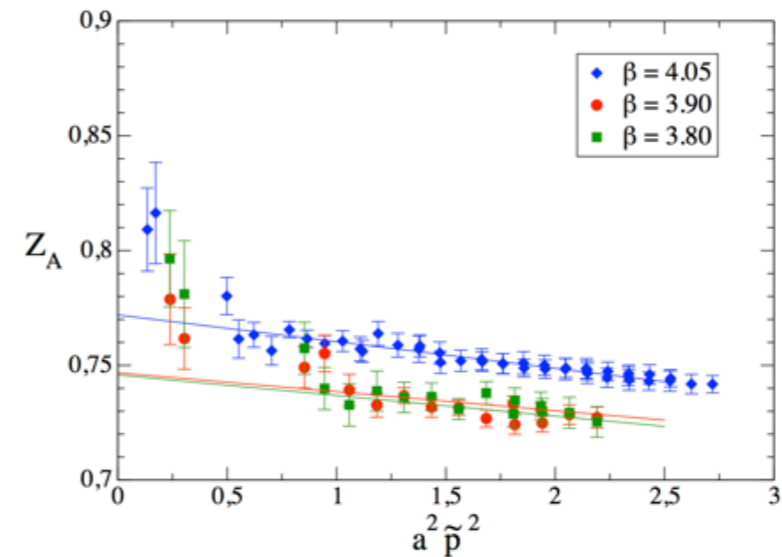
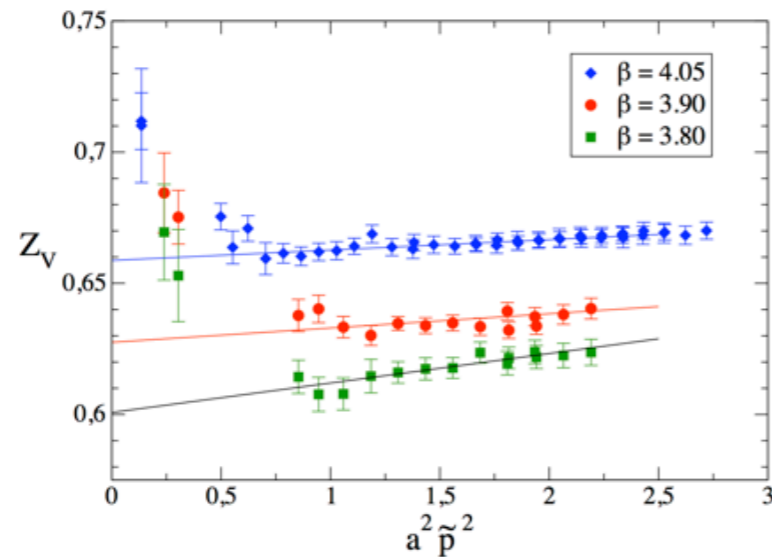
The standard best practice (from ETMC JHEP1008(2010)068)



In order to control finite lattice spacing, one-loop irrelevant contributions are (if possible, i.e. if known) subtracted

The standard best practice (2)

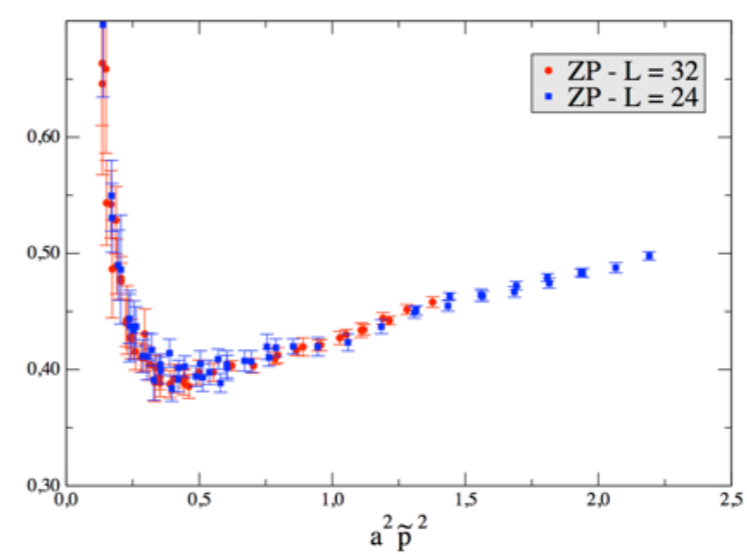
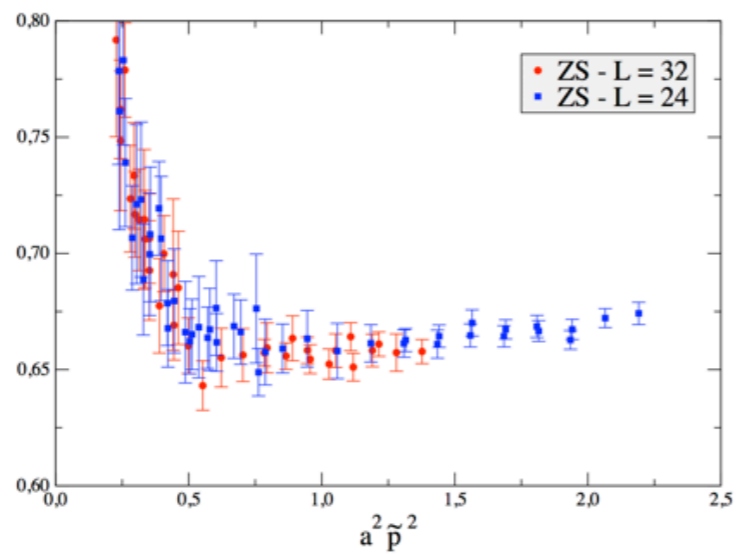
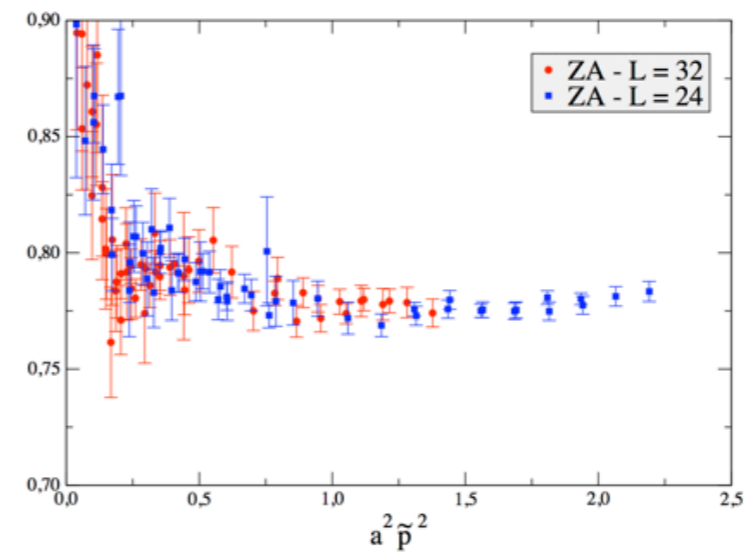
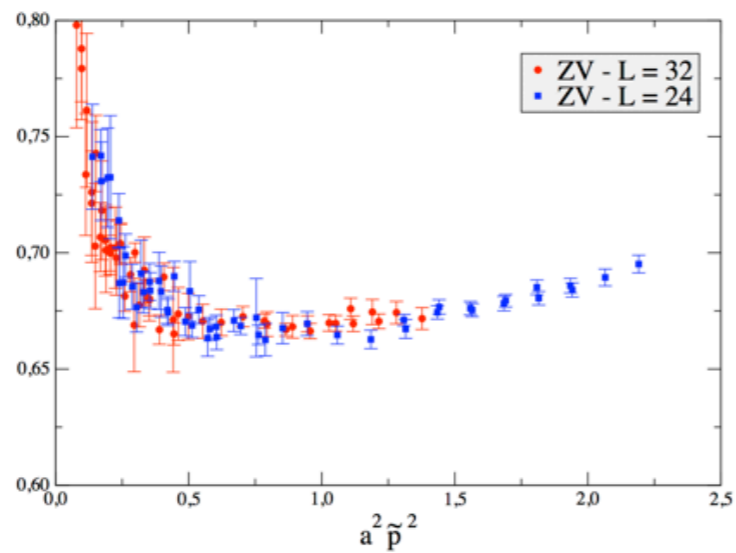
An extra extrapolation can be needed, i.e. an effective higher-order fit ...



... and here **FINITE SIZE EFFECTS** could pop in, since one is pinching the **IR** ...

The standard best practice (3)

A posteriori one looks for confirmation of not-significant **FINITE SIZE EFFECTS**



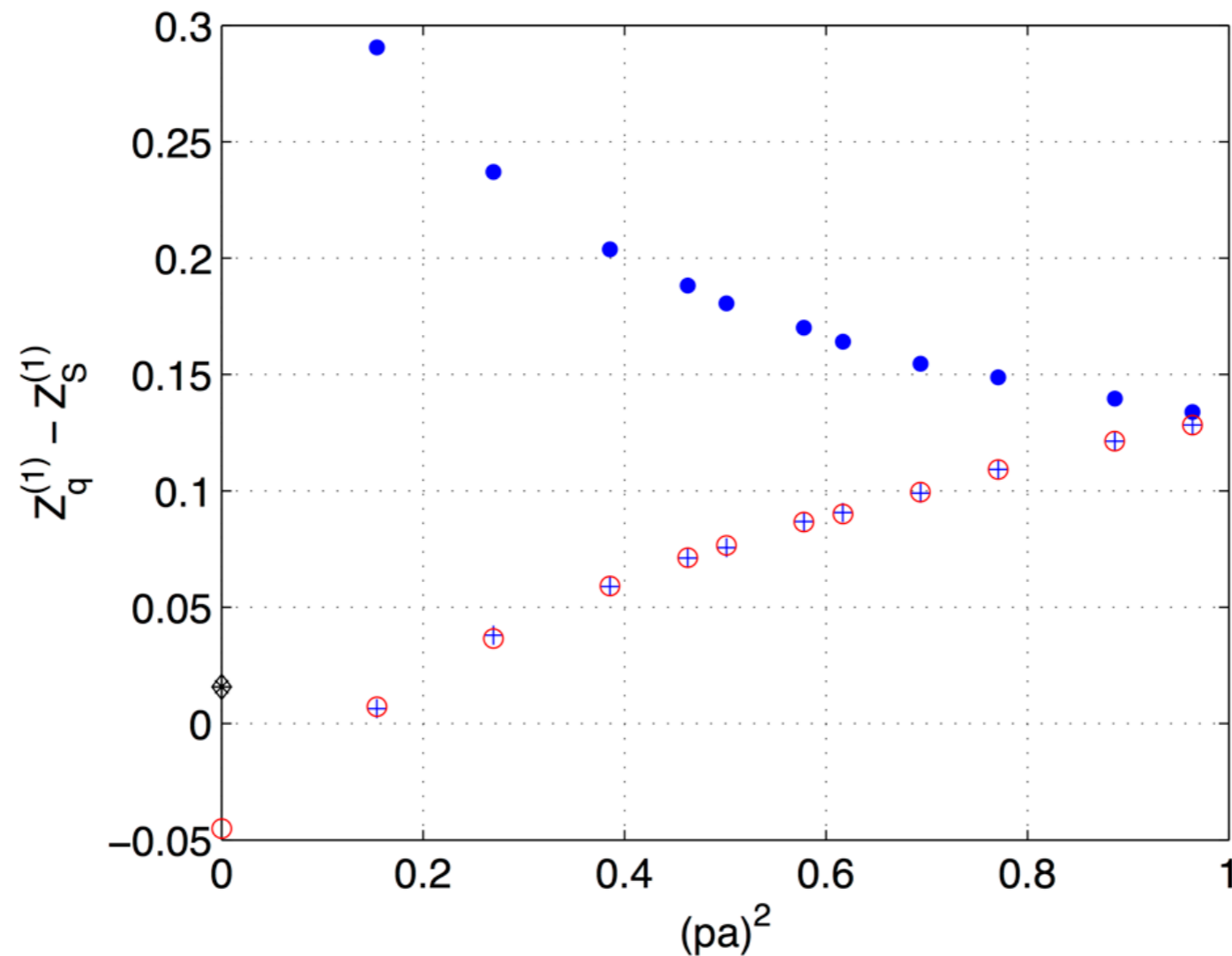
... our claim is that **there could be more information** buried into than looked at ...

NSPT for RI'-MOM renormalization constants
and FINITE SIZE EFFECTS

How FSE popped into NSPT computations (Parma group EPJC51(2007)645)

In order to fit the relevant part in $Z_{O_\Gamma}(\mu a, g(a))Z_q^{-1}(\mu a, g(a))O_\Gamma(pa)|_{p^2=\mu^2} = 1$

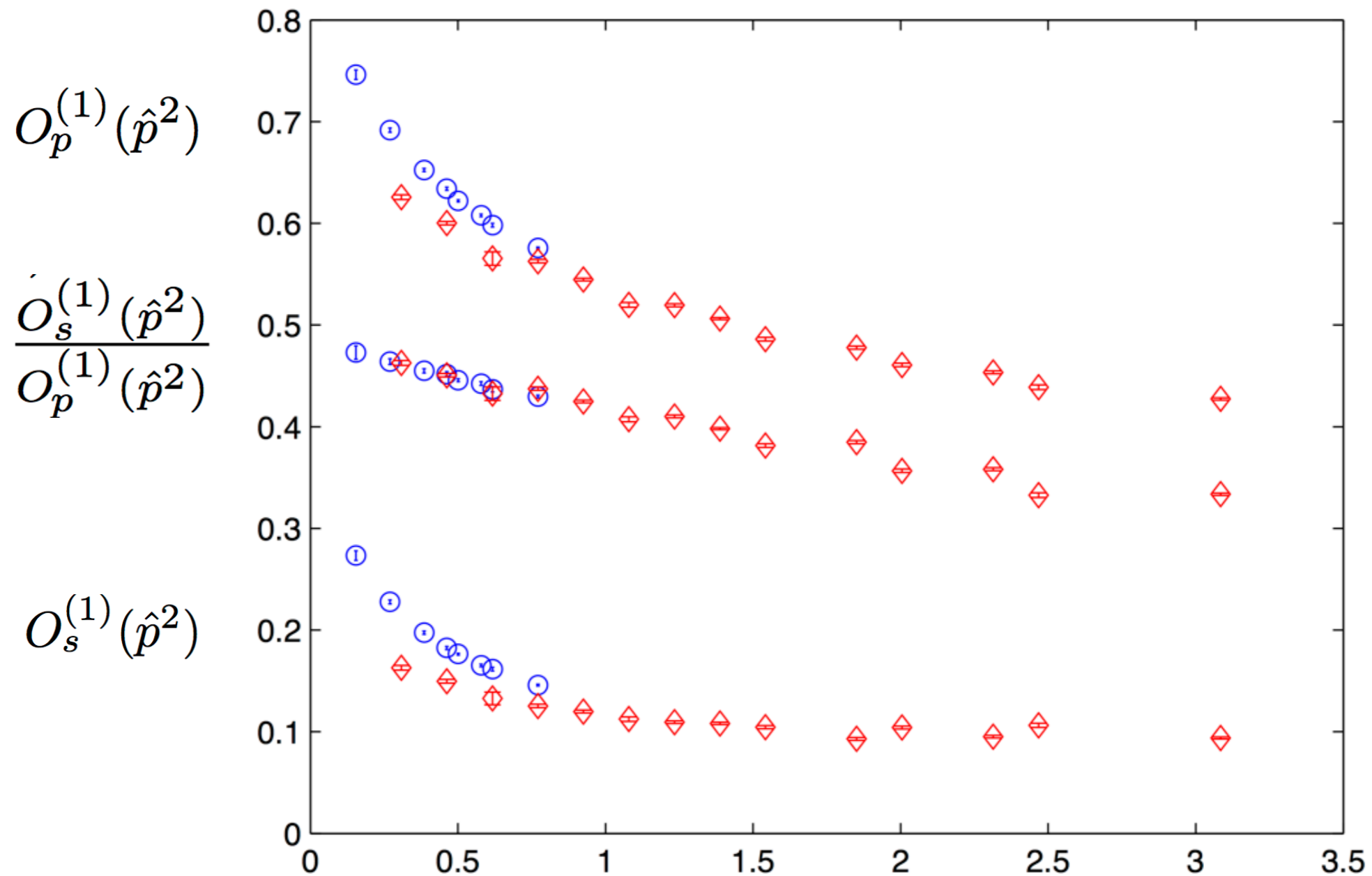
we need to subtract the logs (known) $z_q^{(1)} - z_s^{(1)} = O_s^{(1)}(\hat{p}^2) - \gamma_s^{(1)} \log(\hat{p}^2)$



... but we were failing!

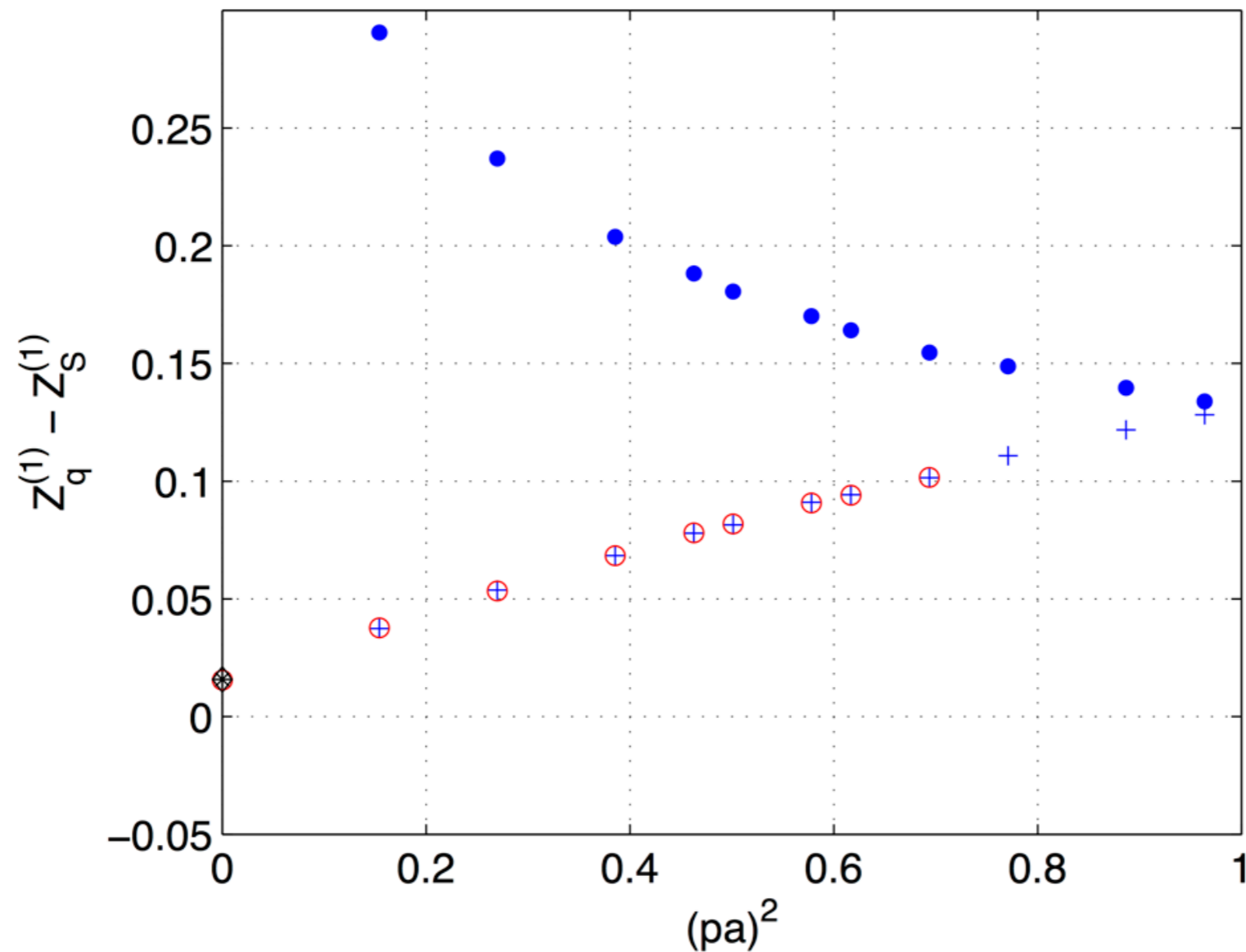
How FSE popped into NSPT computations (2)

By comparing (unsubtracted) results at different volumes, it was clear that it was due to **FINITE SIZE EFFECTS** ...



We had a first solution at that time (Parma group EPJC51(2007)645)

By computing what is left of a logarithm in a finite volume, we were at that time able to recover the correct 1 loop result ...



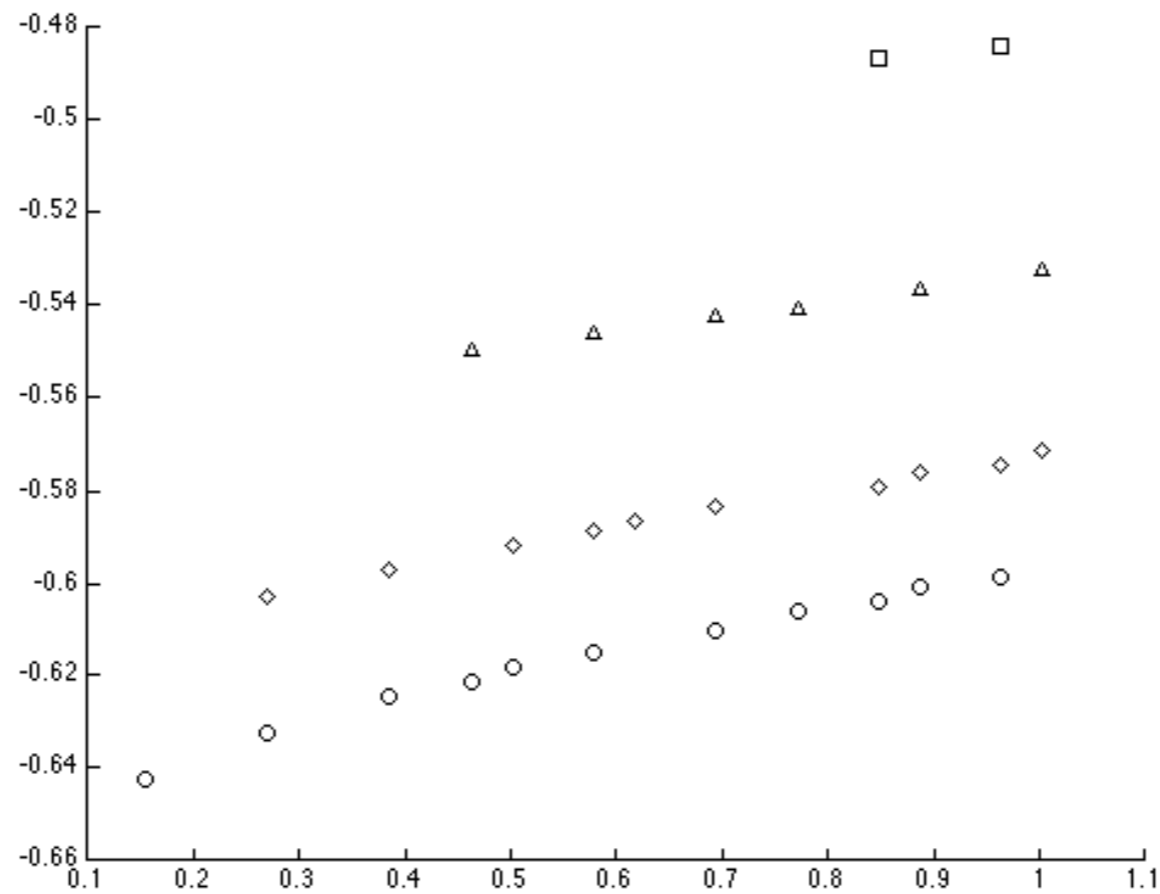
... but that was de facto only workin at 1 loop ...

Taming FINITE SIZE EFFECTS in NSPT

There is a big information built-in the hypercubic group!

Let's consider the **QUARK SELF ENERGY** $a\Gamma_2(\hat{p}, \hat{m}_{cr}, \beta^{-1}) = aS(\hat{p}, \hat{m}_{cr}, \beta^{-1})^{-1}$
 $= i\hat{p} + \hat{m}_W(\hat{p}) - \hat{\Sigma}(\hat{p}, \hat{m}_{cr}, \beta^{-1})$

$$\hat{\Sigma}(\hat{p}, \hat{m}_{cr}, \beta^{-1}) = \hat{\Sigma}_c(\hat{p}, \hat{m}_{cr}, \beta^{-1}) + \hat{\Sigma}_V(\hat{p}, \hat{m}_{cr}, \beta^{-1}) + \hat{\Sigma}_o(\hat{p}, \hat{m}_{cr}, \beta^{-1})$$



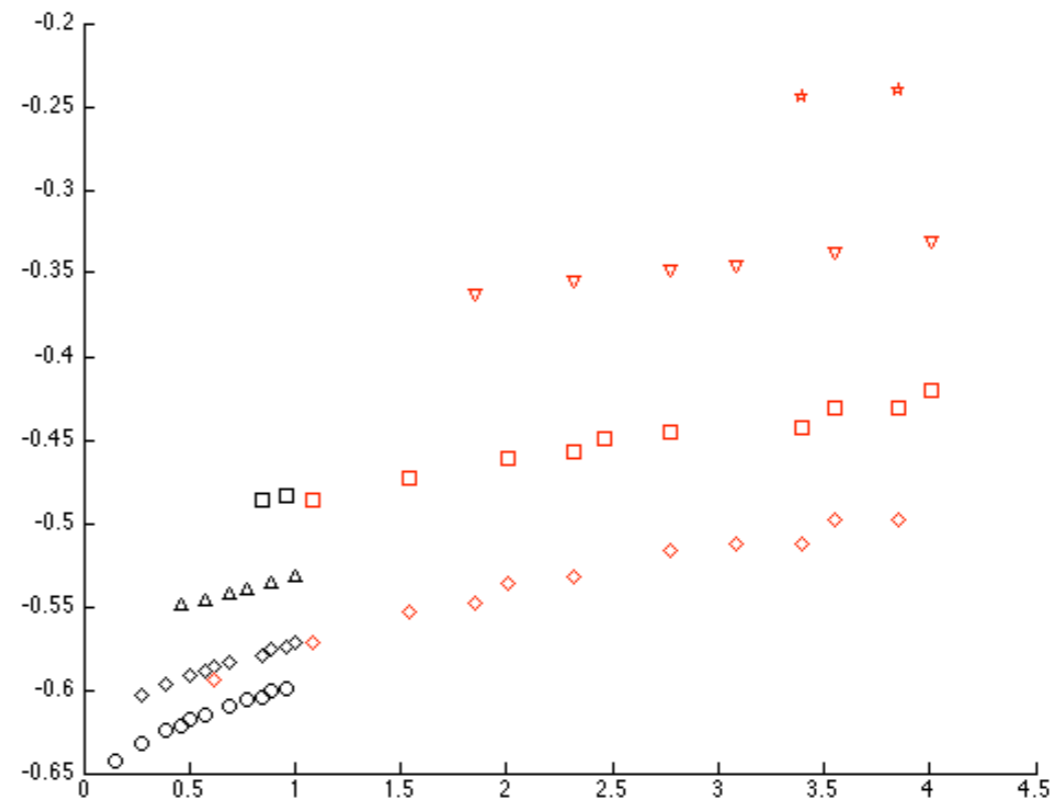
There is a clear sign of symmetry $\hat{\Sigma}_V = i \sum_u \gamma_\mu \hat{p}_\mu \left(\hat{\Sigma}_V^{(0)} + \hat{p}_\mu^2 \hat{\Sigma}_V^{(1)} + \hat{p}_\mu^4 \hat{\Sigma}_V^{(2)} + \dots \right)$

Notice that $\hat{\Sigma}_V^{(n)} = \alpha_1^{(n)} 1 + \alpha_2^{(n)} \sum_v \hat{p}_v^2 + \alpha_3^{(n)} \sum_v \hat{p}_v^4 / p^2 + \mathcal{O}(a^4)$

... and there was an important ingredient missing ...

Let's have a first look at 2 volumes

(32^4 and 16^4)



The real break-through was realizing that there was something missing

(first put at work in [Di Renzo, Ilgenfritz, Perlt, Schiller, Torrero NPB831\(2010\)262](#))

Expect rather

$$\hat{\Sigma}_V^{(n)}(\hat{p}, pL) = \hat{\Sigma}_V^{(n)}(\hat{p}, \infty) + \left(\hat{\Sigma}_V^{(n)}(\hat{p}, pL) - \hat{\Sigma}_V^{(n)}(\hat{p}, \infty) \right)$$
$$= \hat{\Sigma}_V^{(n)}(\hat{p}, \infty) + \Delta \hat{\Sigma}_V^{(n)}(\hat{p}, pL)$$

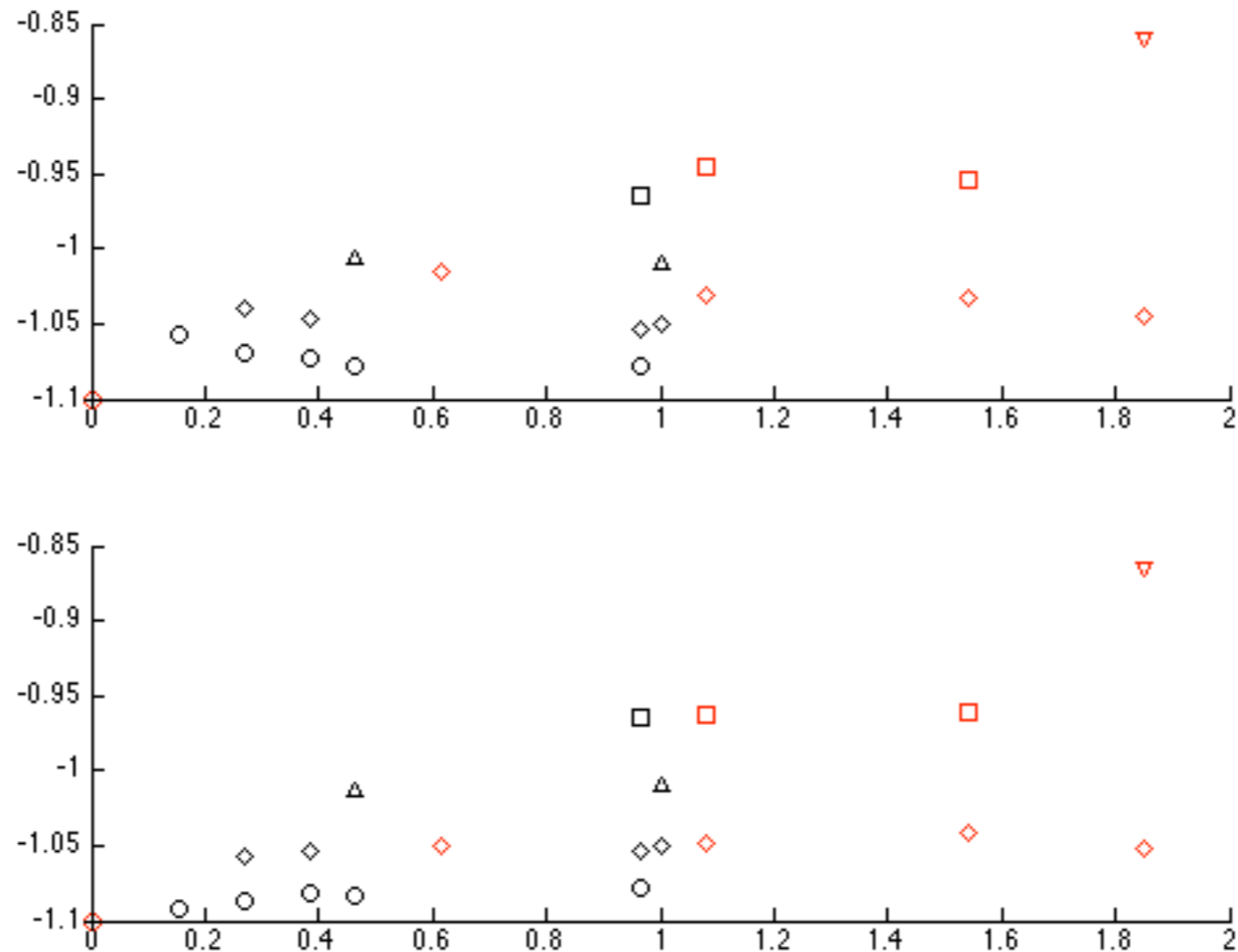
and a better fit is

$$\hat{\Sigma}_V^{(n)}(\hat{p}, pL) = \alpha_1^{(n)} 1 + \alpha_2^{(n)} \sum_{\mathbf{v}} \hat{p}_{\mathbf{v}}^2 + \alpha_3^{(n)} \sum_{\mathbf{v}} \hat{p}_{\mathbf{v}}^4 / p^2 +$$
$$+ \Delta \hat{\Sigma}_V^{(n)}(\hat{p}, pL) + \dots$$

What we really learnt of FSE from NSPT

Let's solve for a single renorm. constant $Z_{O_\Gamma}(\mu a, g(a)) Z_q^{-1}(\mu a, g(a)) O_\Gamma(pa) \Big|_{p^2=\mu^2} = 1$

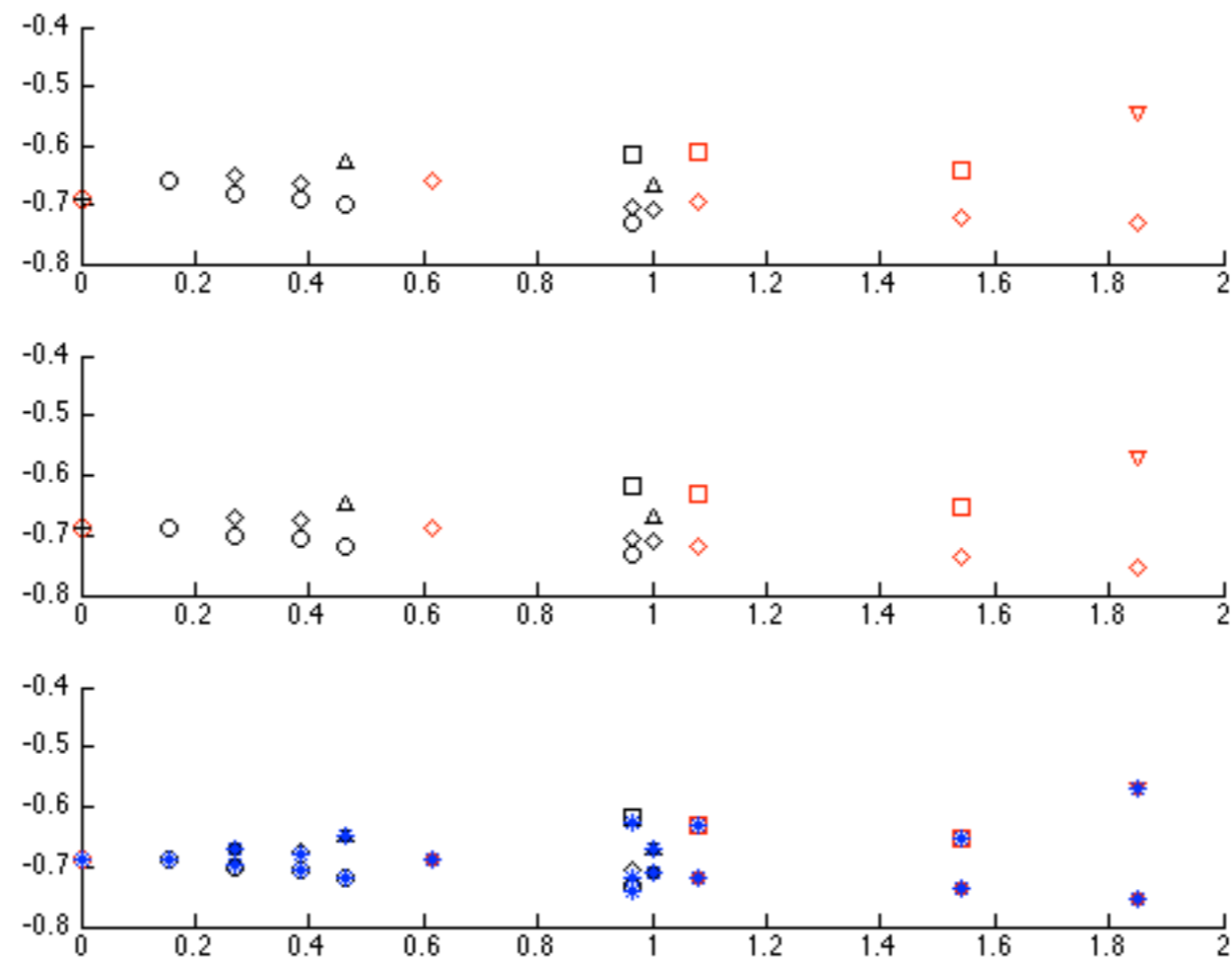
And let's plot 32^4 and 16^4 data together with and without the new fitting form (Z_p)
 (Brambilla, Di Renzo, Hasegawa to be published - see PoS LATTICE2012 242)



If $\Delta \hat{\Sigma}_V^{(n)}(\hat{p}, pL) \sim \Delta \hat{\Sigma}_V^{(n)}(pL)$ then $p_\mu L = \frac{2\pi n_\mu}{L} L = 2\pi n_\mu$

and we only have to care for one number for a given $\{n_1, n_2, n_3, n_4\}$

But if it works in NSPT (Brambilla, Di Renzo, Hasegawa to be published)...



... why don't we try to go the same way in a NP approach?

We need to:

- “really” work in momentum space
- always check the hyperbolic effects
- put different volumes together

We have no conclusions ...

... but we want to advertize work going on

- in assessing finite- a and finite- V effects in NSPT computations
- in implementing the same method of looking for finite size effects also in a NON PERTURBATIVE framework

... well aware that PT and NP will be very different, but confident that there is more information in RI'-MOM computations that we usually look into.

We can think of a “better best practice”

A possible approach can entail:

- a high order NSPT computation of renormalization constants
- a check (and possibly) correction for finite size effects for both NSPT and NP
- a better subtraction of irrelevant effects

Let me finally advertize the poster by Jakob Simeth