



Chiral Symmetry Breaking from Center Vortices

in coop. with M. Faber, U. M. Heller and T. Schweigler

arXiv: 1212.3737 & 1304.1277

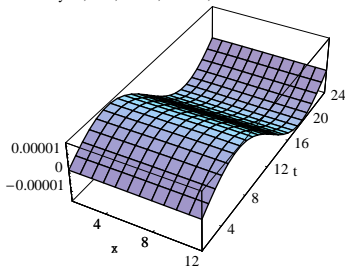
→ 't Hooft 1979, Nielsen, Ambjorn, Olesen, Cornwall, 1979
Mack, 1980; Feynman, 1981

- QCD vacuum is a condensate of closed magnetic flux-lines, they have topology of tubes (3D) or surfaces (4D),
- magnetic flux corresponds to the center of the group,
- Vortex model may explain ...
 - **Confinement** → piercing of Wilson loop \equiv crossing of static electric flux tube and moving closed magnetic flux
 - **Topological charge**: intersection points, writhing points and color structure
 - Engelhardt, Reinhardt (2000), Jordan, R.H., Faber, Heller (2007)
 - **Spontaneous chiral symmetry breaking**: also center-projected configurations show χ SB
 - R.H., Faber, Greensite, Heller, Olejnik (2008)

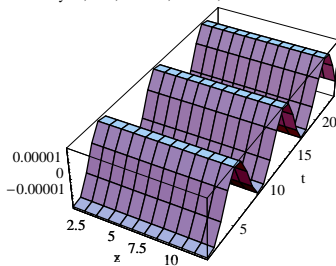
- Instanton Liquid Model:
 - action minima localized in space-time
 - carry $Q = \pm 1$ and attract zero modes according to Atiyah-Singer index theorem
 - overlapping would-be zero modes lead to near-zero modes
 - chiral symmetry breaking via Banks-Casher relation
 - *Diakonov, Petrov (1984)*
- spherical vortices behave like instantons
 - *Schweigler, R.H., Faber, Heller (2012)*
- also intersection points give chiral condensate
 - *R.H., Faber, Heller, Schweigler (2013)*

Overlap fermions on trivial gauge configurations:

$y=6, z=6, \text{chi}=0, n=1-1, \text{max}=0.0000149759$



$y=6, z=6, \text{chi}=-1, n=9-9, \text{max}=0.0000174434$



$$\rho_5 = \psi^\dagger_\pm \gamma_5 \psi_\pm = \frac{1}{2} (\chi^\dagger_R \chi_R - \chi^\dagger_L \chi_L) = \rho_+ - \rho_-$$

Thick Spherical SU(2)-Vortices

Chiral
Symmetry
Breaking
from Center
Vortices

Roman
Höllwieser

Vortices

χ SB

Free Fermions

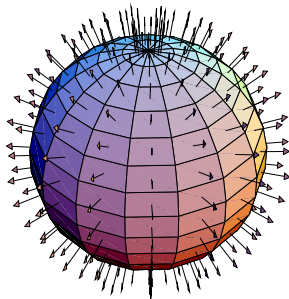
Spherical
Vortex

Interactions

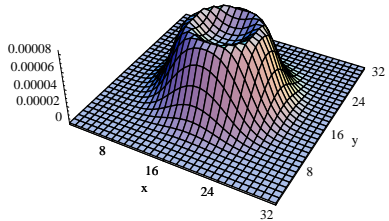
Plane Vortices

Asqtad Modes

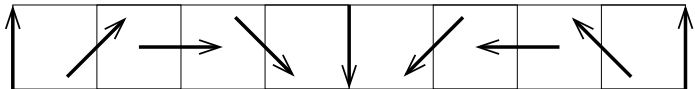
Conclusions

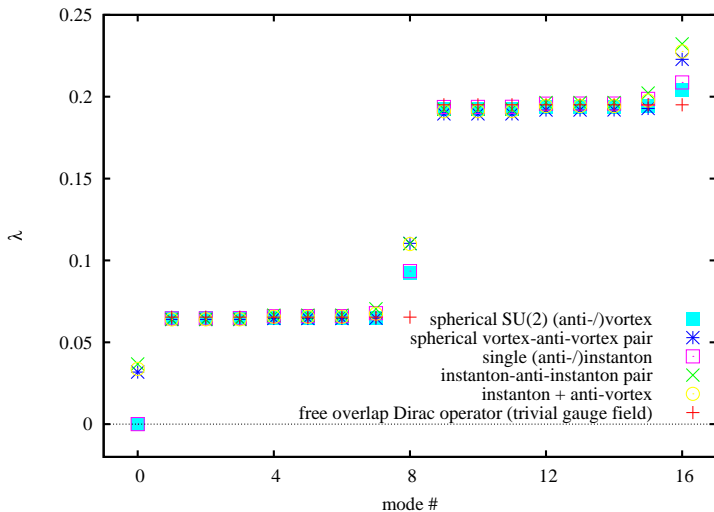


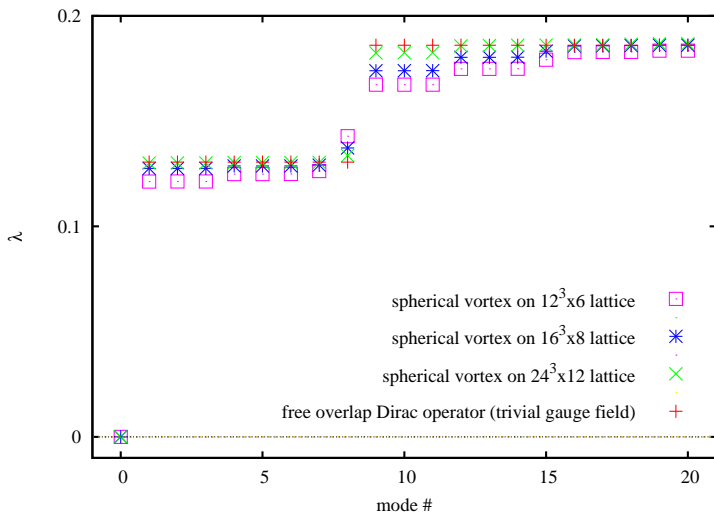
$z=16, t=1, \text{chi}=-1, n=0-0, \text{max}=0.0000822192$



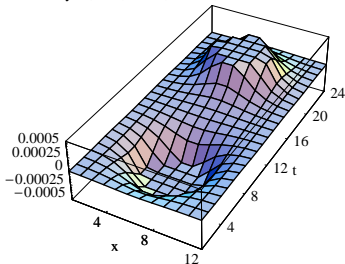
$$U_{\mu}(\vec{r}) = \begin{cases} \exp\{i\alpha(r)\frac{\vec{r}}{r}\vec{\sigma}\}, & t = 1, \mu = 4 \\ 1 & \text{else} \end{cases}$$



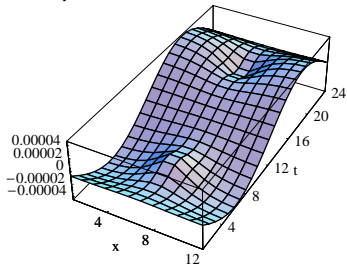




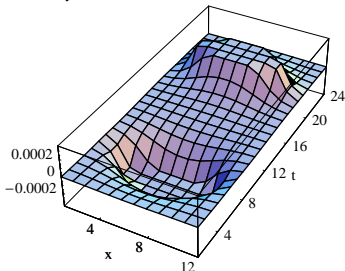
$y=6, z=6, \chi=0, n=0-0, \max=0.00212786$



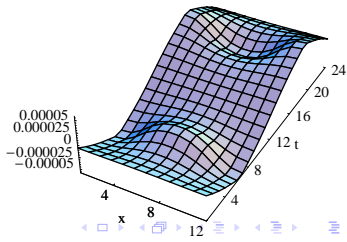
$y=6, z=6, \chi=0, n=1-1, \max=0.000043833$

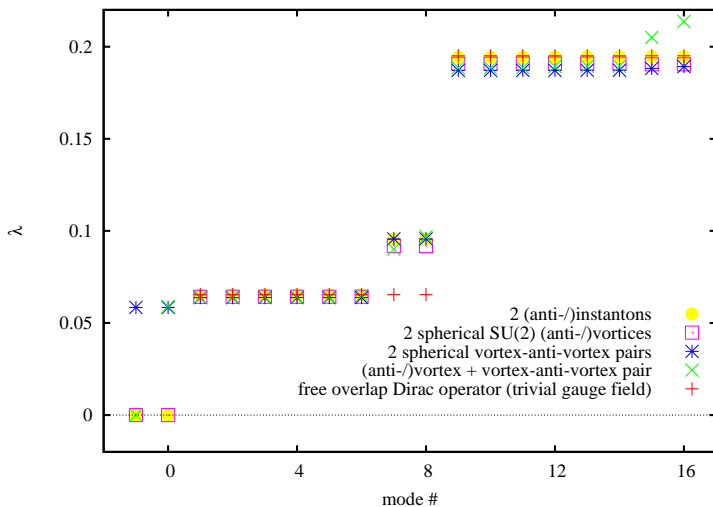


$y=6, z=6, \chi=0, n=0-0, \max=0.0012622$



$y=6, z=6, \chi=0, n=1-1, \max=0.000055155$





Thick, Planar SU(2)-vortices

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Free Fermions

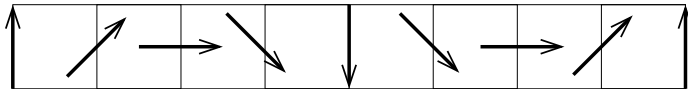
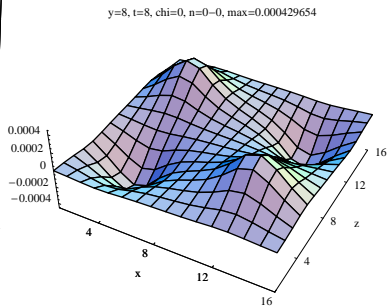
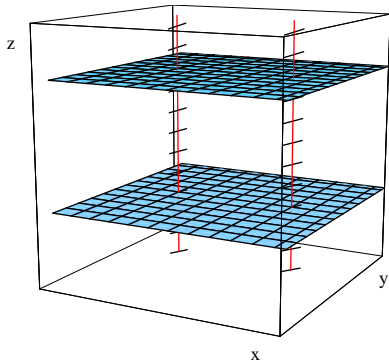
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Vortex

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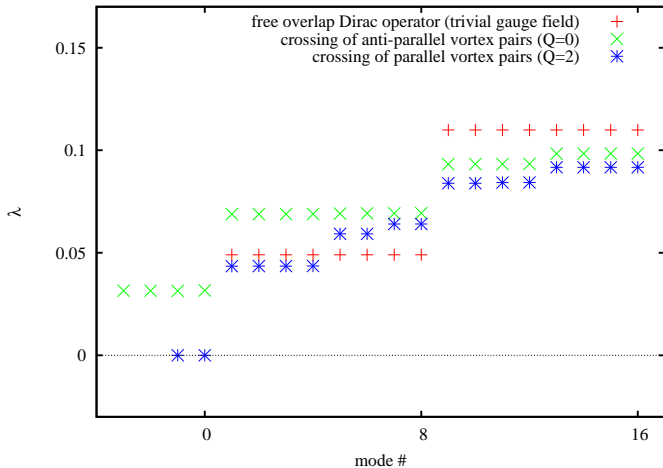
Plane Vortices

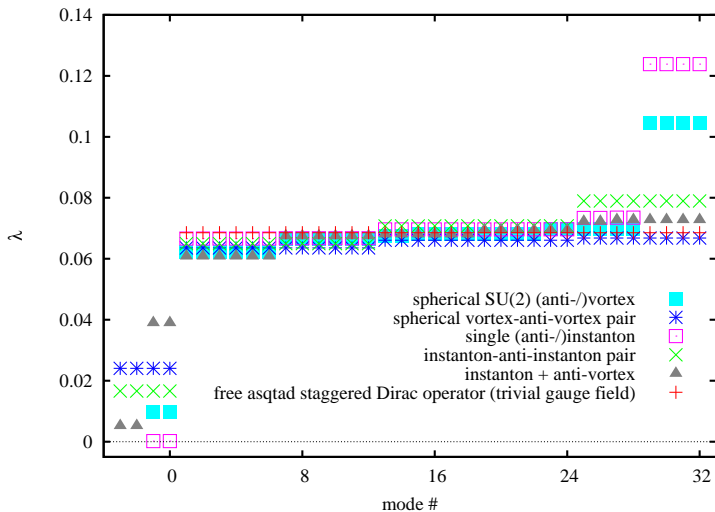
Asqtad Modes

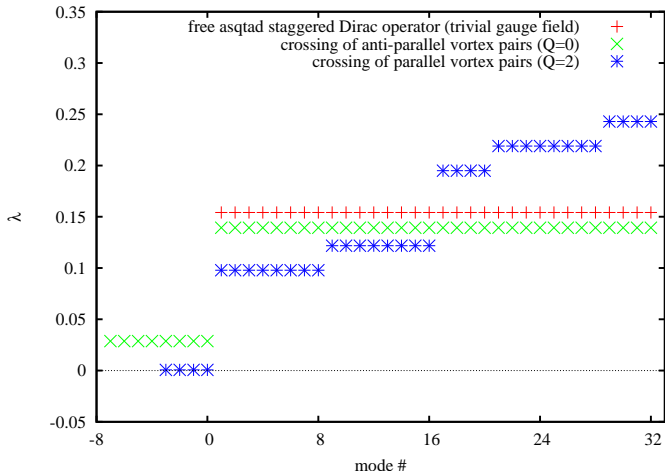
Conclusions



Overlap modes for plane vortices







Chiral Densities of Asqtad Modes

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Symmetry
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Vortices

χ_{SB}

Free Fermions

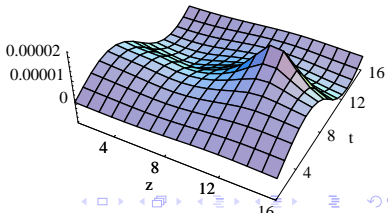
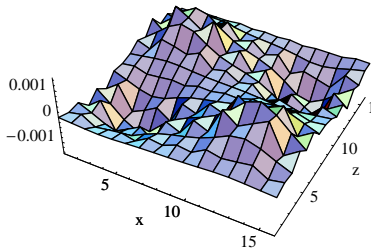
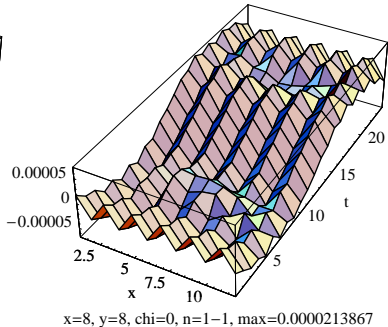
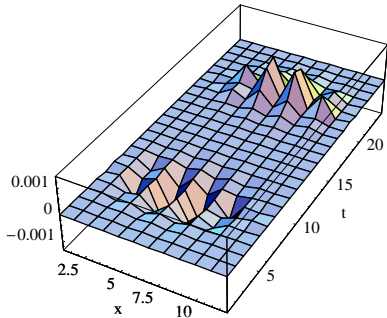
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Conclusions



- easy identification of near-zero modes for classical configs
- similar fermion modes for instantons and spherical vortices
- instanton liquid model can be applied to spherical vortices
- also vortex intersections contribute to chiral condensate
- not an exclusive picture of chiral symmetry breaking
- any source of topological charge can contribute (monopoles, instantons, merons, bions, calorons,...)
- random interactions of quarks with the vortex background
- confining interaction by itself could break chiral symmetry

Thank you for your attention!
Questions?

