# Neutron and proton EDM in $N_f = 2+1$ domain-wall fermion

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### Outline

• Motivation of lattice calculation of EDM in  $\theta$  term

- Strategy and method
- (Preliminary) results
- Summary and future work

## 1. Introduction Motivation

- CP violation (CPV) in QCD and BSM
  - EDM is sensitive to CPV in BSM
    - ▶ SM contribution is extremely small, 10<sup>-33</sup> e<sup>-</sup>cm
  - Strong CP problem
  - There are many experimental plans of EDM pEDM experiment @ BNL, nEDM experiment @ ORNL, ILL, FRM-2, FNAL, PSI/KEK/TRIUMF, ... Charged particle (d, p)EDM @ COSY Lepton EDM @ J-PARC, FNAL aim to 10<sup>-29</sup> e<sup>•</sup> cm (present nEDM limit : |d<sub>N</sub><sup>exp</sup>| < 2.9 × 10<sup>-26</sup> e · cm )
  - ▶ BSM (like SUSY) predicts ~10<sup>-27</sup> e·cm
  - Large uncertainty is hadron effect in low-energy scale.
  - Lattice QCD provides important information of BSM prediction.

### 1. Introduction Nucleon EDM in lattice QCD

- θ term contribution
  - Renormalizable
  - Feasible study toward the BSM (quark EDM and chromo EDM) calculation

Bhattacharya et al, Lattice 2012, Lattice 2013

- Nice check of uncertainties in quark model, BChPT
- Some difficulties
  - Statistical noise
    - Gauge background (topological charge, sea quark) which are intrinsically noisy.
    - Disconnected diagram (flavor singlet) should include (In SU(3) limit this is vanishing. Study including disc. diagram is future work).
  - Systematic study
    - Finite volume effect may be significant. (e.g. BChPT discussion)

O'Connell, Savage, PLB633, 319(2006), Guo, Meissner, 1210.5887

• Chiral behavior( $d_N \sim O(m)$ ) is also important check.

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Chiral symmetry on the lattice, and control the O(a) lattice artifact. Good control of chiral behavior without counterterm (cf.Wilson-clover)

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- Error reduction techniques
  - All-mode-averaging (AMA)
  - Efficient way to reduce the statistical error of correlator

Blum, Izubuchi, ES, arXiv:1208.4349 [hep-lat], ES (lattice 2012)

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- EDM form factor
  - Extraction from matrix element including CPV of  $\theta$  angle
  - Extrapolation into physical kinematics,  $-q^2 = 0$  and m = m<sub>phy</sub>

## 2. EDM calculation on the lattice **Domain-wall fermion**

- Chiral symmetry
  - L, R fermion is localized on boundaries
    - $\Rightarrow$  Chiral symmetry is realized (if  $L_s \rightarrow \infty$ ).
  - Remaining good chiral symmetry m<sub>res</sub> ~ exp(-L<sub>s</sub>)
  - Reasonable computational cost
- RBC/UKQCD collaboration
  - Generation of N<sub>f</sub> = 2+1 configurations.
  - Various lattice size, lattice cut-off, quark mass are available.
  - Many studies of Keon physics, nucleon physics, finite temperature , ...



[Blum Soni, (97), CP-PACS(99), RBC(00), RBC/UKQCD. (05--) ]

Blum, Izubuchi, ES, 1208.4349, ES (lattice 2012), Chulwoo, plenary on Fri

- Covariant approximation averaging (CAA)
  - For original correlator O, (unbiased) improved estimator is defined as

$$\mathcal{O}^{(\text{imp})} = \mathcal{O}^{(\text{rest})} + \frac{1}{N_G} \sum_{g \in G} \mathcal{O}^{(\text{appx}),g}, \quad \mathcal{O}^{(\text{rest})} = \mathcal{O} - \mathcal{O}^{(\text{appx})}$$

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  - $P_n(\lambda)$  is polynomial approximation of  $I/\lambda$ 
    - Low mode part :# of eigen mode
    - Mid-high mode : degree of poly.



ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\langle n(P_1)|J_{\mu}^{\text{EM}}|n(P_2)\rangle_{\theta} = \bar{u}_N^{\theta} \Big[\underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N}\gamma_5\sigma_{\mu\nu}Q_{\nu}}_{\text{P,T-odd}} + \underbrace{F_1\gamma_{\mu} + \frac{F_2}{2m_N}\sigma_{\mu\nu}Q_{\nu}}_{\text{P,T-even}} + \cdots \Big]u_N^{\theta}$$

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$$\langle n(P_1) | J_{\mu}^{\text{EM}} | n(P_2) \rangle_{\theta} = \bar{u}_N^{\theta} \Big[ \underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N} \gamma_5 \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-odd}} + \underbrace{F_1 \gamma_{\mu} + \frac{F_2}{2m_N} \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-even}} + \cdots \Big] u_N^{\theta} \\ \sum_s u_N^{\theta}(s) \bar{u}_N^{\theta}(s) = \frac{ip \cdot \gamma + m_N e^{i\alpha_N^{\theta} \gamma_5}}{2E_N} \\ \end{bmatrix}$$
 CPV phase  $\alpha_N$  in nucleon propagator

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

$$\begin{split} \langle n(P_1) | J_{\mu}^{\text{EM}} | n(P_2) \rangle_{\theta} &= \bar{u}_N^{\theta} \Big[ \underbrace{\frac{F_3^{\theta}(Q^2)}{2m_N} \gamma_5 \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-odd}} + \underbrace{F_1 \gamma_{\mu} + \frac{F_2}{2m_N} \sigma_{\mu\nu} Q_{\nu}}_{\text{P,T-even}} + \cdots \Big] u_N^{\theta} \\ &\sum_s u_N^{\theta}(s) \bar{u}_N^{\theta}(s) = \frac{ip \cdot \gamma + m_N e^{i\alpha_N^{\theta} \gamma_5}}{2E_N} \Big] \quad \text{CPV phase } \alpha_{\text{N}} \text{ in nucleon propagator} \\ \langle \theta | \eta_N J_{\mu}^{\text{EM}} \bar{\eta}_N | \theta \rangle &= \langle 0 | \eta_N J_{\mu}^{\text{EM}} \bar{\eta}_N | 0 \rangle + i\theta \langle 0 | \eta_N J_{\mu}^{\text{EM}} Q \bar{\eta}_N | 0 \rangle \end{split}$$

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$$\begin{split} \langle n(P_{1})|J_{\mu}^{\mathrm{EM}}|n(P_{2})\rangle_{\theta} &= \bar{u}_{N}^{\theta} \Big[ \underbrace{\frac{F_{3}^{\theta}(Q^{2})}{2m_{N}}\gamma_{5}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-odd}} + \underbrace{F_{1}\gamma_{\mu} + \frac{F_{2}}{2m_{N}}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-even}} + \cdots \Big] u_{N}^{\theta} \\ &\sum_{s} u_{N}^{\theta}(s)\bar{u}_{N}^{\theta}(s) = \frac{ip\cdot\gamma + m_{N}e^{i\alpha_{N}^{\theta}\gamma_{5}}}{2E_{N}} \Big] \\ \mathcal{C}\mathsf{PV} \text{ phase } \alpha_{\mathsf{N}} \text{ in nucleon propagator} \\ \langle \theta|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|\theta\rangle &= \langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|0\rangle + i\theta\langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}Q\bar{\eta}_{N}|0\rangle \\ \langle 0|\eta_{N}(t_{1})J_{\mu}^{\mathrm{EM}}(t)Q\bar{\eta}_{N}(t_{0})|0\rangle & \qquad \\ = \frac{\alpha_{N}}{2}\gamma_{5}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} + \frac{1+\gamma_{4}}{2}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{\alpha_{N}}{2}\gamma_{5} & \qquad \\ \end{bmatrix} \text{ Computation} \\ + \frac{1+\gamma_{4}}{2}\Big[F_{3}\frac{q_{\nu}\gamma_{5}\sigma_{\mu\nu}}{2m_{N}} + F_{A}(iq^{2}\gamma_{\mu}\gamma_{5} - 2m_{N}q_{\mu}\gamma_{5})\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} & \qquad \\ \end{bmatrix} \underbrace{k_{N}}^{ip\cdot\gamma + m_{N}}_{ip\cdot\gamma + m_{N}} \\ = \underbrace{k_{N}}^{ip\cdot\gamma + m_{N}}_{ip\cdot\gamma + m$$

ES, et al., Phys. Rev. D72, 014504 (2005), Berruto, et al., Phys. Rev. D73, 05409 (2006).

#### EM Matrix element

$$\begin{split} \langle n(P_{1})|J_{\mu}^{\mathrm{EM}}|n(P_{2})\rangle_{\theta} &= \bar{u}_{N}^{\theta} \Big[ \underbrace{\frac{F_{3}^{\theta}(Q^{2})}{2m_{N}}\gamma_{5}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-odd}} + \underbrace{F_{1}\gamma_{\mu} + \frac{F_{2}}{2m_{N}}\sigma_{\mu\nu}Q_{\nu}}_{\mathrm{P,T-even}} + \cdots \Big] u_{N}^{\theta} \\ &\sum_{s} u_{N}^{\theta}(s)\bar{u}_{N}^{\theta}(s) = \frac{ip\cdot\gamma + m_{N}e^{i\alpha_{N}^{\theta}\gamma_{5}}}{2E_{N}} \Big] \\ \mathcal{C}\mathsf{PV} \text{ phase } \alpha_{\mathsf{N}} \text{ in nucleon propagator} \\ \langle \theta|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|\theta\rangle &= \langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}\bar{\eta}_{N}|0\rangle + i\theta\langle 0|\eta_{N}J_{\mu}^{\mathrm{EM}}Q\bar{\eta}_{N}|0\rangle \\ \langle 0|\eta_{N}(t_{1})J_{\mu}^{\mathrm{EM}}(t)Q\bar{\eta}_{N}(t_{0})|0\rangle & \qquad \\ = \frac{\alpha_{N}}{2}\gamma_{5}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} + \frac{1+\gamma_{4}}{2}\Big[F_{1}\gamma_{\mu} + F_{2}\frac{q_{\nu}\sigma_{\mu\nu}}{2m_{N}}\Big]\frac{\alpha_{N}}{2}\gamma_{5} & \qquad \\ \end{bmatrix} \text{ Computation} \\ + \frac{1+\gamma_{4}}{2}\Big[F_{3}\frac{q_{\nu}\gamma_{5}\sigma_{\mu\nu}}{2m_{N}} + F_{A}(iq^{2}\gamma_{\mu}\gamma_{5} - 2m_{N}q_{\mu}\gamma_{5})\Big]\frac{ip\cdot\gamma + m_{N}}{2E_{N}} & \qquad \\ \end{bmatrix} \underbrace{ip\cdot\gamma + m_{N}}{2E_{N}} & \qquad \\ \end{bmatrix} \text{ Extraction} \end{split}$$

- Subtraction of CP-odd phase,  $\alpha_{\rm N}$ , in n propagator and CP-even part  ${\sf F}_{\rm I,2}$ 

$$d_N = \lim_{Q^2 \to 0} F_3(Q^2) / 2m_N$$

## 3. (Preliminary) Results **Parameters**

DWF

- >  $24^3 \times 64$  lattice,  $a^{-1} = 1.73$  GeV (~3 fm<sup>3</sup> lattice)
- $L_{\rm s} = 16$  and  $am_{\rm res} = 0.003$
- m = 0.005, 0.01 corresponding to  $m_{\pi} = 0.33, 0.42$  GeV
- Two temporal separation of N sink and source in 3 pt. function

$$t_{sep} = 12 (t_{source} = 0, t_{sink} = 12), t_{sep} = 8 (t_{source} = 0, t_{sink} = 8)$$

# configs = 751 (m=0.005), 700 (m=0.01) [ $t_{sep}$  = 12] # configs = 180 (m=0.005) [ $t_{sep}$  = 8]

AMA

- # of low-mode :  $N_{\lambda} = 400 \text{ (m=0.005)}, 180 \text{ (m=0.01)}$
- Stopping condition, |r| < 0.003</p>
- ▶  $N_G = 32$  (2 separation for spatial, 4 separation for temporal direction of source localtion) → effectively O(10<sup>4</sup>) statistics

#### 3. (Preliminary) Results

### $\alpha_N$

- Projection with  $\gamma_5$  for 2 pt with Q charge, global fitting with  $\operatorname{tr}\left[\gamma_5 \langle N(t)\bar{N}(0)Q \rangle\right] = Z_N \frac{2m_N}{E_N} \alpha_N e^{-E_N t} + O(e^{-E_N *})$
- By using AMA, this factor is determined within 15 % error.
- It does not depend on smearing and momentum, but mass dependence is not so clear.



### 3. (Preliminary) Results Subtraction term and 3pt function

- Splitting EDM form factor into two parts:  $F_3 = F_Q + F_{\alpha}, F_Q = C(m_N) \langle N J_t^{\text{EM}} \bar{N} Q \rangle, F_{\alpha} = F(\alpha_N, F_{1,2})$
- $F_{\alpha}$  is good precision, and fluctuation of  $F_{O}$  is large.



## 3. (Preliminary) Results Comparison with $\mu = t, z$

- EDM form factor is given from two directions of EM current
- Two signals are consistent, and data in t direction is much stable.



### 3. (Preliminary) Results Comparison with different t<sub>sep</sub>

- The sink and source separation in 3pt function enables us to control the statistical noise and excited state contamination
  - Short: statistical fluctuation < excited state contamination</p>
  - Long: statistical fluctuation > excited state contamination
- Comparison
  - $t_{sep} = 12 \text{ (blue),}$ [ $N_{conf} = 751$ ]  $t_{sep} = 8 \text{ (green)}$ [ $N_{conf} = 180$ ]
  - Good consistency between them.
  - Precision in t<sub>sep</sub>=8 is much better.



## 3. (Preliminary) Results $-q^2 = 0$ extrapolation

- Fitting data of EDM form factor at each momenta.
- Open(t<sub>sep</sub> = 8),
   filled (t<sub>sep</sub> = 12)
- Fitting function
  - 3 point linear :
  - $-q^2 < 0.55 \text{ GeV}^2$
  - 2 point linear: -q<sup>2</sup> < 0.4 GeV<sup>2</sup>
- Estimate of systematic error of extrapolation



## 3. (Preliminary) Results Mass dependence

#### Comparison with full QCD results

- DWF results are in the lightest quark mass.
- Statistical error is dominant rather than systematic one.
- Central value is 10 times larger than models.



## 3. (Preliminary) Results Statistical error

- Comparison between AMA error reduction and number of configurations.
- Number of configurations : reduce stat. error and relating to Q distribution
   AMA error reduction : reduce stat. error



### 4. Summary Summary and future plan

- Nucleon EDM in  $N_f = 2+1$  DWF in  $\theta$  vacuum
  - Signal of EDM within 40% statistical error using AMA techniques.
  - ▶ 3-pt function is still noisy.
  - Short t<sub>sep</sub> allows us to reduce the statistical error without large excited state contamination effect.

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  - Signal of EDM within 40% statistical error using AMA techniques.
  - ▶ 3-pt function is still noisy.
  - Short t<sub>sep</sub> allows us to reduce the statistical error without large excited state contamination effect.
- (Near) physical point of DWF configurations
  - Ensembles near physical points and large volume are available.
  - > AMA with Möbius-DWF approximation is helpful.
  - Remove chiral extrapolation  $\rightarrow$  less than 10% precision

Lattice size	Physical size	а	L <sub>s</sub>	Gauge action	Pion mass
$32^{3} \times 64$	4.6 fm <sup>3</sup>	0.135 fm	32	DSDR	171 241 MeV
48 <sup>3</sup> × 96	5.5 fm <sup>3</sup>	0.115 fm	16	Iwasaki	135 MeV

#### Thank you for your attention !

### Backup

### Nucleon EDM in the SM

- CKM phase in EDM
  - No CPV phase in I-loop (|V<sub>dq</sub>|<sup>2</sup> : no phase) and 2-loop diagram (unitarity).



• Three-loop order(short) or pion loop correction (long)  $d_N^{\text{KM short}} \sim -10^{-34} \,\text{e} \cdot \text{cm}, \quad d_N^{\text{KM long}} \sim 10^{-30} - 10^{-32} \,\text{e} \cdot \text{cm}$ 

 $\Rightarrow d_N^{\rm KM} = d_N^{\rm KM\,short} + d_N^{\rm KM\,long} \simeq 10^{-30} - 10^{-32}\,{\rm e\cdot cm}$ 

6-order magnitude below the experimental upper limit.

•  $\theta$  term in the QCD Lagrangian

Renormalizable and CP-violating from in topological charge density.

$$\mathcal{L}_{\theta} = \bar{\theta} \frac{1}{64\pi^2} G \widetilde{G}, \quad \bar{\theta} = \theta + \arg \det M, \quad \arg \det M \sim \eta \frac{m_u m_d m_s}{m_u m_d + m_d m_s + m_s m_u}$$

- m heta term has been estimated as  $ar heta < 10^{-9\pm 1}$  Crewther, et al. (1979), Ellis, Gaillard (1979)
- Unnatural cancellation (strong CP problem)

### Nucleon EDM in the BSM

- Possible higher dimension operators
  - In supersymmery (SUSY) model there is CPV phase from 1-loop (Im(g<sub>L</sub> g<sub>R</sub><sup>\*</sup>)≠0)



 $\blacktriangleright$  CPV effective Hamiltonian with higher dimension than  $\theta$  term

$$H_{CP} = \sum_{k} C_{k}(\mu) \mathcal{O}_{k} \qquad \qquad \mathcal{O}_{qEDM} = d_{q} \bar{q} (\sigma \cdot F) \gamma_{5} q \quad : \text{Quark-photon (5-dim)} \\ \mathcal{O}_{cEDM} = d_{q}^{c} \bar{q} (\sigma \cdot G) \gamma_{5} q \quad : \text{Quark-gluon (5-dim)} \\ \mathcal{O}_{Weinberg} = d^{G} G G \tilde{G} \quad : \text{Pure gluonic (6-dim)} \end{cases}$$

Contribution to nEDM in low energy model

**BChPT**:

QCD sum rule:



### Lattice methods on EDM

#### Spectrum method

- I. S.Aoki and A. Gocksch, Phys. Rev. Lett. 63, 1125 (1989).
- 2. S.Aoki, A. Gocksch, A.V. Manohar, S. R. Sharpe, Phys. Rev. Lett. 65, 1092 (1990), in which they discussed about the possible lattice artifact in ref.1 results
- 3. ES, et al., for CP-PACS collaboration, Phys. Rev. D75, 034507 (2007)
- 4. ES, S. Aoki, Y. Kuramashi, Phys. Rev. D78, 014503 (2008)

#### Form factor

- I. ES, et al., for CP-PACS collaboration, Phys. Rev. D72, 014504 (2005).
- 2. Berruto, et al. for RBC collaboration, Phys. Rev. D73, 05409 (2006).
- 3. ES et al., Lattice 2008.

#### Imaginary θ

- I. T. Izubuchi, Lattice 2007.
- 2. Horsley et al., arXiv:0808.1428 [hep-lat]

### EM form factor



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