

Vortex liquid in superconducting vacuum of QCD induced by strong magnetic field

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Numerical part based on: arXiv:1104.3767, arXiv:1301.6590 + ongoing work

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Conventional superconductivity

- ① Zero electrical resistance
- ② An enemy of magnetic field
 - Meissner effect (expulsion of a weak magnetic field from a superconductor)
 - Stronger magnetic field penetrates into superconductor in terms of Abrikosov vortices (type-II superconductor)
 - Strong enough magnetic field kills superconductivity

The claim:

In a background of the strong magnetic field the vacuum of QCD becomes a superconductor(due to the condensation of charged ρ -mesons).

M.Chernodub, arXiv:1008.1055, arXiv:1101.0117

Key players: ρ -mesons and vacuum

ρ -mesons:

- electrically charged($\rho^\pm : q = \pm e$) and neutral($\rho^{(0)} : q = 0$) particles
- vector particles($s = 1$)
- Quark content: $\rho^+ : u\bar{d}$, $\rho^- : d\bar{u}$, $\rho^{(0)} : \frac{u\bar{u}-d\bar{d}}{\sqrt{2}}$
- mass: $m = 775.5\text{MeV}$
- lifetime: $c\tau = 1.35\text{fm}$

Vacuum:

- QCD+QED, zero temperature and density

Superconductivity of QCD in strong magnetic field

- 1 Spontaneously emerges above the critical value of magnetic field:

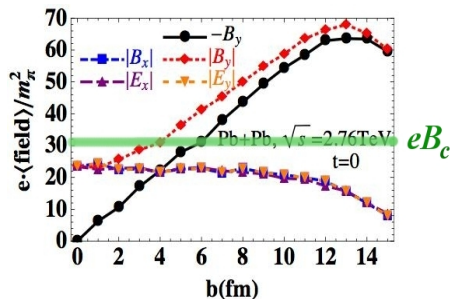
$$B_c \approx 10^{16} \text{ Tesla}$$

$$eB_c \approx m_\rho^2 \approx 0.6 \text{ GeV}^2$$

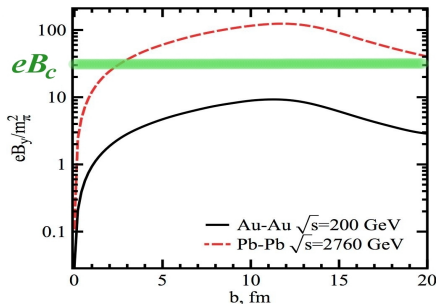
- 2 No Meissner effect (but vortices are formed)
- 3 Zero resistance along the the magnetic field
- 4 Insulator in other (perpendicular) directions

Strong magnetic fields

Strong magnetic fields ($eB_c = m_\rho^2 \approx 31m_\pi^2 \approx 0.6\text{GeV}^2$):



W. T. Deng and X. G. Huang,
Phys. Rev. C85 (2012) 044907,
arXiv:1201.5108



A. Bzdak and V. Skokov,
Phys.Lett. B710 (2012) 171,
arXiv:1111.1949

Approaches to the problem

- ① General ideas;
- ② Effective bosonic model (M. Chernodub, arXiv:1008.1055);
- ③ Effective Nambu-Jona-Lasinio fermionic model (M. Chernodub, arXiv:1101.0117);
- ④ Gauge/gravity duality (N. Callebaut, D. Dudal, H. Verschelde, arXiv:1105.2217; M. Ammon, J. Erdmenger, P. Kerner, M. Strydom, arXiv:1106.4551; ...)
- ⑤ Numerical calculations (ITEP Lattice Group, M. Müller-Preussker)
- ⑥ Ongoing discussions
 - Y. Hidaka and A. Yamamoto, arXiv:1209.0007
 - M. Chernodub, arXiv:1209.3587
 - Chuan Li, Qing Wang, arXiv:1301.7009

Energy of a ρ -meson in the magnetic field:

$$E^2 = m_\rho^2 + eB(2n + 1 - 2S_z) + p_z^2$$

$$\left. \begin{array}{l} \text{Lowest Landau level:} \\ \text{Not moving along magnetic field:} \\ \text{Spin along magnetic field:} \end{array} \right\} \begin{array}{l} n = 0 \\ p_z = 0 \\ S_z = 1 \end{array} \quad E^2 = m_\rho^2 - eB$$

If $eB > eB_c = m_{\rho^2} \approx 0.6 \text{ GeV}^2 \Rightarrow E^2 < 0 \Rightarrow$ Condensation

Lagrangian

$$\begin{aligned} L = & -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{2}\rho^{\mu\nu}\rho_{\mu\nu}^\dagger + m_\rho^2\rho^\mu\rho_\mu^\dagger \\ & - \frac{1}{4}\rho_{\mu\nu}^{(0)}\rho^{(0)\mu\nu} + \frac{m_\rho^2}{2}\rho_\mu^{(0)}\rho^{(0)\mu} + \frac{e}{2g_s}F^{\mu\nu}\rho_{\mu\nu}^{(0)} \end{aligned}$$

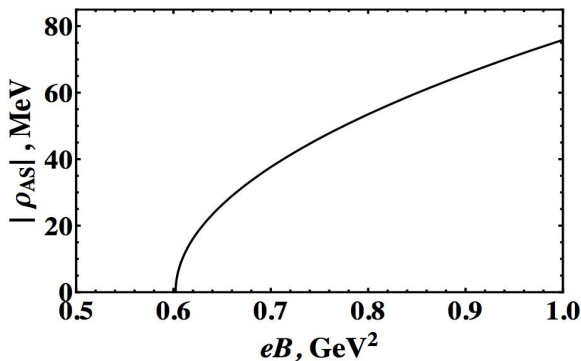
Here

$$\begin{aligned} F_{\mu\nu} &= \partial_\mu A_\nu - \partial_\nu A_\mu \\ \rho_{\mu\nu}^{(0)} &= \partial_\mu \rho_\nu^{(0)} - \partial_\nu \rho_\mu^{(0)} - ig_s(\rho_\mu^\dagger \rho_\nu - \rho_\mu \rho_\nu^\dagger) \\ \rho_{\mu\nu} &= D_\mu \rho_\nu - D_\nu \rho_\mu \\ D_\mu &= \partial_\mu + ig_s \rho_\mu^{(0)} - ieA_\mu \end{aligned}$$

Absolute value of condensate

Condensate(absolute value):

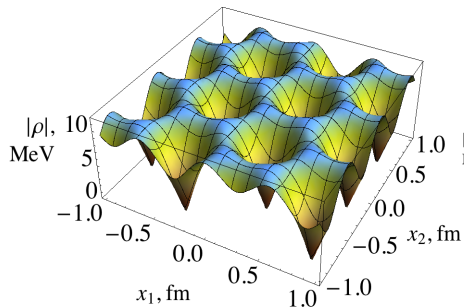
$$|\rho|_0 = \begin{cases} \sqrt{\frac{e(B_{\text{ext}} - B_c)}{2g_s^2}}, & \text{if } B_{\text{ext}} > B_c. \\ 0, & \text{if } B_{\text{ext}} < B_c \end{cases}$$



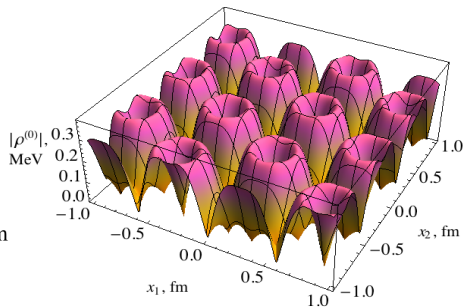
Structure of condensate

Effective bosonic model

$$B = 1.01 B_c$$



Charged ρ -mesons



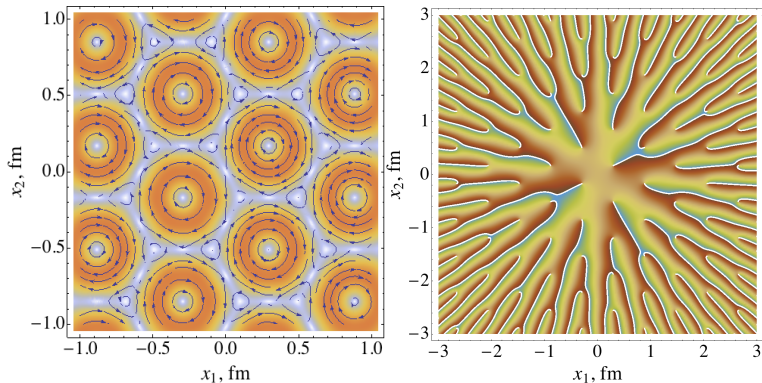
Neutral ρ -mesons

M. Chernodub, J. Doorsselaere, H. Verschelde, arXiv:1111.4401.

Similar results in holographic approach: Y.-Y. Bu, J. Erdmenger, J. P. Shock, M. Strydom, arXiv:1210.6669

Structure of condensate

Superconducting current and the phase of the condensate:



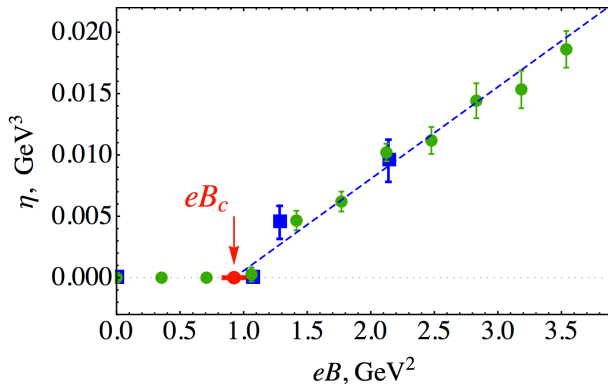
M. Chernodub, J. Doorselaere, H. Verschelde, arXiv:1111.4401.

Numerical calculations in quenched QCD

- Two quark flavours: u, d
- Rho-meson operator: $\rho_\mu = \bar{u}\gamma_\mu d$
- Spin ± 1 along magnetic field: $\rho_\pm = \frac{1}{2}(\rho_1 \pm i\rho_2)$
- Equal time correlator: $G_\pm(z) = \langle \rho_\pm^\dagger(0) \rho_\pm(z) \rangle$
- Condensate: $\lim_{|z| \rightarrow \infty} G_+(z) = |\langle \rho \rangle|^2$
- Spin -1 : $\lim_{|z| \rightarrow \infty} G_-(z) = 0$

Superconducting condensate

Quenched QCD with overlap fermions



Critical field:

$$eB_c = 0.924(77) \text{GeV}^2$$

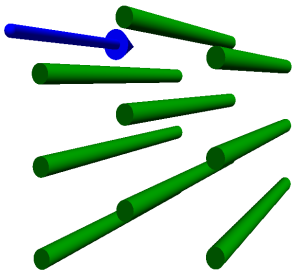
$$m_\rho^2 \sim 1.1 \text{GeV}^2$$

Superconducting vortices.

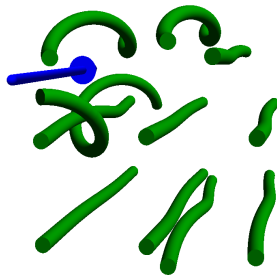
Observables:

- Rho-meson field: $\rho_\mu = \bar{u}\gamma_\mu d$ - cannot be computed in our approach!
- Correlator: $\rho_+(x) \rightarrow \phi(x) = \langle \rho_+^\dagger(0) \rho_+(x) \rangle_f$
- Energy density: $E(x) = \frac{|D_\mu \phi(x)|^2}{|\phi(x)|^2}, D_\mu = \partial_\mu - ieA_\mu$
- Electric current: $j_\mu(x) = \frac{\phi^* D_\mu \phi - \phi (D_\mu \phi)^*}{2i|\phi(x)|^2}$
- Vortex density: $v(x) = \text{sing arg } \phi(x) = \frac{\epsilon^{ab}}{2\pi} \frac{\partial}{\partial x_a} \frac{\partial}{\partial x_b} \arg \phi(x)$

Expectations

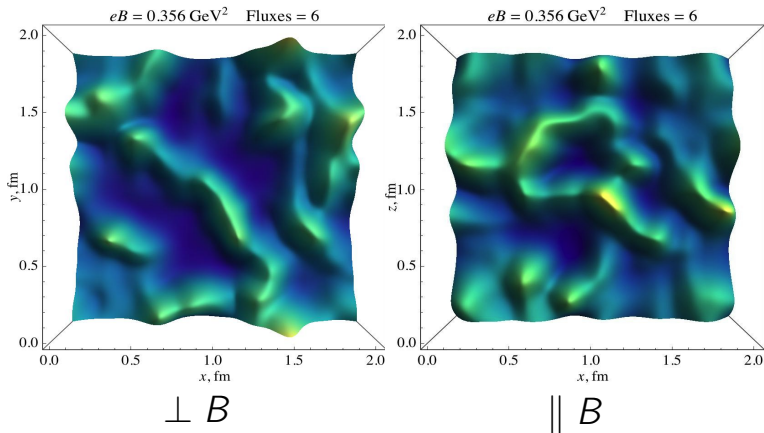


Effective model



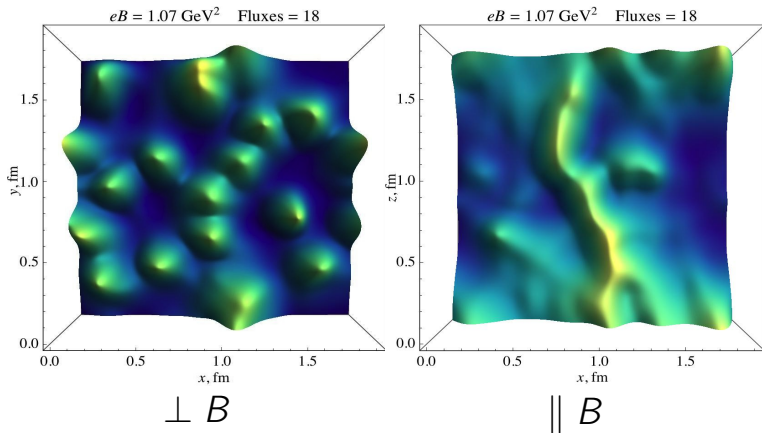
Quantum fluctuations
move and disturb vortices

Superconducting vortices. Energy density. $eB = 0.36 \text{ GeV}^2$

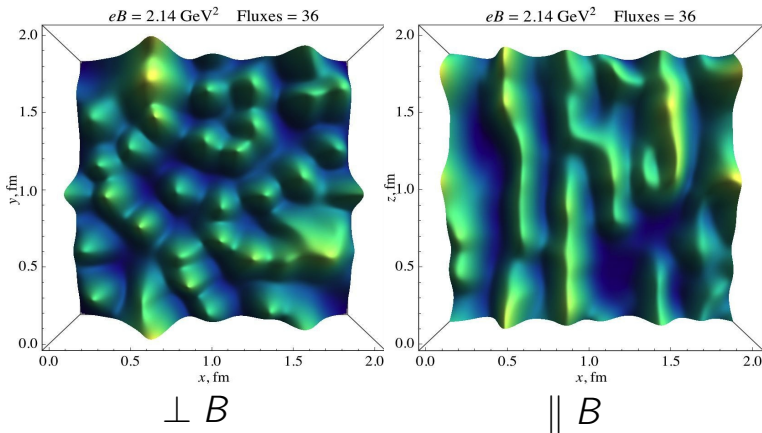


arXiv:1301.6590

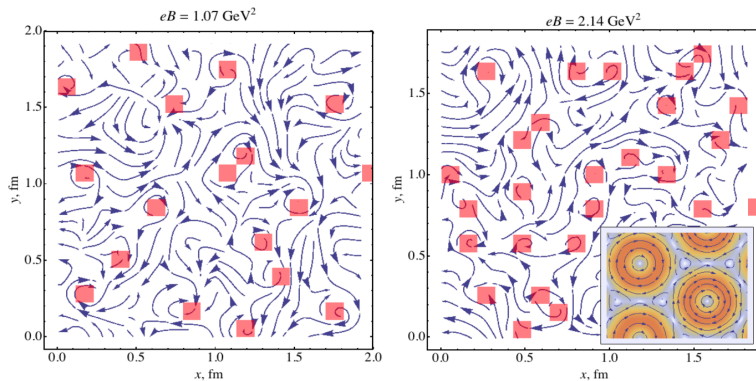
Superconducting vortices. Energy density. $eB = 1.07 \text{ GeV}^2$



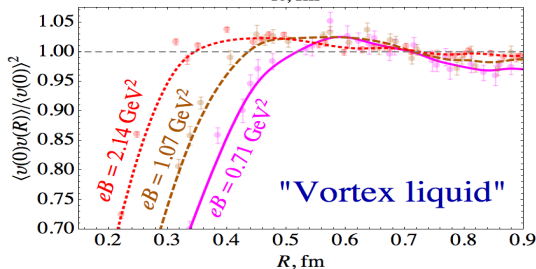
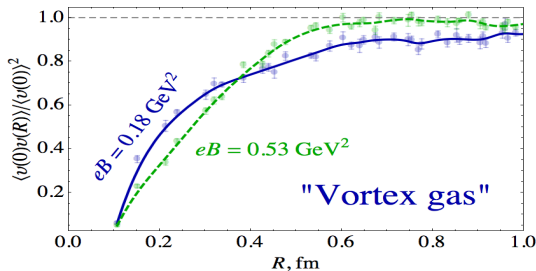
Superconducting vortices. Energy density. $eB = 2.14 \text{ GeV}^2$



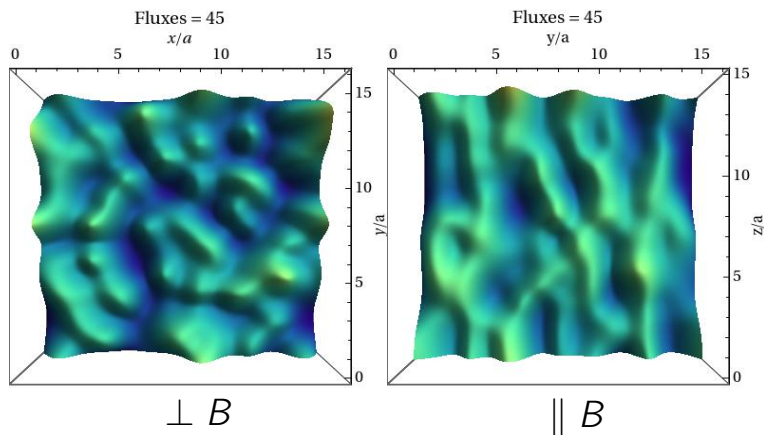
Electric current around vortices



Vortex density-density correlators



First simulations with dynamical fermions



$$eB \sim 1.5 \text{ GeV}^2$$

- In a sufficiently strong magnetic field ρ -meson condensate is formed simultaneously.
- New type of topological defects, " ρ -vortices", emerge.
- Liquid of ρ -vortices is observed in quenched lattice calculations(cf. theory: trigonal lattice)
- Plan: simulations with dynamic fermions.