

# The anomalous mass dimension from the techniquark propagator in Minimal Walking Technicolor

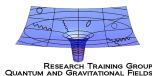
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[August, Maas, 2013: JHEP 1307 (2013) 001, arXiv:1304.4423 [hep-lat]]

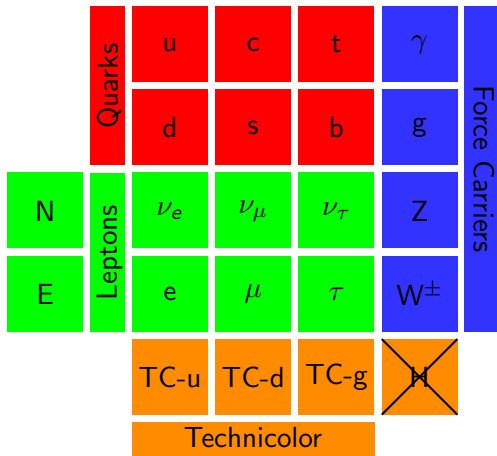
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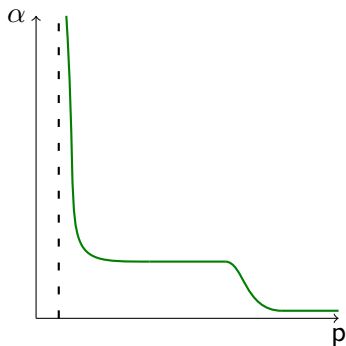
# Minimal Walking Technicolor



MWT:

- gauge group SU(2)
- two adjoint flavours
- extra pair of Leptons

# Walking



walking dynamics

$$\langle \bar{Q}Q \rangle_{ETC} \sim \left( \frac{\Lambda_{ETC}}{\Lambda_{TC}} \right)^\gamma \langle \bar{Q}Q \rangle_{TC}$$

$\gamma$  is the anomalous mass dimension of the quark propagator

# Calculation

## Method

- Wilson fermions on the lattice
- inversion with biconjugate gradient method
- lattice corrections
- configurations provided by L. D. Debbio, B. Lucini, A. Patella, C. Pica, A. Rago

## Setting

- gauge fixing: minimal Landau gauge  $\partial_\mu A_\mu = 0$
- one lattice per value

# Quantities

Propagator

$$S(p) = \delta^{ab} \frac{1}{-A(p^2)i\not{p} + B(p^2)} = \delta^{ab} Z(p^2) \frac{i\not{p} + M(p^2)}{p^2 + M(p^2)^2}$$

$$Z(p^2) = \frac{1}{A(p^2)}, \quad M(p^2) = \frac{B(p^2)}{A(p^2)}$$

Renormalization

$$S(p = \mu) = \frac{1}{-i\not{p} + m_{PCAC}}$$

Schwinger functions

$$\Delta_v = \frac{1}{\pi} \int_0^\infty dp \cos(tp) \frac{Z(p^2)}{p^2 + M(p^2)^2}$$

$$\Delta_s = \frac{1}{\pi} \int_0^\infty dp \cos(tp) \frac{Z(p^2)M(p^2)}{p^2 + M(p^2)^2},$$

# Lattice Corrections

$$\rho A_L = \frac{\text{tr} \gamma_0 S_L}{(\text{tr} \gamma_0 S_L)^2 + (\text{tr} S_L)^2}$$

$$B_L = \frac{\text{tr} S_L}{(\text{tr} \gamma_0 S_L)^2 + (\text{tr} S_L)^2}$$

free case:

$$\rho A_L^{\text{free}} = \sin\left(\frac{2\pi P_0}{L}\right)$$

$$B_L^{\text{free}} = m + 1 - \cos\left(\frac{2\pi P_0}{L}\right)$$

→ momentum dependent lattice artifacts

$$A = \frac{A_L}{A_L^{\text{free}}}$$

$$B = \frac{B_L \cdot m}{B_L^{\text{free}}}$$

# Mass function

explicit chiral symmetry breaking

$$M(p) = M(\mu) \left( \omega \ln \left( \frac{p^2}{\mu^2} \right) + 1 \right)^{-\gamma}$$

spontaneously broken

$$M(p) = \frac{2\pi^2\gamma_f}{3} \frac{-\langle \bar{\Psi}\Psi \rangle}{p^2 \left( \frac{1}{2} \ln \frac{p^2}{\Lambda^2} \right)^{1-\gamma_f}}$$

→ fit function which interpolates both

$$M(p) = \frac{2\pi^2(1-\gamma)}{3} \frac{-\langle \bar{\Psi}\Psi \rangle}{(p+a^2)^{2b} \left( \frac{1}{2} \ln \frac{p^2+c^2}{\Lambda^2} \right)^\gamma}$$

everything except p is a free fit parameter!

## Wrong phase

- most interesting physics for small quark masses

- observed spatial center-breaking transition

L. D. Debbio et al.,2010:arXiv:1004.3206 [hep-lat]

bounds: below  $am = -0.975$  for  $N_s = 8$ , below  $am = -1.05$  for  $N_s = 12$  and between  $am = -1.05$  and  $am = -1.15$  for  $N_s = 16$

- this suggests changing bare mass at fixed  $\beta$  changes the lattice spacing  $\rightarrow$  volume shrinks toward the chiral limit

- supported by gluonic observables A. Maas,2011: arXiv:1102.5023

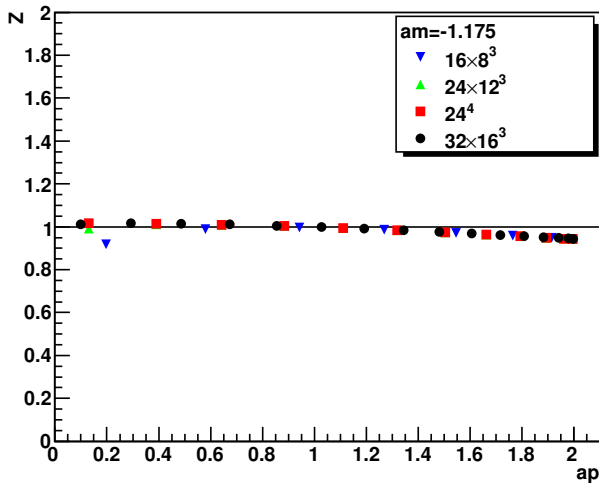
- phase transition due the shrinking volume like in YM-theory

$\rightarrow$  we assume this is a pure lattice artifact

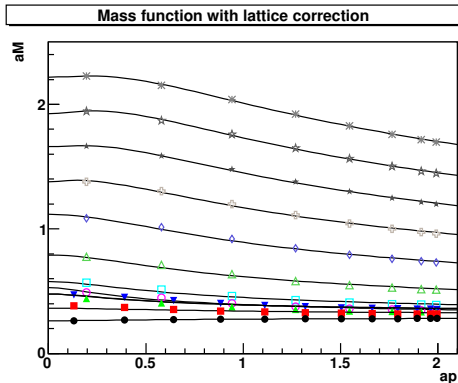


# Wave function renormalization

Wave function renormalization with lattice correction



# Mass function



✱	$16 \times 8^3$ , $am=0.5$ , $am_{PCAC}=2.02$
☆	$16 \times 8^3$ , $am=0.25$ , $am_{PCAC}=1.74$
★	$16 \times 8^3$ , $am=0$ , $am_{PCAC}=1.46$
⊕	$16 \times 8^3$ , $am=-0.25$ , $am_{PCAC}=1.18$
◇	$16 \times 8^3$ , $am=-0.5$ , $am_{PCAC}=0.897$
△	$16 \times 8^3$ , $am=-0.75$ , $am_{PCAC}=0.617$
□	$16 \times 8^3$ , $am=-0.9$ , $am_{PCAC}=0.449$
○	$16 \times 8^3$ , $am=-0.95$ , $am_{PCAC}=0.393$
▼	$24 \times 12^3$ , $am=-0.95$ , $am_{PCAC}=-0.393$
▲	$16 \times 8^3$ , $am=-0.975$ , $am_{PCAC}=-0.365$
■	$24 \times 12^3$ , $am=-1$ , $am_{PCAC}=-0.337$
●	$24 \times 12^3$ , $am=-1.05$ , $am_{PCAC}=-0.277$

$$M(p) = \frac{2\pi^2(1-\gamma)}{3} \frac{-\langle \bar{\Psi}\Psi \rangle}{(p+a^2)^{2b} \left( \frac{1}{2} \ln \frac{p^2+c^2}{\Lambda^2} \right)^\gamma}$$

## Fit parameters

$am_0$	$N_t$	$2b$	$\gamma$
0.5	16	0.18(2)	0.033(6)
0.25	16	0.18(2)	0.028(8)
0	16	0.18(2)	0.020(8)
-0.25	16	0.17(2)	0.009(7)
-0.5	16	0.15(2)	0.01(2)
-0.75	16	0.11(2)	0.03(2)
-0.9	16	0.06(3)	0.04(2)
-0.95	16	0.01(7)	0.07(6)
-0.95	24	0.02(3)	0.04(2)
-0.975	16	0.04(11)	0.11(11)
-1	24	0.00(2)	0.03(2)
-1.05	24	0.02(4)	0.00(3)

$$M(p) = \frac{2\pi^2(1-\gamma)}{3} \frac{-\langle \bar{\Psi} \Psi \rangle}{(p+a^2)^{2b} \left( \frac{1}{2} \ln \frac{p^2+c^2}{\Lambda^2} \right)^\gamma}$$

# Fit parameters

$am_0$	$N_t[\text{TeV}]$	$\gamma$
-0.90	16	0.50(1)
-0.95	16	0.33(1)
-0.95	24	0.23(1)
-1.00	24	0.10(1)

$$M(p) = \frac{1}{a} (\omega \ln(p) + k)^{-\gamma}$$

L. D. Debbio et al.,2010:arXiv:1004.3206 [hep-lat]

A. Patella,2012: arXiv:1204.4432 [hep-lat]

J. Giedt, E. Weinberg,2012: arXiv:1201.6262 [hep-lat]

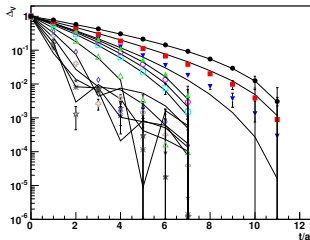
F. Bursa et al.,2009:, arXiv:0910.4535 [hep-ph]

T. DeGrand et al.,2011: arXiv:1102.2843 [hep-lat]

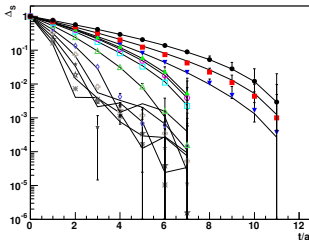
B. Lucini,2009: arXiv:0911.0020 [hep-ph]

# Schwinger functions

Vector Schwinger function



Scalar Schwinger function



bent down  $\rightarrow$  unphysical particle

# Conclusion

- direct calculation of the anomalous mass dimension of the quark propagator in MWT
- for small masses no spontaneously broken symmetry
- small anomalous mass dimension for explicit chiral symmetry breaking is in line with previous results
- techniquark unphysical particle