

Quark localization in QCD above T_c

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Dirac operator

Continuum

- $\not{D} = \gamma_\mu \partial_\mu + m$
- Quark propagator \sim inverse of \not{D}
- Banks-Casher $\langle \bar{\psi} \psi \rangle = \lim_{m \rightarrow 0} \lim_{V \rightarrow \infty} \frac{\pi \rho(0)}{V}$
- Importance of the low lying eigenvalues

Symmetries

- Anti-hermitian: $D^\dagger = -D$
- Chiral symmetry: $\{D, \gamma_5\} = 0$
- γ_5 hermiticity: $D^\dagger = \gamma_5 D \gamma_5$
- Eigenvalues come in complex conjugate pairs



Dirac spectrum below and above T_c

- $T < T_c$: Random Matrix statistics in the ε regime
[Shuryak, Verbaarschot(1993)]
- $T \simeq T_c$: Localized eigenmodes
[Garcia-Garcia, Osborn(2006)]
- $T > T_c$: lowest eigenmodes of the $SU(2)$ overlap Dirac operator follow Poisson statistics [T.G.Kovacs(2010)]
- $SU(2)$ quenched: localized-delocalized transition
[T.G.Kovacs, FP(2010)]

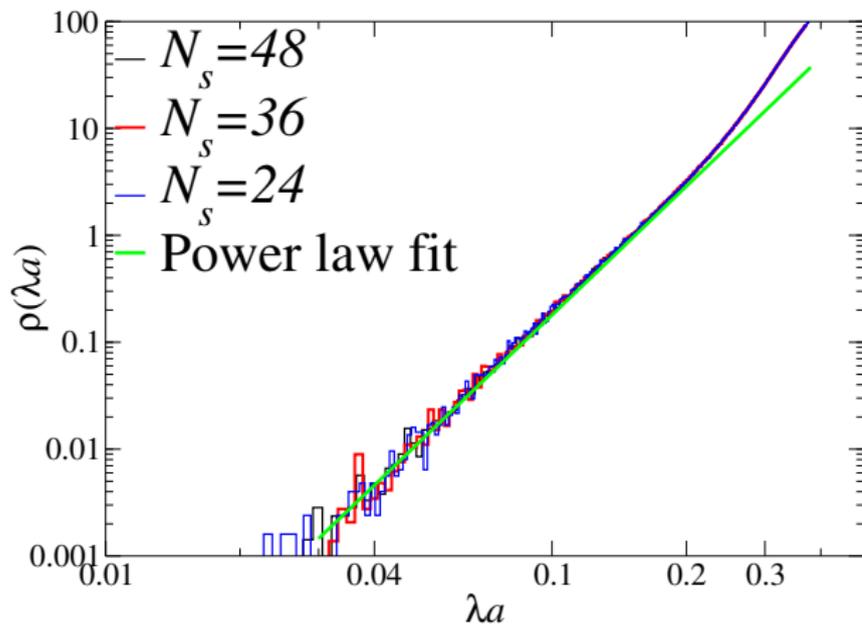
QCD

- $SU(3)$ gauge theory 2 + 1 flavors of dynamical quarks with physical masses [Wuppertal Budapest]
- Determinant suppresses the lowest quark modes



Spectral density of the Dirac operator

$$N_f = 4, a = 0.125\text{fm}, V = 27\text{fm}^3 \rightarrow 215\text{fm}^3, T = 2.6T_c$$



- Chiral symmetry $\lim_{\lambda \rightarrow 0} \rho(\lambda) = 0$
- Scales properly with the volume
- No evidence of spectral gap
- $\rho(\lambda) = c\lambda^4$ for small λ



The spectral statistics and localization of D above T_c

RMT type spectrum

- The spectral density is large
- The eigenvectors have spatial overlap: delocalized modes
- Typical fluctuations in the gauge field **mix** the eigenmodes
- Eigenvalues are correlated: Wigner-Dyson statistics

Poisson type spectrum

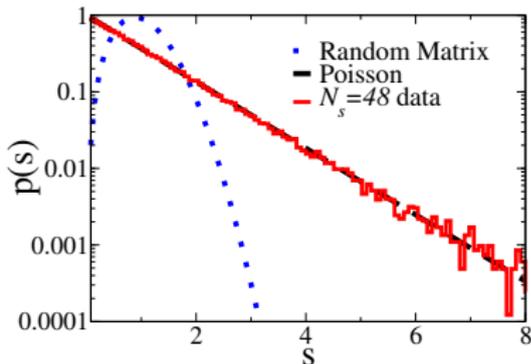
- The spectral density is small
- The eigenvectors do not have spatial overlap: localized modes
- Typical fluctuations in the gauge field **do not mix** the eigenmodes
- Eigenvalues are uncorrelated: Poisson statistics



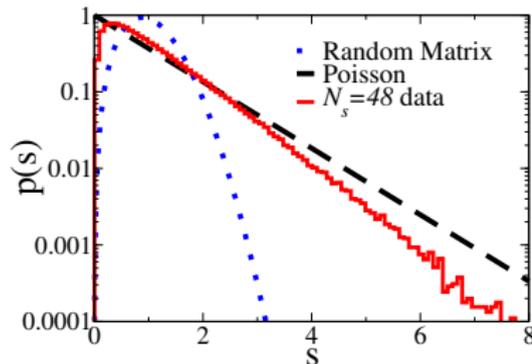
Unfolded level spacing distribution (ULSD)

$$N_t = 4, N_s = 48, \quad s = \frac{\lambda_{n+1} - \lambda_n}{\langle \lambda_{n+1} - \lambda_n \rangle}$$

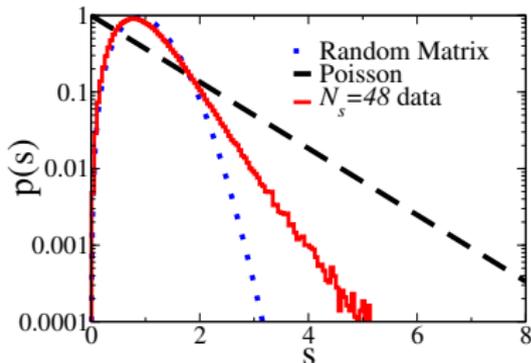
$0.15 < \lambda a < 0.19$



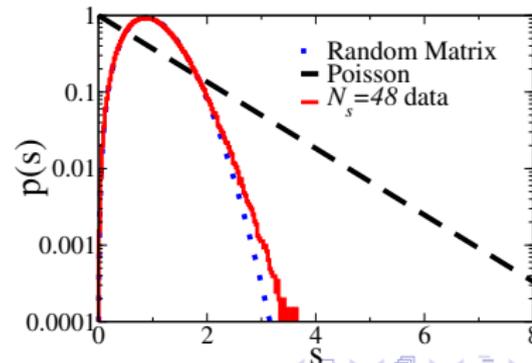
$0.31 < \lambda a < 0.32$



$0.34 < \lambda a < 0.35$

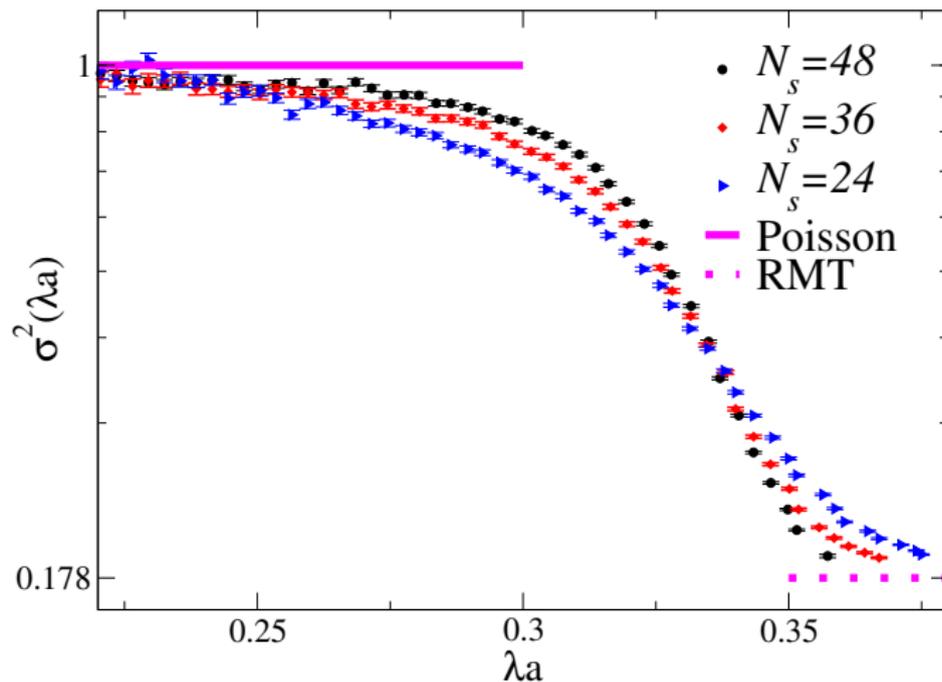


$0.36 < \lambda a < 0.373$



Thermodynamic limit

Variance of the $ULSD$ in the spectrum for several volumes, $N_t = 4$, $a = 0.125\text{fm}$

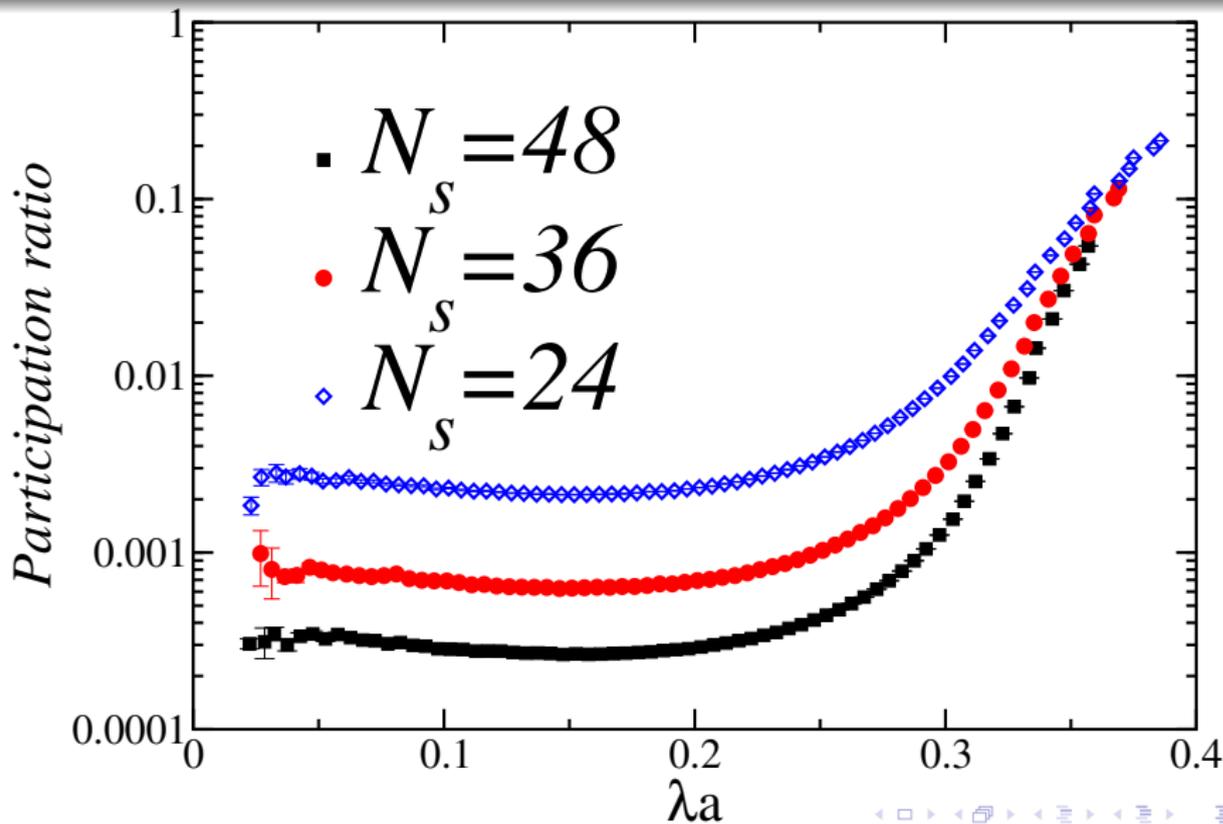


- We can define a mobility edge (λ_c) for all V
- $\lambda_c \equiv$ Location of the inflection point of the curves
- ? Is it a real phase transition?

✓ M. Giordano's talk



QCD Dirac spectrum: Localization-Delocalization transition in terms of eigenvector statistics



Effects of localized modes

- quark propagator \sim inverse of D
- Hadronic correlators: build up from quark propagators
- Lowest modes: largest weight
- They have negligible contribution to the long distance correlators

Questions:

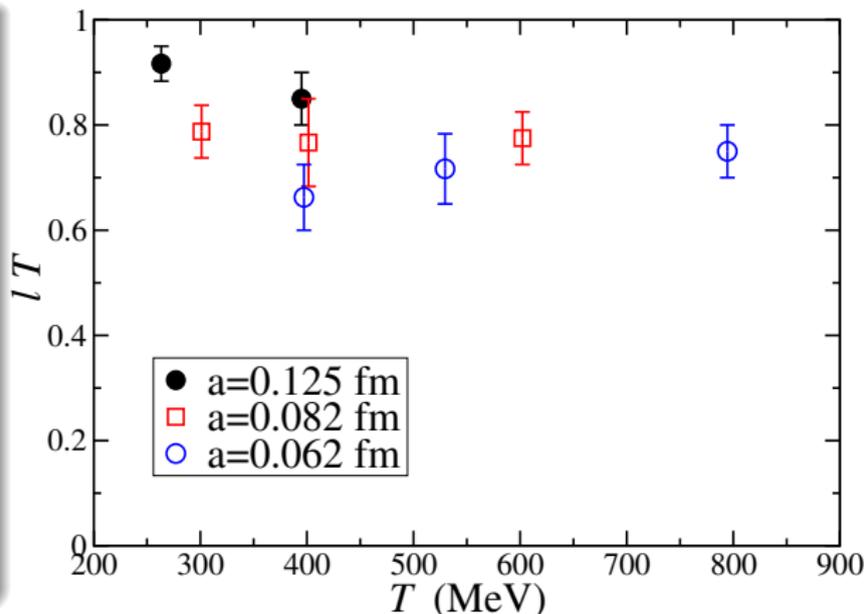
- How large is the localization length?
- How far up in the spectrum are the eigenmodes localized?



Localization length

What determines the localization length?

- Temperature squeezes the modes in the time direction
- Localized modes are squeezed in the same way in the spatial directions



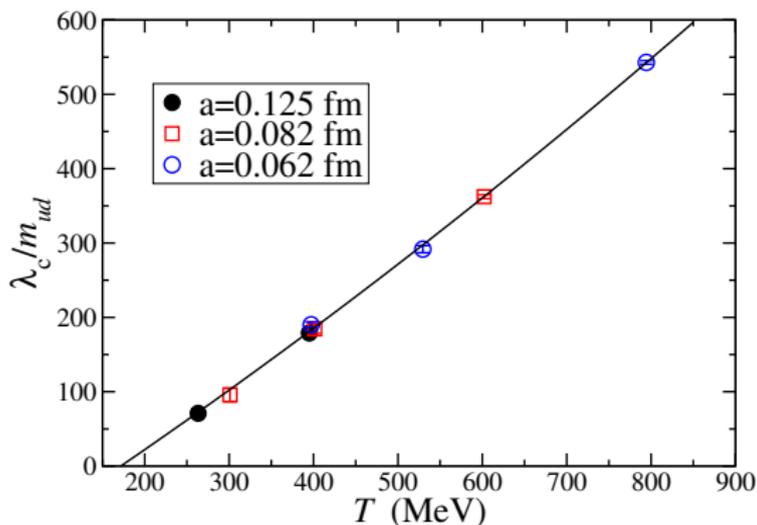
- Localization length \sim inverse temperature



Mobility edge in the continuum limit

How λ_c behaves in the continuum?

- λ_c introduces an effective gap in the spectrum
- Its renormalization is similar to the quark mass
- λ_c/m_{ud} tends to finite value in the continuum limit
- Localization is physical



- $\lambda_c(T_c) = 0 \rightarrow T_c \simeq 170$ MeV



Conclusion

- Lowest part of the QCD Dirac operator consists of Poisson modes
- The Poisson modes are localized to a distance scale set by the inverse temperature
- Mobility edge is effectively a gap in the spectrum
- Is this a real phase transition? See M. Giordano's talk



Thank you for your attention!



References

- [Shuryak, Verbaarschot(1993)] E. V. Shuryak and J. J. M. Verbaarschot, Nucl. Phys. A **560**, 306 (1993) [hep-th/9212088].
- [Garcia-Garcia, Osborn(2006)] A. M. Garcia-Garcia and J. C. Osborn, Phys. Rev. D **75**, 034503 (2007) [hep-lat/0611019].
- [T.G.Kovacs(2010)] T. G. Kovacs, Phys. Rev. Lett. **104**, 031601 (2010) [arXiv:0906.5373 [hep-lat]].
- [T.G.Kovacs,FP(2010)] T. G. Kovacs and F. Pittler, Phys. Rev. Lett. **105**, 192001 (2010) [arXiv:1006.1205 [hep-lat]].

